

# Monte Carlo Methods for Rough Energy Landscapes

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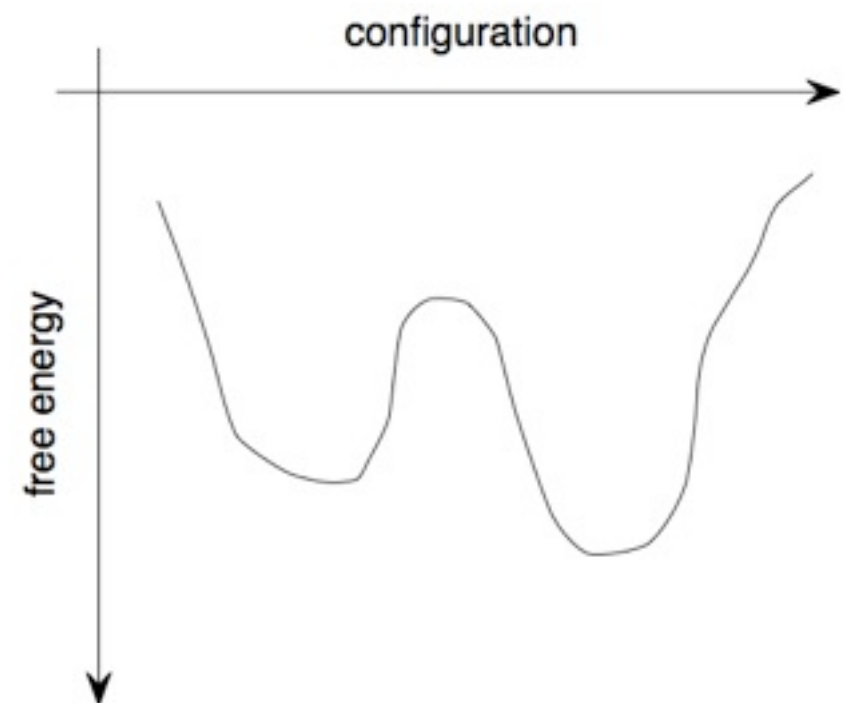
SFI, August 9, 2013

# Collaborators

- Stephan Burkhardt
- Chris Pond
- Jingran Li
- Richard Ellis

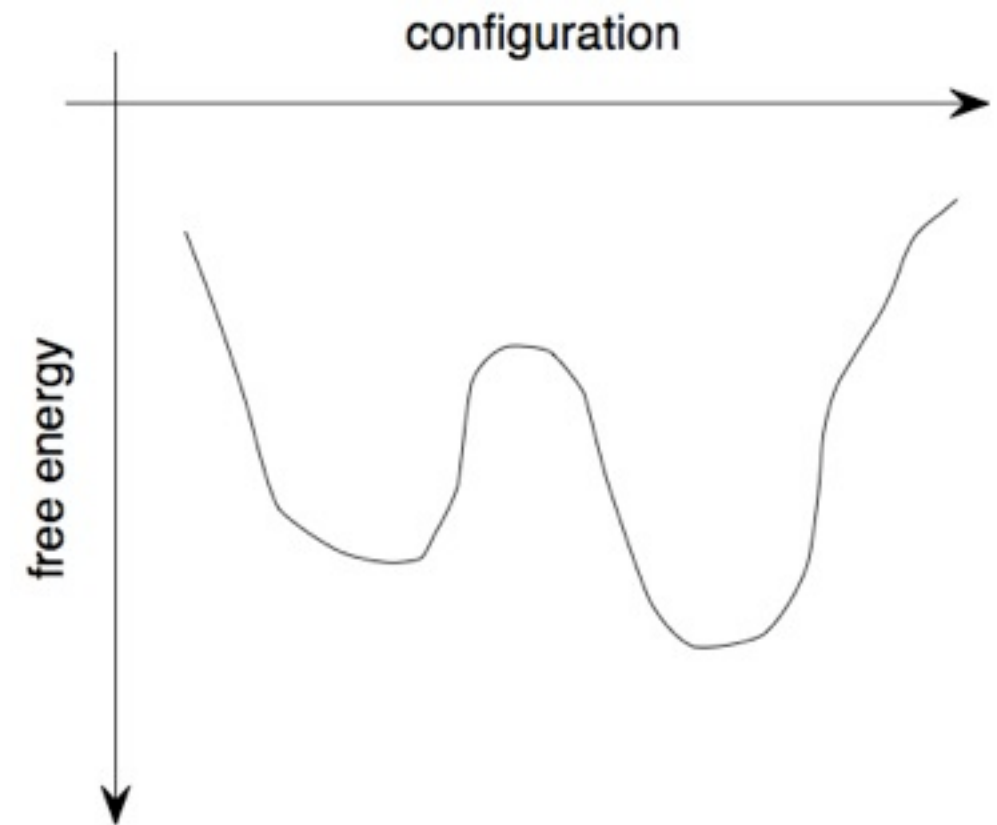
# Motivation

- How to sample equilibrium states of systems with rugged free energy landscapes, e.g. spin glasses, configurational glasses, proteins, combinatorial optimization problems.



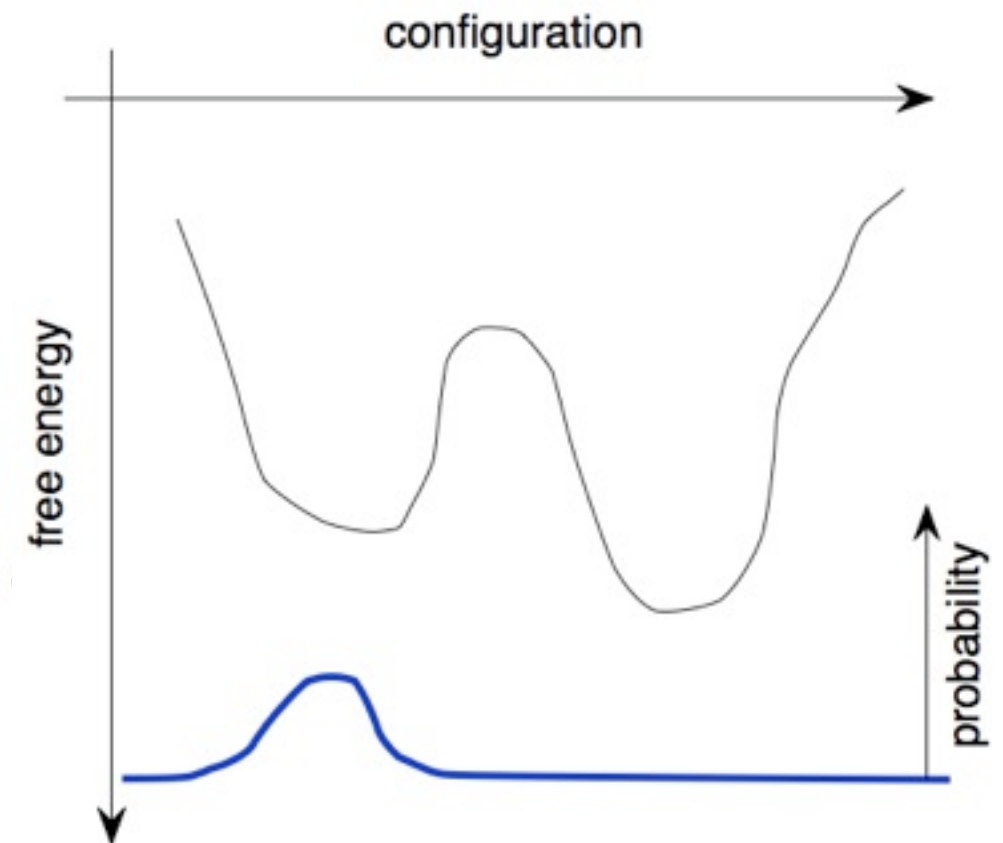
# Problem

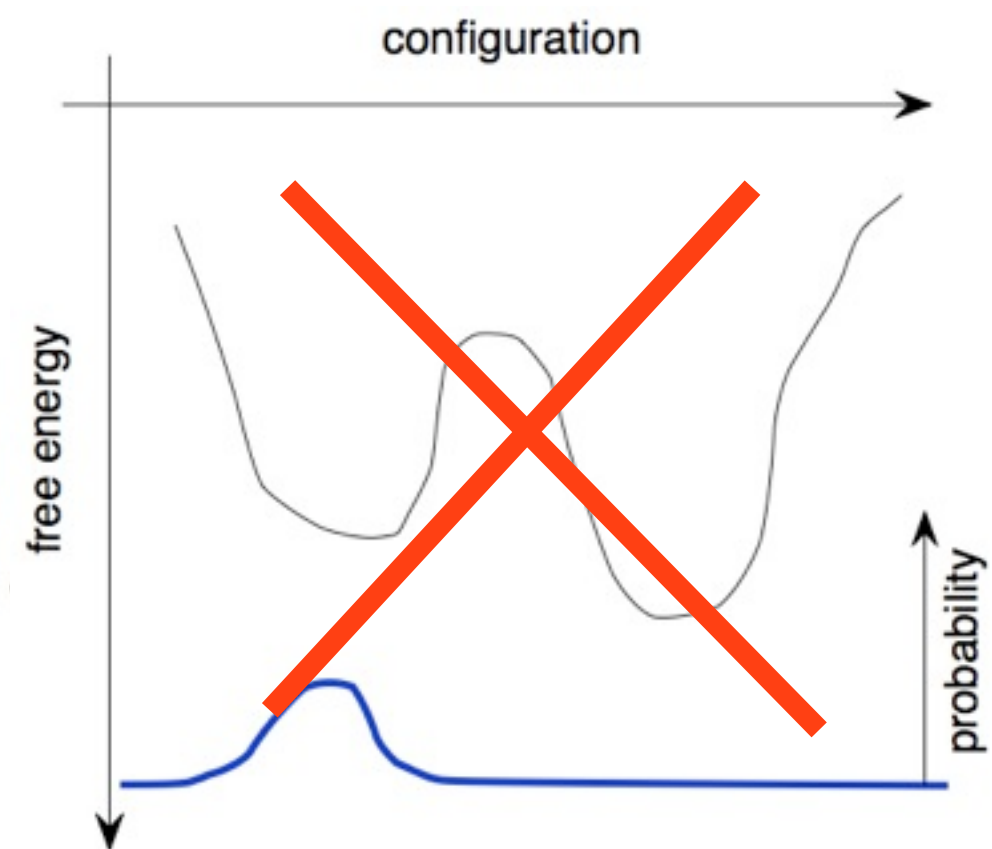
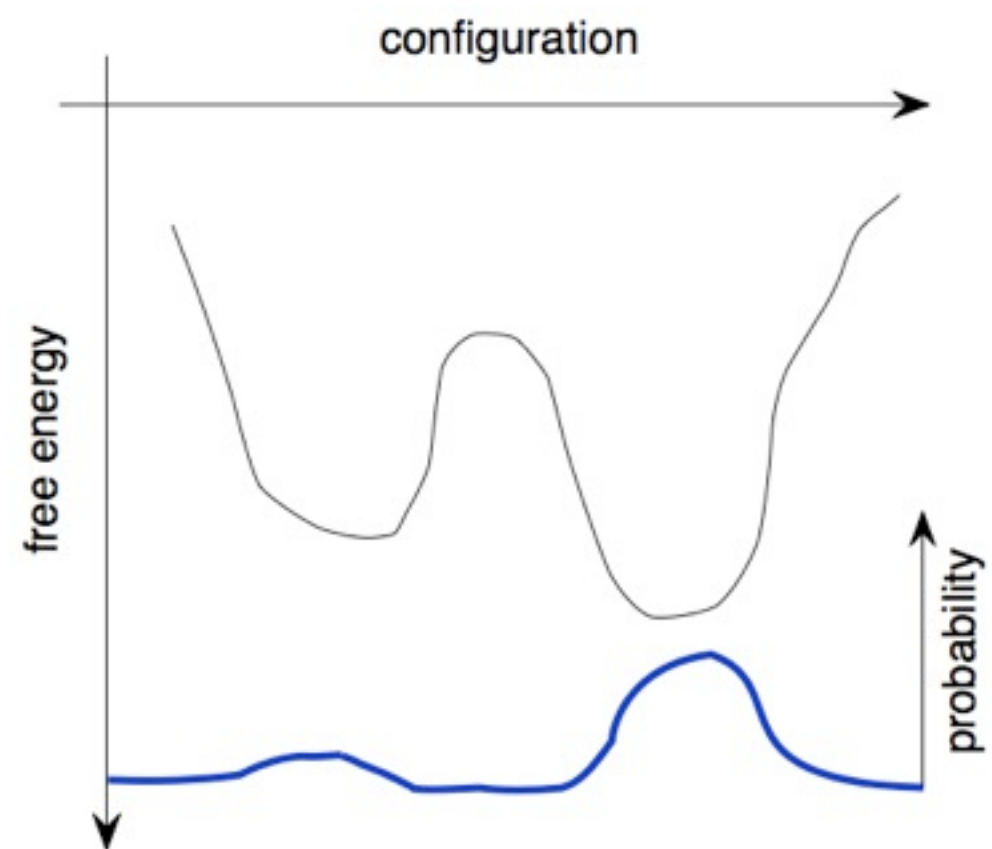
- Markov chain Monte Carlo (MCMC) at a single temperature such as the Metropolis algorithm gets stuck in local minima.



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# Outline

- **Population Annealing**
- Parallel Tempering
- PA vs PT
- Conclusions



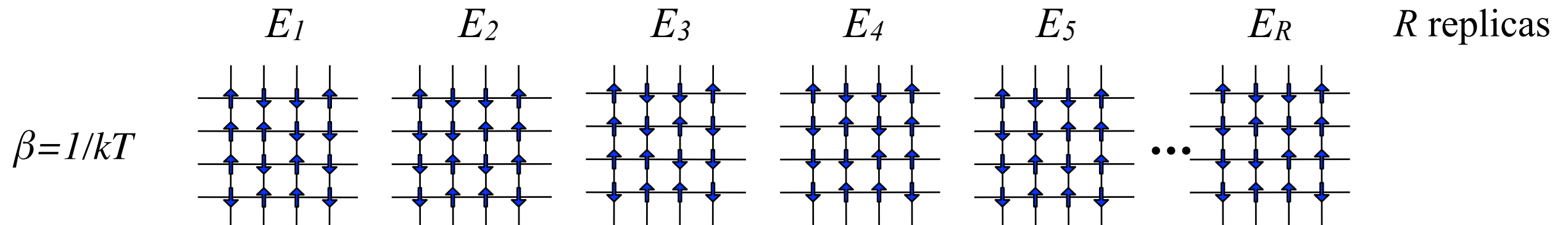
# Population Annealing

K. Hukushima and Y. Iba, in *THE MONTE CARLO METHOD IN THE PHYSICAL SCIENCES: Celebrating the 50th Anniversary of the Metropolis Algorithm*, edited by J. E. Gubernatis (AIP, 2003), vol. 690, pp. 200–206.

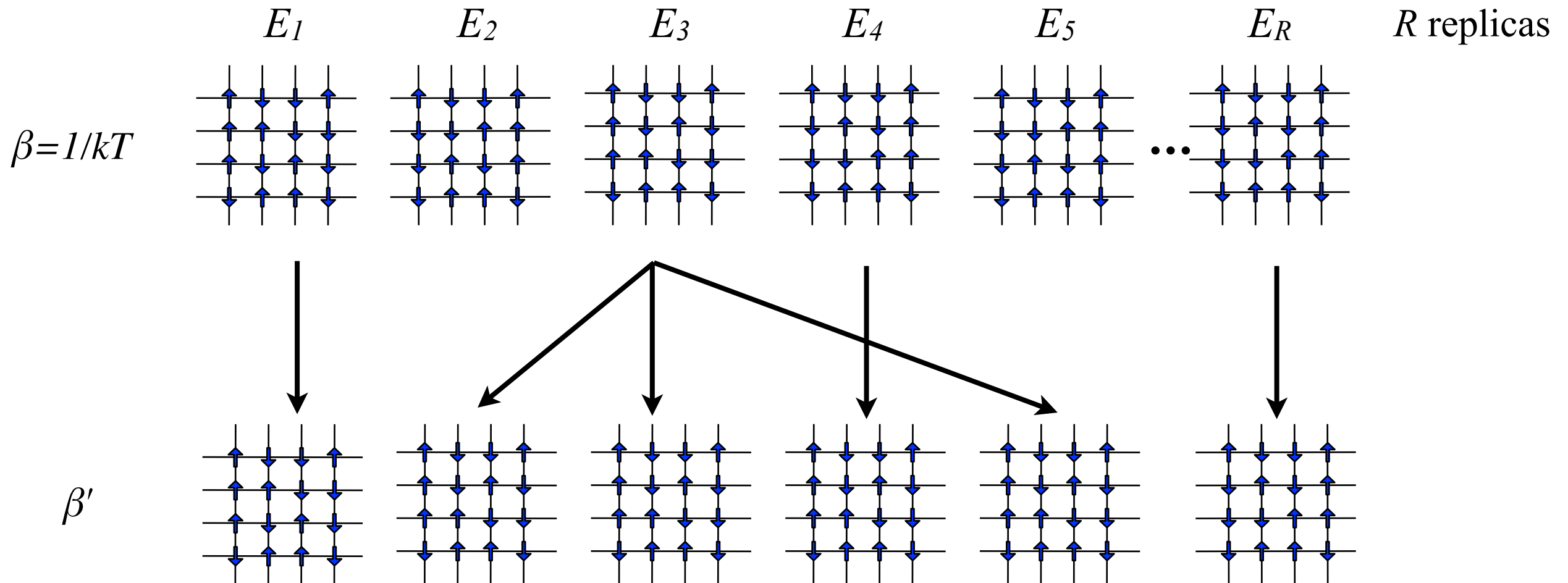
- Modification of simulated annealing for equilibrium sampling.
- A population of replicas is cooled according to an annealing schedule. Each replica is acted on by a MCMC (e.g. Metropolis) at the current temperature.
- During each temperature step, the population is differentially resampled according to Boltzmann weights to maintain equilibrium.



# Population Annealing

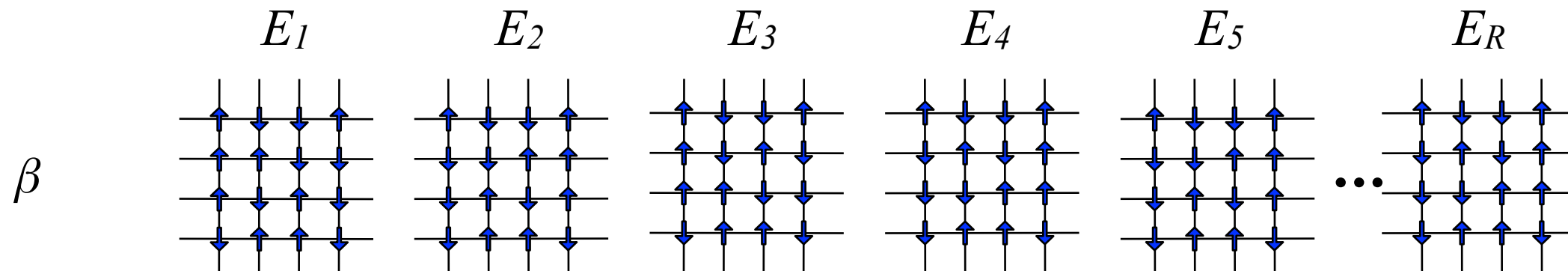


# Population Annealing



Population annealing = simulated annealing with differential reproduction (resampling) of replicas

# Temperature Step

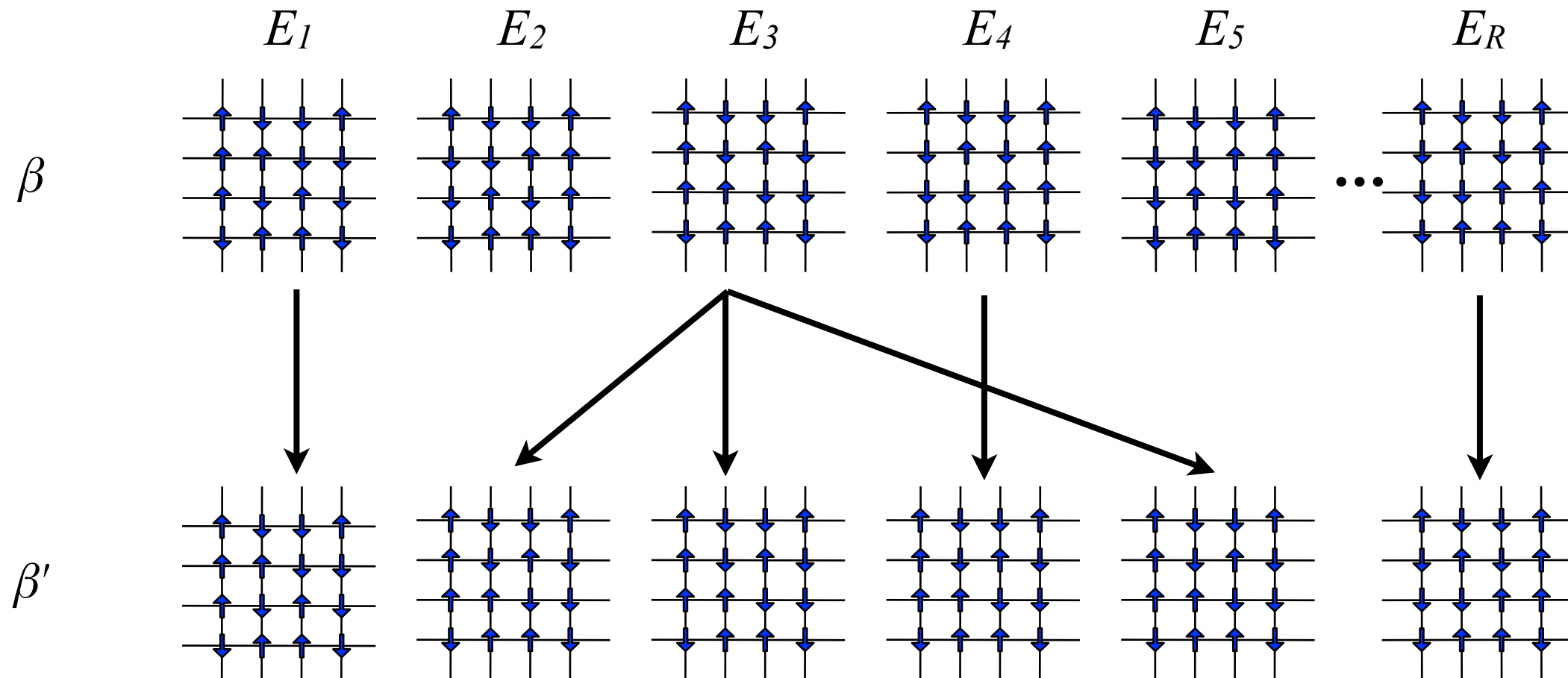


$$\tau_j(\beta, \beta') = \frac{\exp [-(\beta' - \beta)E_j]}{Q(\beta, \beta')}$$

$$Q(\beta, \beta') = \frac{\sum_{j=1}^R \exp [-(\beta' - \beta)E_j]}{R}$$

Replica  $j$  is reproduced  $n_j$  times where  $n_j$  is an integer random variate with mean  $\tau_j$ .

# Temperature Step



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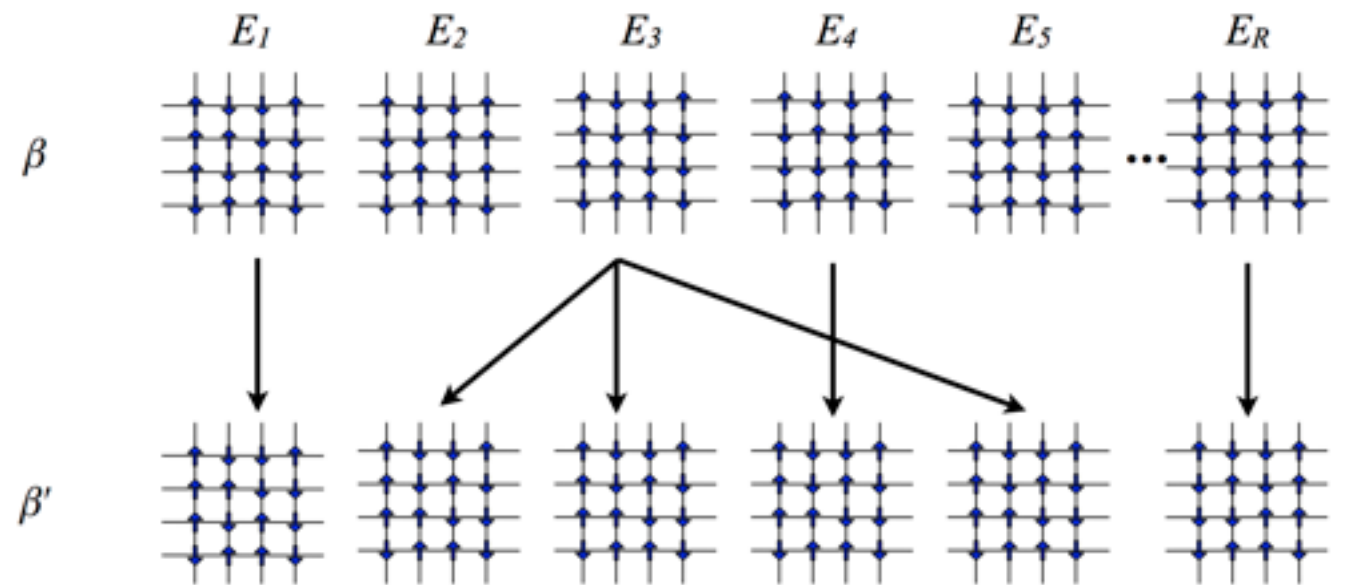
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# Resampling

e.g.  $n_1 = 1, n_2 = 0, n_3 = 3, n_4 = 1$

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Replica  $j$  is reproduced  $n_j$  times where  $n_j$  is an integer random variate with mean  $\tau_j$ .

## Nearest integer resampling

$$\text{Prob}(n = \lfloor \tau \rfloor) = \tau - \lceil \tau \rceil$$

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$$\text{Prob}(n - \tau > 1) = 0$$

Minimizes variance of  $n_j$ .

Maximizes number of ancestors with descendants.

Population size fluctuates.

# Population Annealing is related to...

◆ Simulated annealing

◆ Sequential Monte Carlo

- See e.g. “Sequential Monte Carlo Methods in Practice”, A. Doucet, et. al. (2001)
- aka Particle Filters
- Nested Sampling, Skilling
- Go with the Winners, Grassberger (2002)
- Diffusion (quantum) Monte Carlo
- Nonequilibrium Equality for Free Energy Differences, Jarzynski (1997)
- Histogram Re-weighting, Swendsen and Ferrenberg (1988)

# Convergence to Equilibrium

Suppose the population at  $\beta$  is a correct, iid sample:

$$Y_i = \exp [-(\beta' - \beta)E_i + \lambda A_i]$$

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Suppose the population at  $\beta$  is a correct, iid sample:

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A bias inversely proportional to the population size results from the re-weighting to  $\beta'$

# Convergence to Equilibrium

Resampling complicates the analysis...

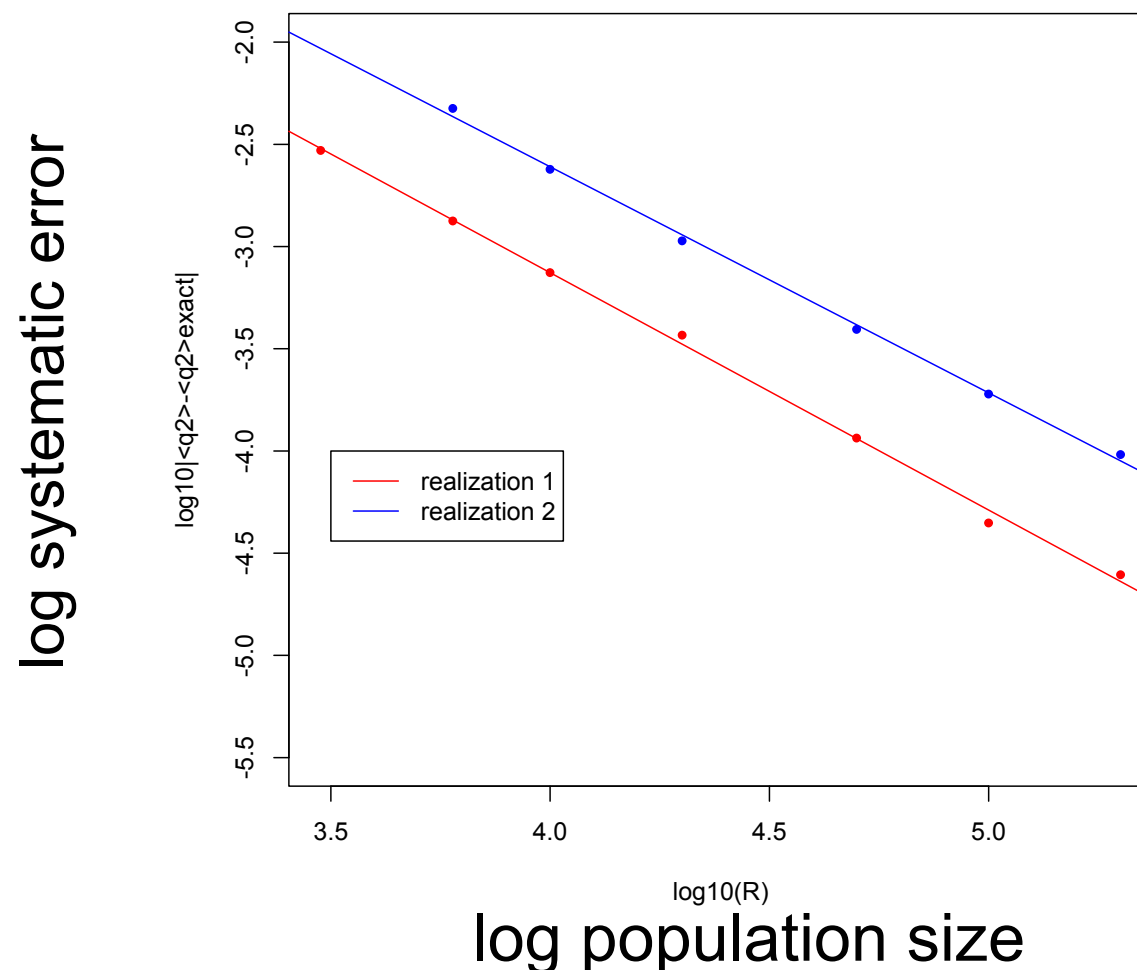
Conjecture: Fixing other parameters of population annealing, observable converge to their equilibrium values inversely in population size.

# Quantifying Convergence to Equilibrium

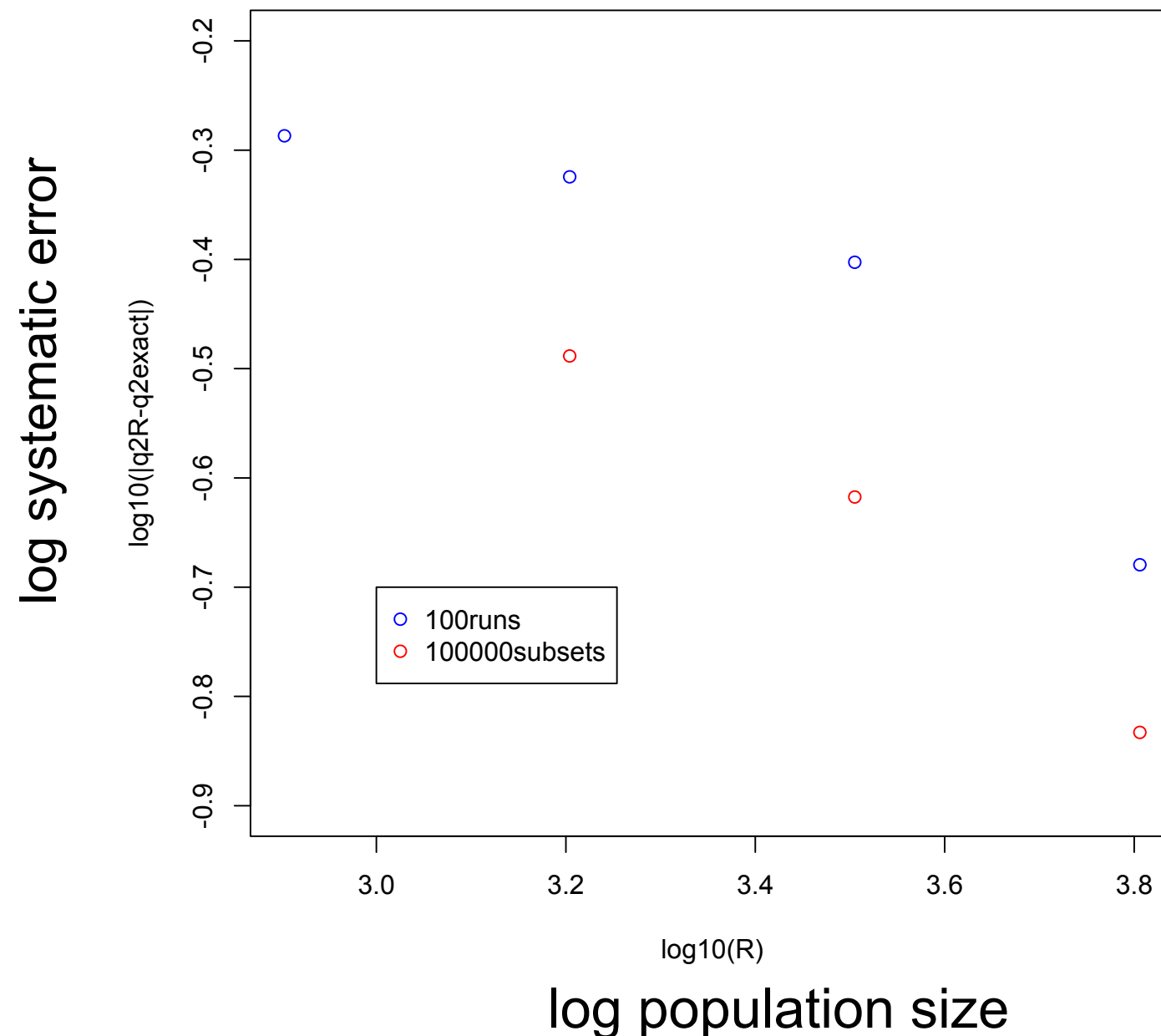
- Do many runs for each of many population sizes. Expensive.
- Is there an analog to the autocorrelation function and the exponential autocorrelation time?

# Quantifying Convergence to Equilibrium

- Idea: resample small populations from a single run with a large population (bootstrap).
  - Choose a small population of parents and measure the observable in their descendants.
  - Repeat many times for different populations sizes.



- Problem: A small sample within a large population is not the same as an isolated small population because the sample within the large populations exists in a more competitive environment.



● =bootstrap  
● =many runs



# Direct Estimate of Free Energies

$$Q(\beta, \beta') = \frac{\sum_{j=1}^R \exp [-(\beta' - \beta) E_j]}{R}$$

$$-\beta_k F(\beta_k) = \sum_{\ell=k}^K \ln Q(\beta_{\ell+1}, \beta_{\ell}) + \beta_K F(\beta_K)$$

**Derivation:**

$$\begin{aligned} \frac{Z(\beta')}{Z(\beta)} &= \frac{\sum_{\gamma} e^{-\beta' E_{\gamma}}}{Z(\beta)} \\ &= \sum_{\gamma} e^{-(\beta' - \beta) E_{\gamma}} \left( \frac{e^{-\beta E_{\gamma}}}{Z(\beta)} \right) \\ &= \langle e^{-(\beta' - \beta) E_{\gamma}} \rangle_{\beta} \\ &\approx \frac{1}{R} \sum_{j=1}^R e^{-(\beta' - \beta) E_j} = Q(\beta, \beta'). \end{aligned}$$

# Weighted Averages

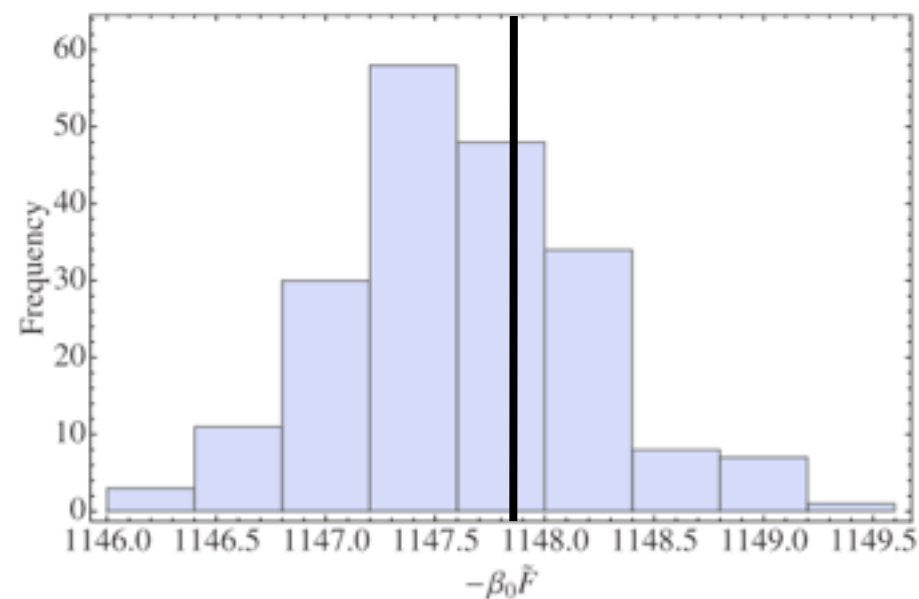
JM, PRE82,26704(2010)

- Results using small population size are biased.
- Results from independent runs can be combined and biases reduced using *weighted* averages.
- Observables from each run weighted by the exponential of the free energy estimator for that run.

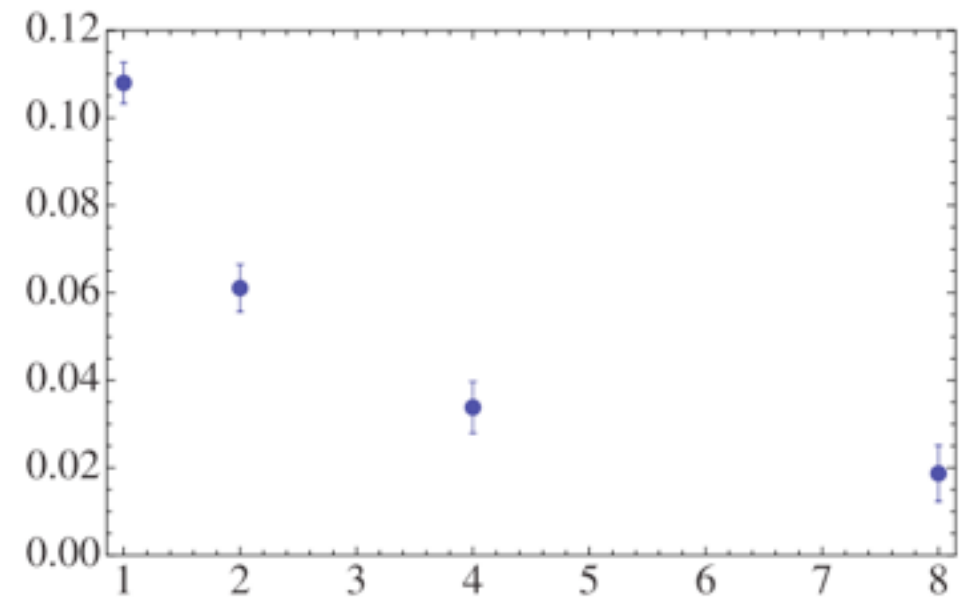
$$\overline{A}(\beta) = \sum_{m=1}^M \tilde{A}_m(\beta) \omega_m(\beta) \quad \omega_m(\beta) = \frac{e^{-\beta \tilde{F}_m(\beta)}}{\sum_{i=1}^M e^{-\beta \tilde{F}_i(\beta)}}$$

# Weighted Averages

- Number of runs needed for unbiased results determined by the width of the distribution of the free energy estimator--an intrinsic measure of equilibration.



systematic error

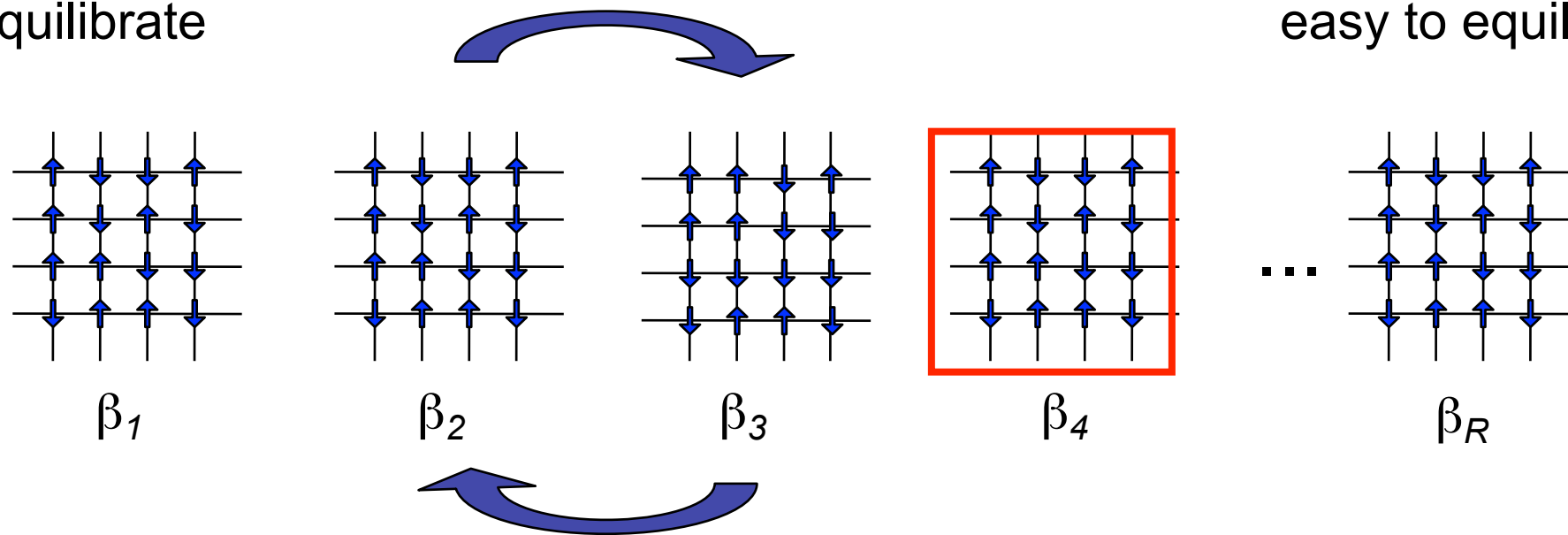


number of runs

# Parallel Tempering

hard to equilibrate

easy to equilibrate

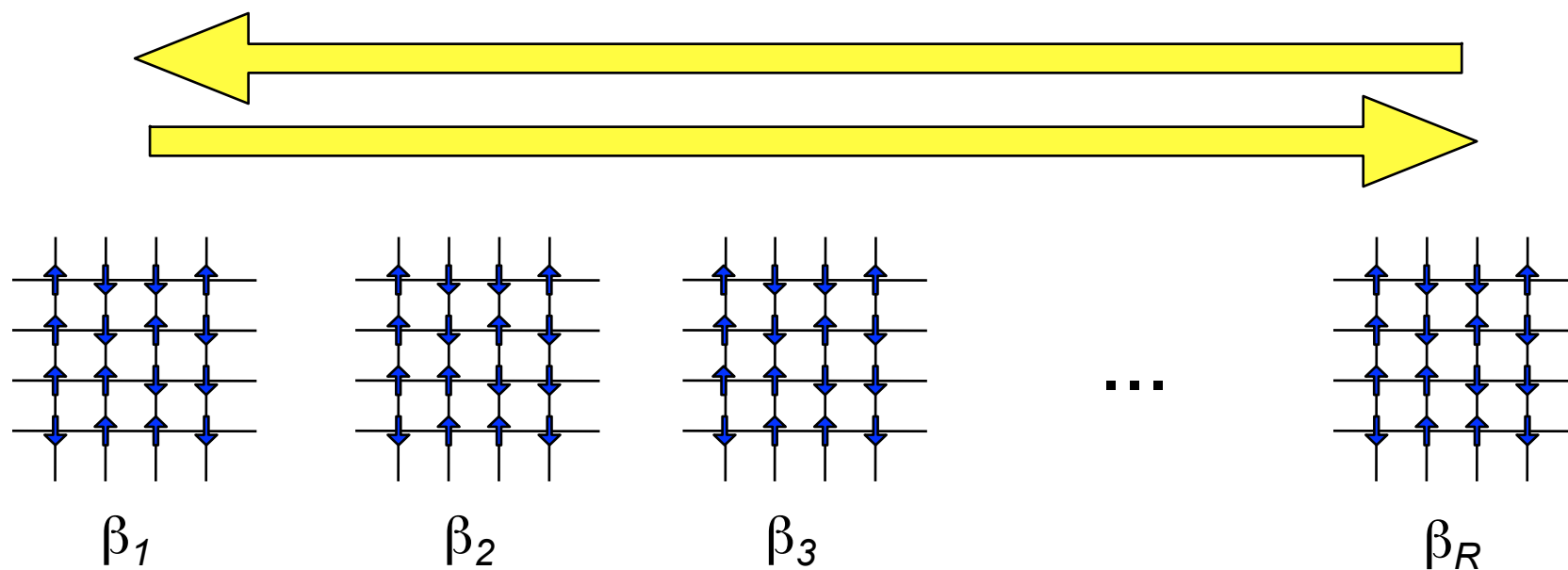


- $R$  replicas at inverse temperatures  $\beta_1 > \beta_2 > \dots > \beta_R$  (each with the same couplings).
- MCMC (e.g. Metropolis) on each replica
- Exchange replicas with energies  $E$  and  $E'$  and temperatures  $\beta$  and  $\beta'$ , with probability:

$$p_{\text{swap}} = \min \left[ 1, e^{(\beta - \beta')(E - E')} \right]$$

# Intuition

- Mixing is accelerated by “round trips” from low to high temperature and back.

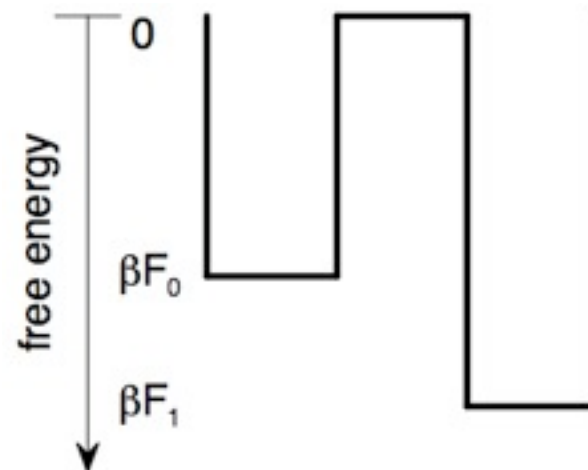


Population Annealing	Parallel Tempering
Sequential Monte Carlo	Markov Chain Monte Carlo
# Replicas ( $R \gg 1$ ) <i>space</i>	#Sweeps ( $t \gg 1$ ) <i>parallel time</i>
#Temperature steps ( $T$ ) <i>parallel time</i>	#Replicas ( $R$ ) <i>space</i>
$work = RT$	$work = Rt$
$(A(R) - A_{eq}) \sim R_0/R$	$(A(t) - A_{eq}) \sim e^{(-t/\tau)}$

# Two-Well landscape

- Consider a toy free energy landscape with two free energy minima separated by a high barrier.

– JM, PRE **80**, 056706 (2009)

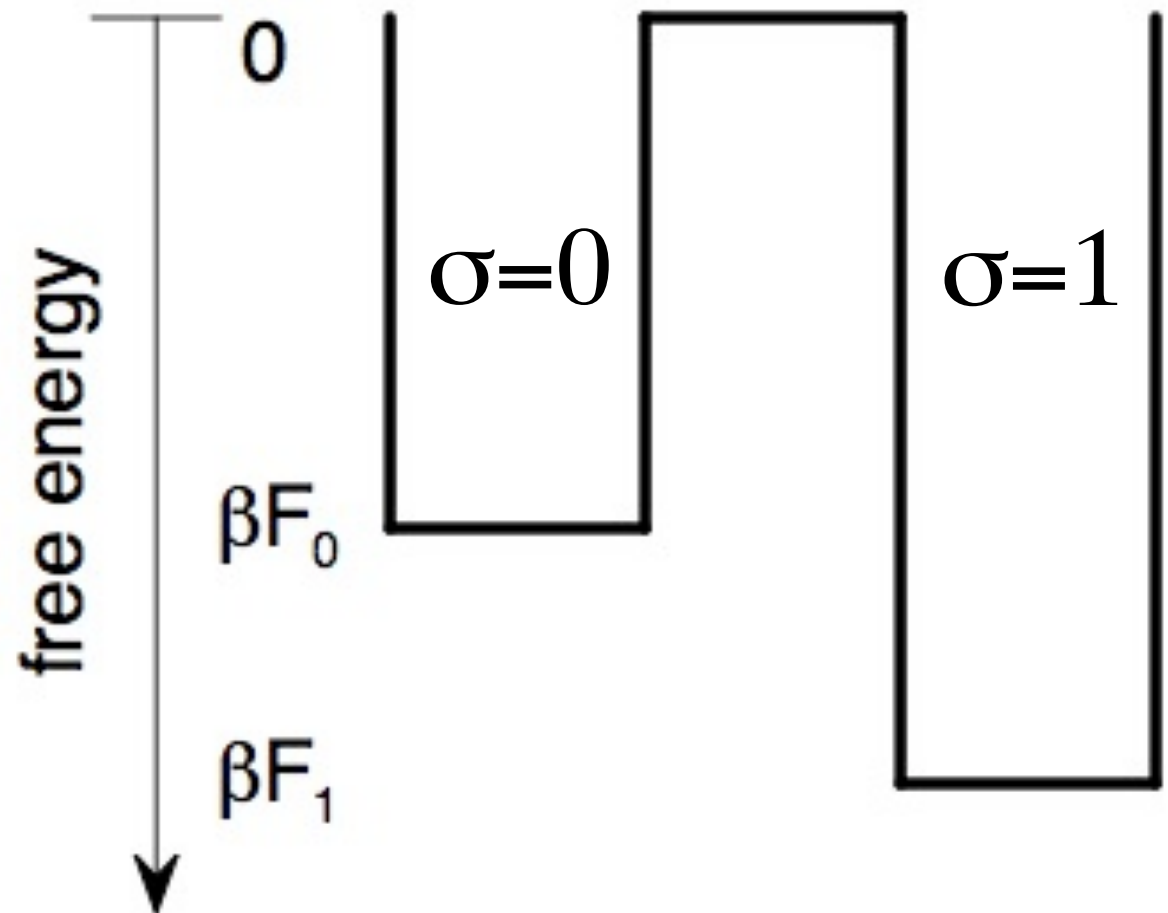


# Two-Well Landscape

$$\beta F_{\sigma}(\beta) = -\frac{1}{2}(\beta - \beta_c)^2(K + H\sigma)$$

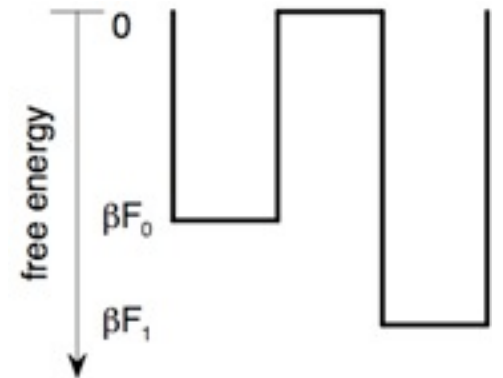
$$\beta \delta F = -\frac{1}{2}(\beta - \beta_c)^2 H$$

$$\text{Prob} [\sigma = +1] = \frac{1}{1 + e^{-\beta \delta F}}$$





# Dynamics of the two-well model



- Assumptions:
  - Fast equilibration within each well.
  - No transitions between wells except at  $\beta_c$  where each well is equally probable.
  - Energy is normally distributed in each well; from thermodynamics:

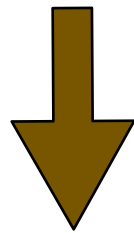
$$\langle E \rangle = -(\beta - \beta_c)(K + H\sigma) \quad \mathbf{Var}(E) = (K + H\sigma)$$

# Replica exchange probabilities

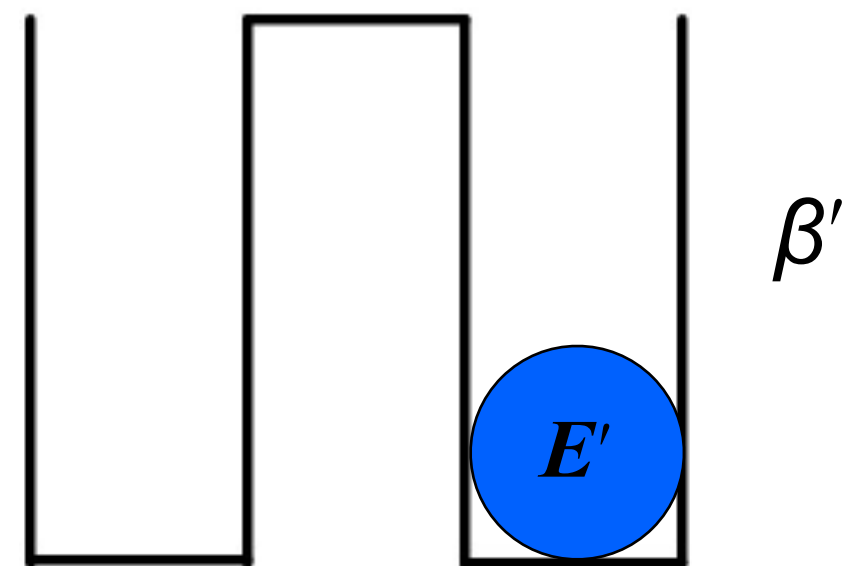
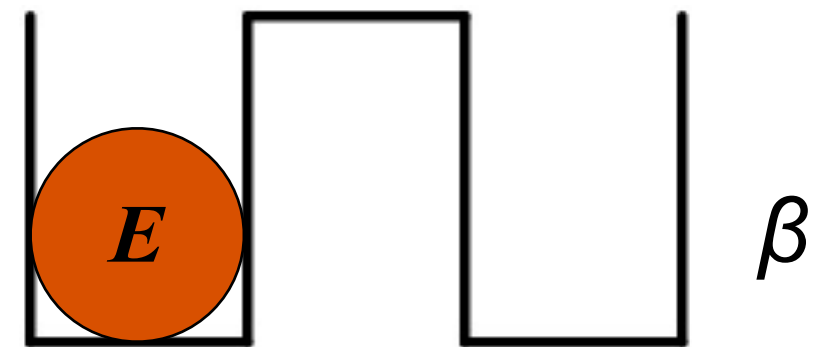
- For the two-well model, replica exchange transition probabilities can be computed exactly. For symmetric wells ( $H=0$ ):

$$p_{\text{swap}} = \min \left[ 1, e^{(\beta - \beta')(E - E')} \right]$$

$$\langle E \rangle = -(\beta - \beta_c)K, \quad \text{Var}(E) = K$$



$$p_{\text{swap}} = \frac{1}{2} \text{Erfc} \left( (\beta - \beta') \sqrt{K} \right)$$

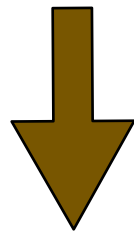


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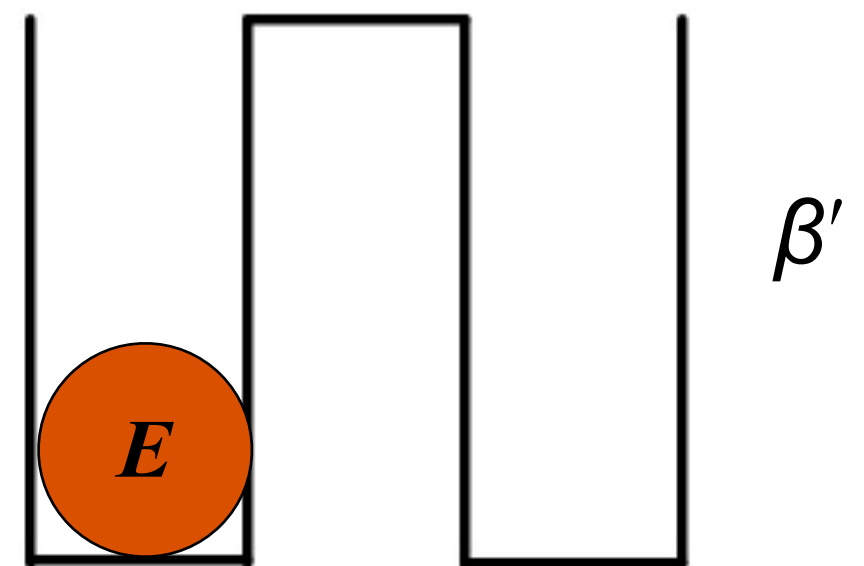
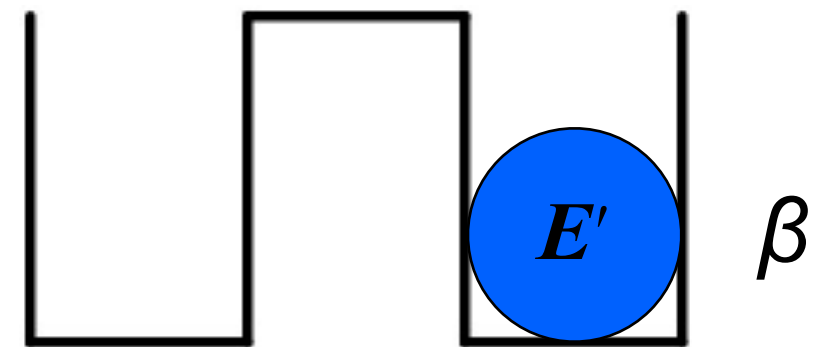
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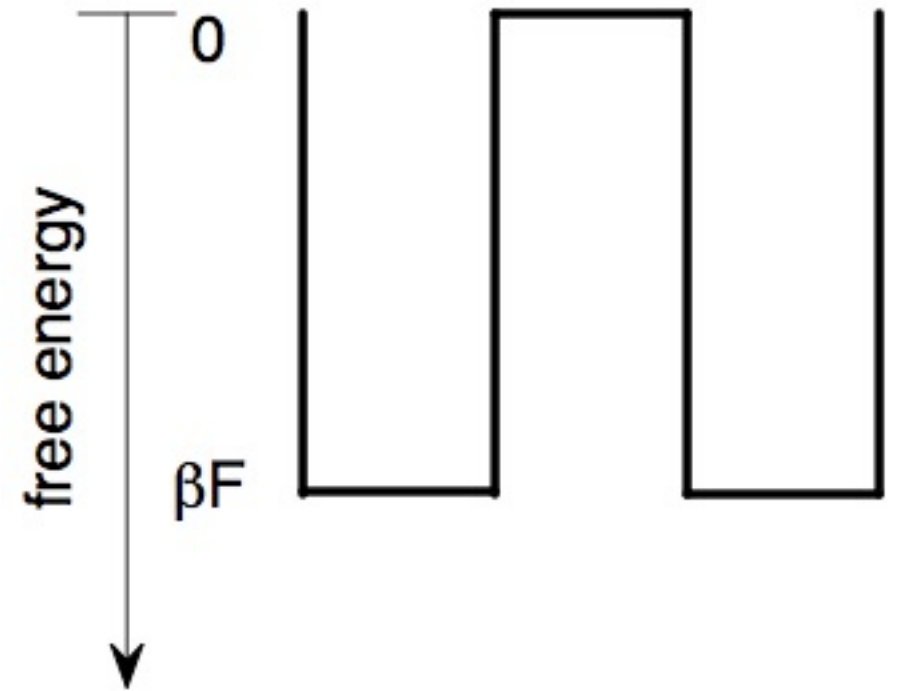


$$p_{\text{swap}} = \frac{1}{2} \text{Erfc} \left( (\beta - \beta') \sqrt{K} \right)$$



# PT for the two-well model

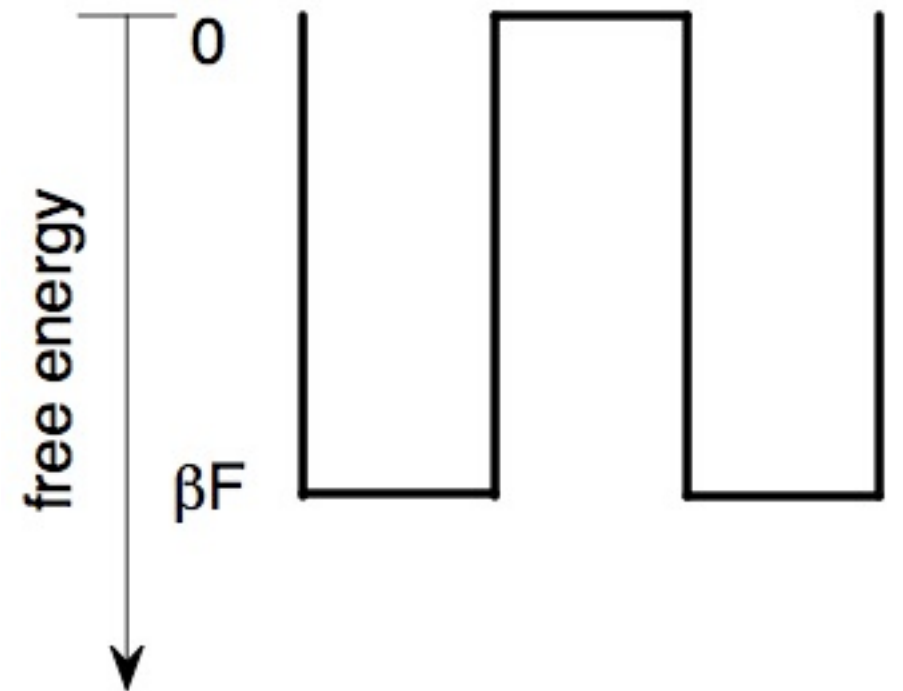
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# PT for the two-well model

- *Diffusion* of replicas

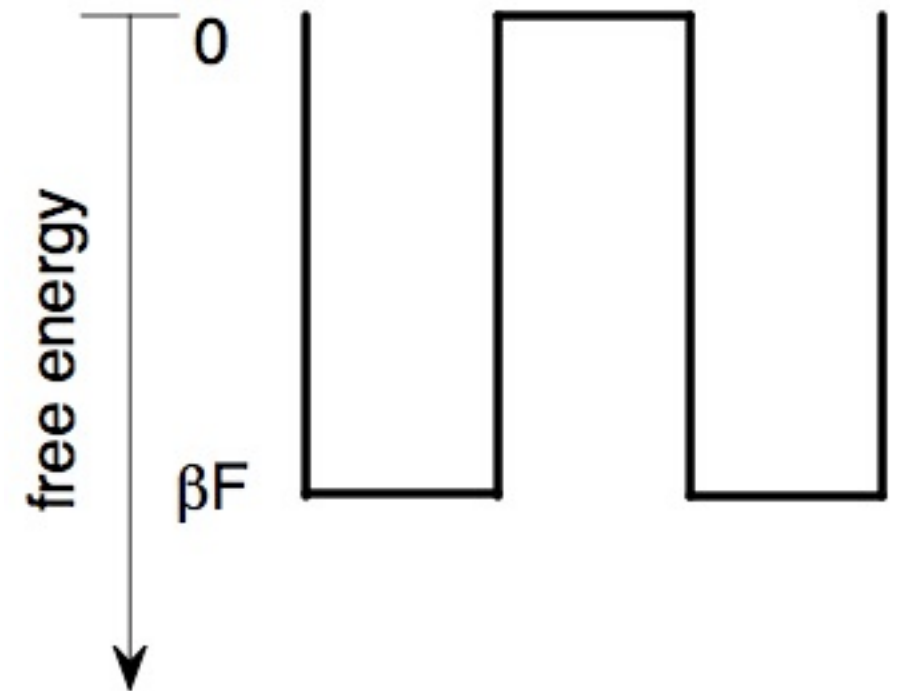
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# PT for the two-well model

- *Diffusion* of replicas
- Equilibration time  $\tau$  is proportional to the mean first passage time for diffusion from  $\beta_0$  to  $\beta_c$  with  $R$  equally spaced replicas

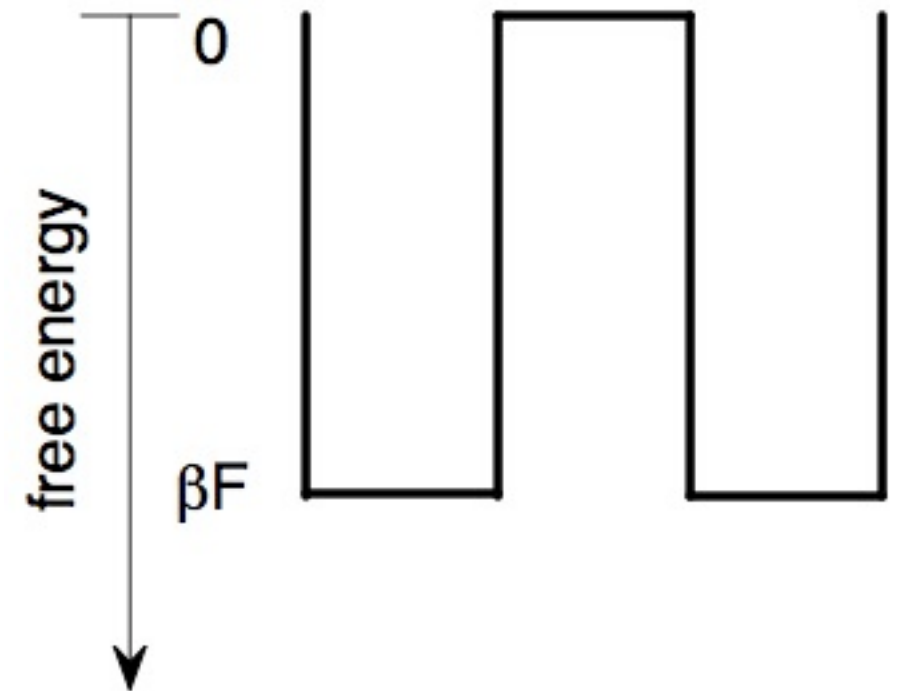
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- Optimum number of replicas balances diffusion time and replica exchange acceptance fraction.

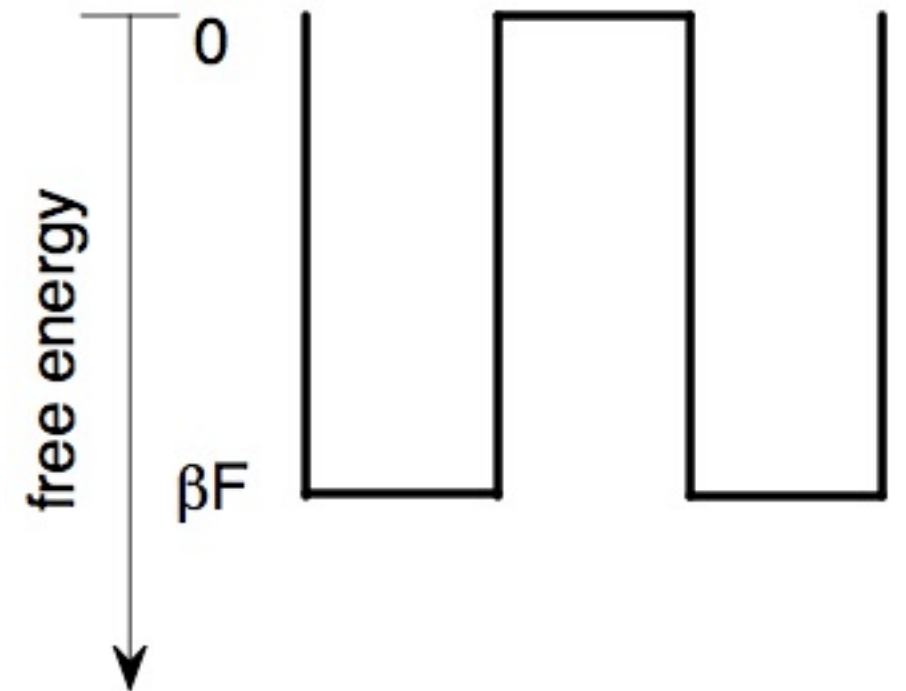
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$$p_{\text{swap}} = \frac{1}{2} \text{Erfc} \left( (\beta - \beta') \sqrt{K} \right)$$



$$R_{\text{opt}} = 1 + 0.594(\beta_0 - \beta_c)\sqrt{K} \sim \sqrt{-\beta F}$$

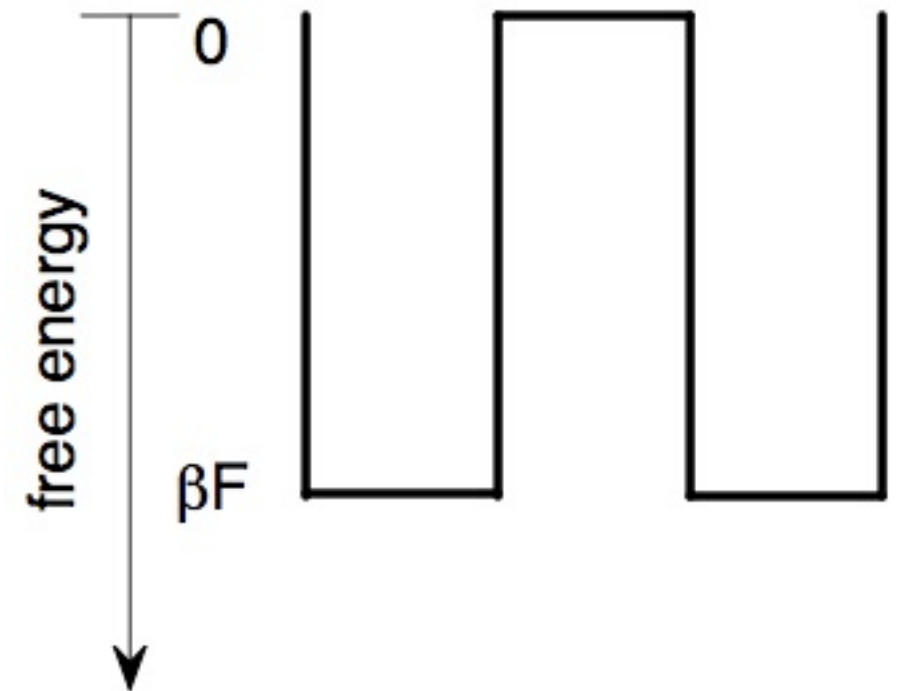


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$$\tau \sim R^2 \sim \beta F$$



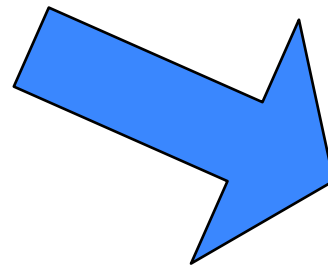
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# PA for the two-well model

$$Y_i = \exp [ -(\beta' - \beta)E_i + \lambda A_i ]$$

$$I(\beta, \beta', \lambda) = \mathbf{E} \log \left( \frac{1}{R} \sum_{j=1}^R Y_j \right) = \log(\mathbf{E}Y) - \frac{1}{2R} \frac{\mathbf{Var}Y}{(\mathbf{E}Y)^2} + O\left(\frac{1}{R^{3/2}}\right)$$

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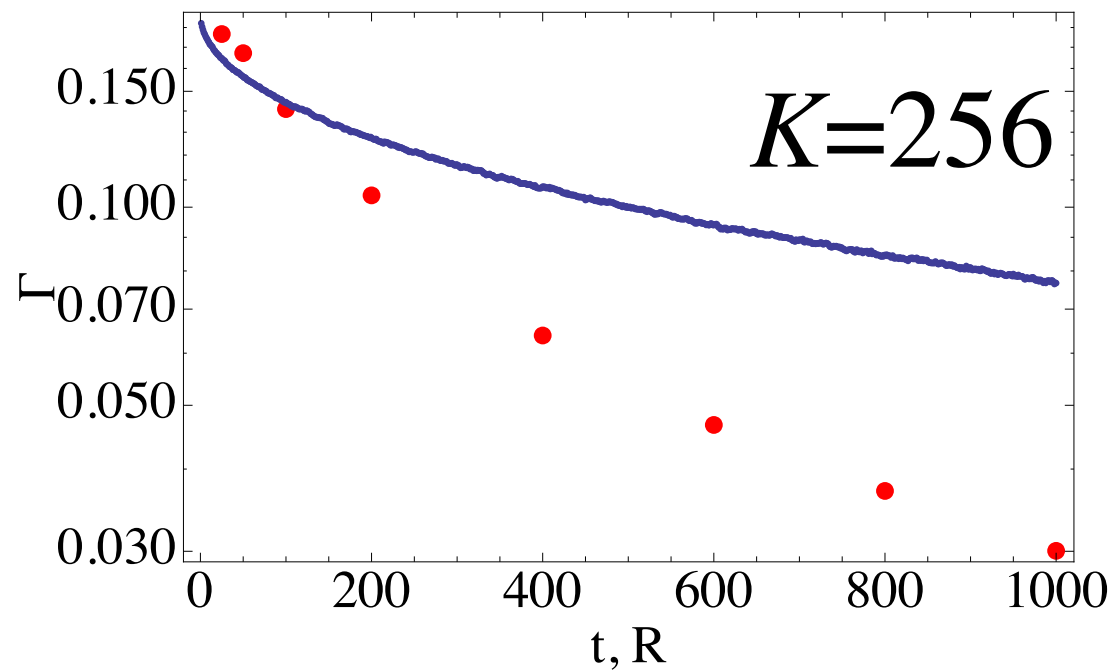
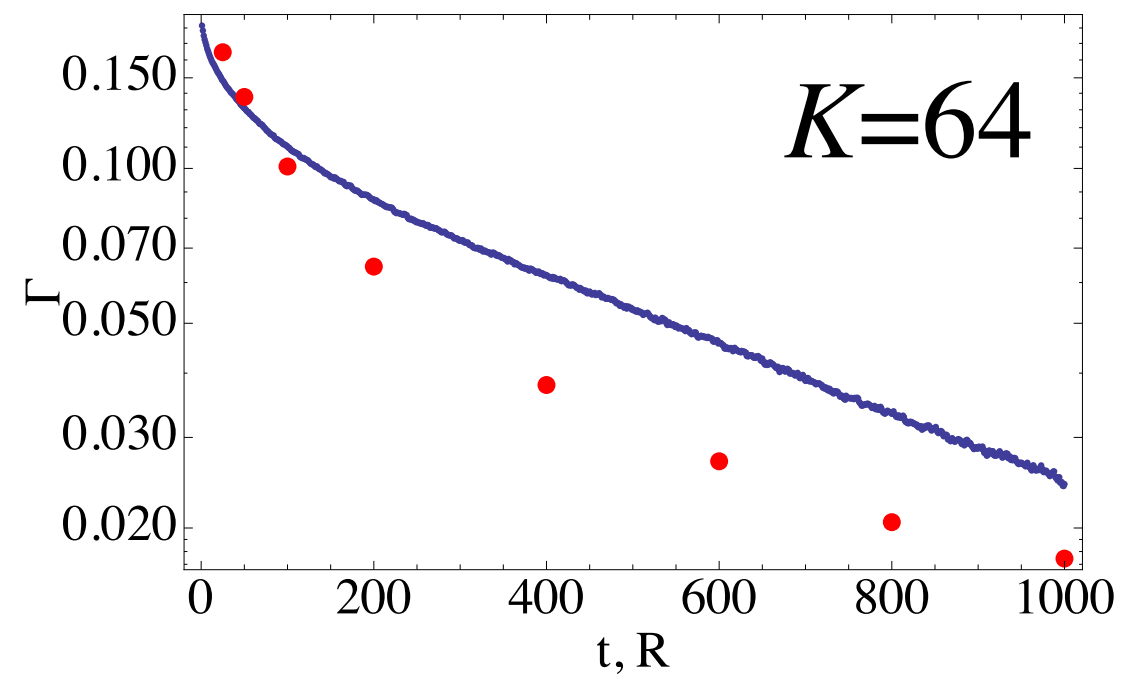
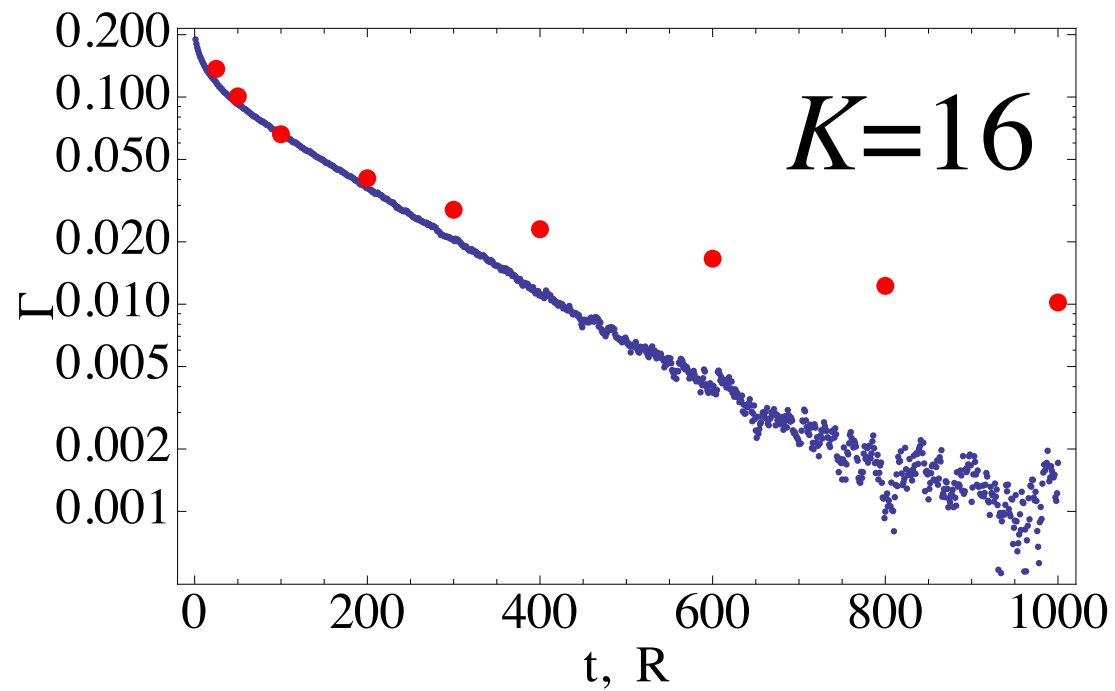
$$(A(R) - A_{eq}) \sim R_0/R$$

$$T \sim K^{1/2}$$

$$R_0 \sim T^a \quad a=1 ?$$

<b>Population Annealing</b>	<b>Parallel Tempering</b>
# Replicas ( $R$ ) <i>space</i>	#Sweeps ( $t$ ) <i>parallel time</i>
#Temperature steps ( $T$ ) <i>parallel time</i>	#Replicas ( $R$ ) <i>space</i>
$(A(R) - A_{eq}) \sim R_0/R$	$(A(t) - A_{eq}) \sim e^{(-t/\tau)}$
$R_0 \sim T \quad T \sim K^{1/2}$	$\tau \sim K \quad R \sim K^{1/2}$
work to eq $\sim R_0 T \sim K$	work to eq $\sim \tau R \sim K^{3/2}$

# Compare PA and PT for the Two-Well Model



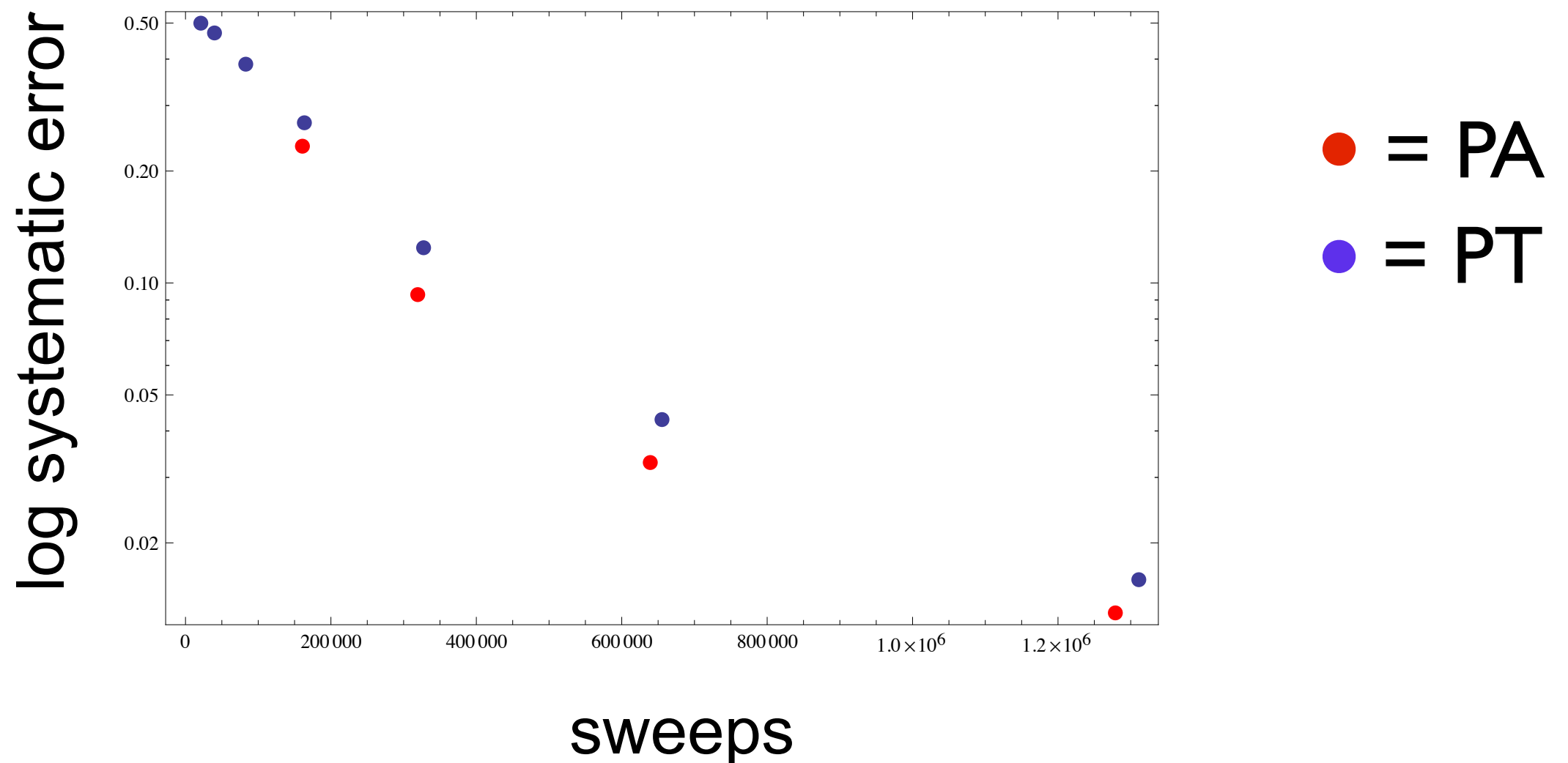
● = PA  
— = PT

$H=0.1$

$$\Gamma = \langle \text{Prob} [\sigma = +1] \rangle - \langle \text{Prob} [\sigma = +1] \rangle_{\text{eq}}$$

# PA vs PT for 3D Spin Glass

- For a small handful of disorder realizations, doing many runs for different population sizes (for PA) or run lengths (for PT), the convergence to equilibrium is comparable as measured in sweeps.



# Conclusions

- Both parallel tempering and population annealing overcome the exponential slowing associated with large free energy barriers.
- Population annealing is comparably efficient to parallel tempering and has several features to recommend it:
  - Massively parallel
  - Direct measurement of free energies
  - Perhaps more efficient?