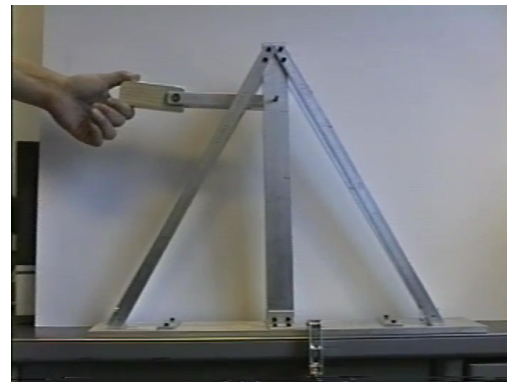


### So far: mostly about *maps*.

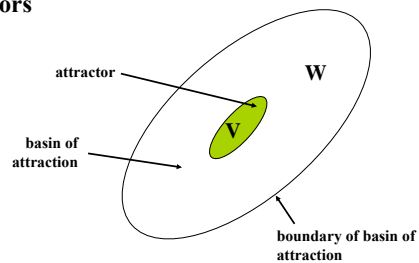
- discrete time systems:
  - time proceeds in clicks
  - “maps”
  - modeling tool: differ~~ence~~*ence* equation

### Next up: *flows*

- continuous time systems:
  - time proceeds smoothly
  - “flows”
  - modeling tool: differ~~ential~~*ential* equations



### Attractors



- Attractors exist only in dissipative systems!
- Dissipation  $\iff$  contraction of state space under the influence of the dynamics
- Can still have chaos if no dissipation...just not chaotic *attractors*



Courtesy of Allison Brown  
Best Poster prize, Experimental Chaos Conference, 2012

### Conditions for chaos in continuous-time systems

#### Necessary:

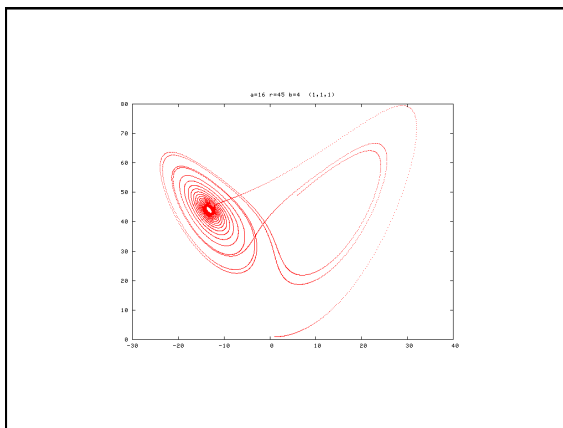
- Nonlinear
- At least three state-space dimensions (NB: only one needed in maps)

#### Necessary and sufficient:

- Cannot be solved in closed form ("nonintegrable," in Hamiltonian parlance)

### Concepts: review

- State variable
- State space
- Initial condition
- Trajectory
- Attractor
- Basin of attraction
- Transient
- Fixed point (un/stable)
- Bifurcation
- Parameter



### Deterministic Nonperiodic Flow<sup>1</sup>

EDWARD N. LORENZ

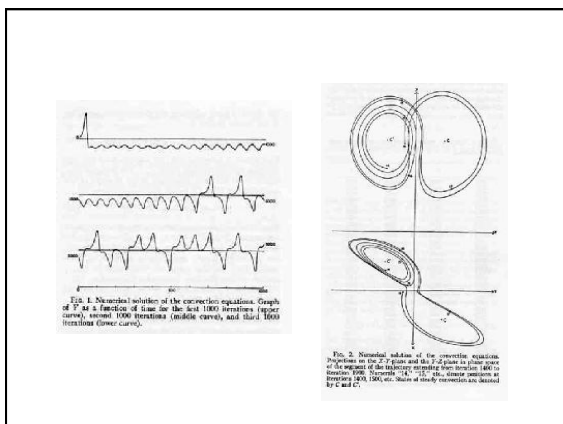
Massachusetts Institute of Technology

(Manuscript received 18 November 1962, in revised form 7 January 1963)

#### ABSTRACT

Finite systems of deterministic ordinary nonlinear differential equations may be designed to represent forced dissipative hydrodynamic flow. Solutions of these equations can be identified with trajectories in phase space. For those systems with bounded solutions, it is found that nonperiodic solutions are ordinarily unstable with respect to small modifications, so that slightly differing initial states can evolve into considerably different states. Systems with bounded solutions are shown to possess bounded numerical solutions. A simple system representing cellular convection is solved numerically. All of the solutions are found to be unstable, and almost all of them are nonperiodic.

The feasibility of very-long-range weather prediction is examined in the light of these results.

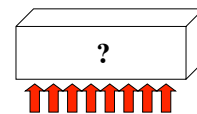


- Equations:

$$x' = a(y - x)$$

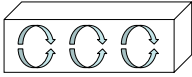
$$y' = rx - y - xz$$


$$z' = xy - bz$$

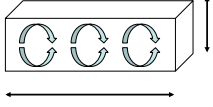
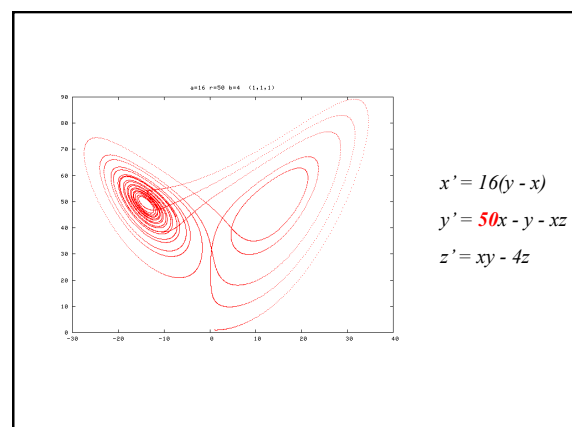
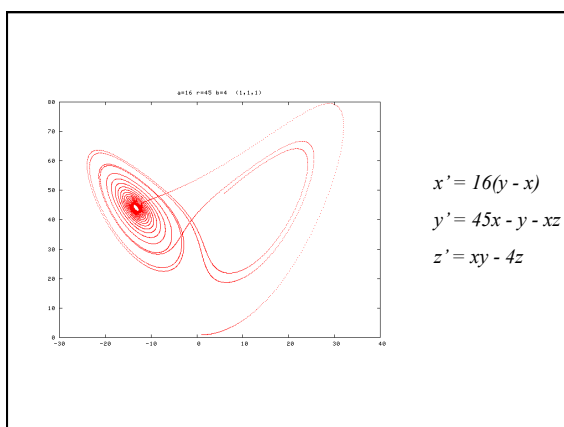
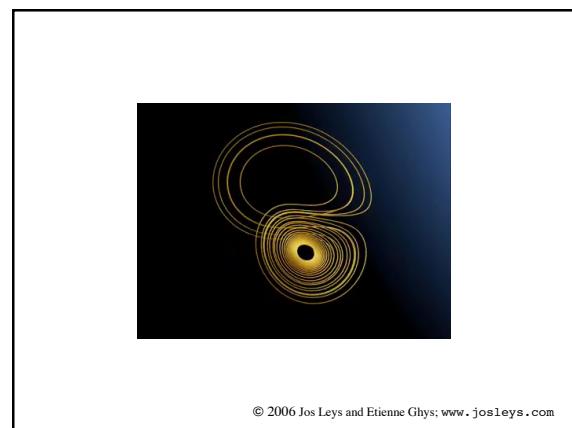
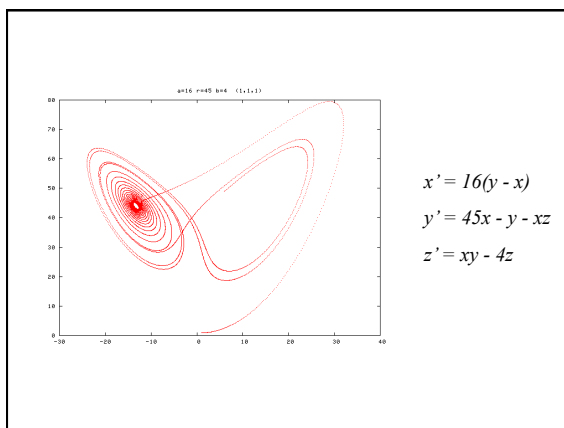


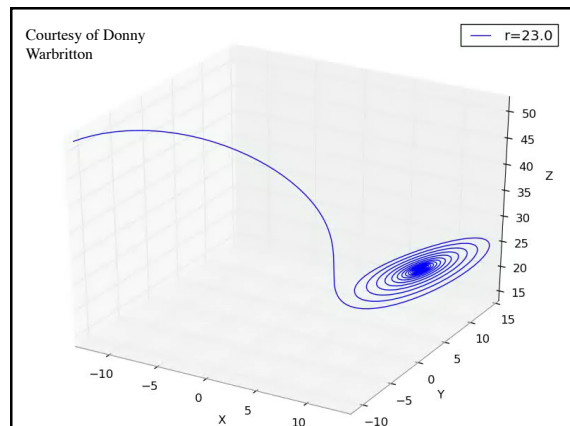
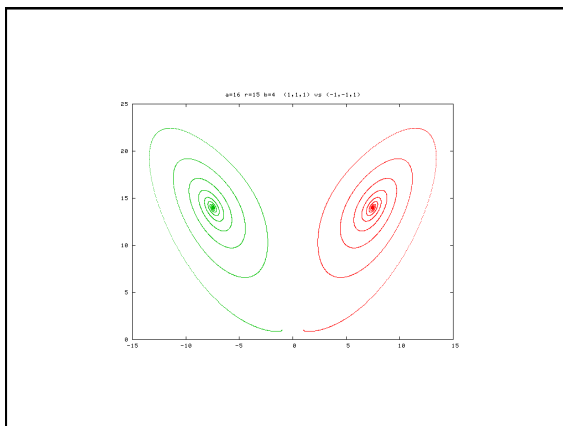
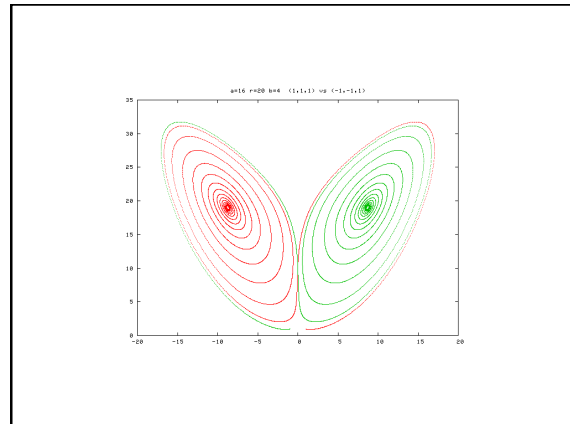
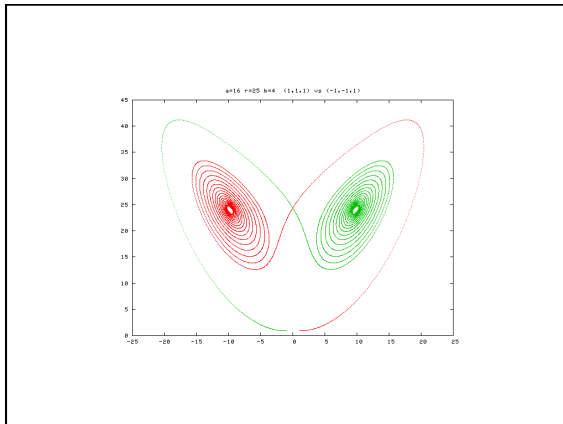
(first three terms of a Fourier expansion of the Navier-Stokes eqns)

- State variables:
  - $x$  convective intensity
  - $y$  temperature
  - $z$  deviation from linearity in the vertical convection profile



- Parameters:
  - $a$  Prandtl number - fluids property
  - $r$  Rayleigh number - related to  $\Delta T$  
  - $b$  aspect ratio of the fluid sheet



## Before we leave Lorenz...

### Deterministic Nonperiodic Flow<sup>1</sup>

EDWARD N. LORENZ

Massachusetts Institute of Technology

(Manuscript received 18 November 1962, in revised form 7 January 1963)

#### ABSTRACT

Finite systems of deterministic ordinary nonlinear differential equations may be designed to represent forced dissipative hydrodynamic flow. Solutions of these equations can be identified with trajectories in phase space. For those systems with bounded solutions, it is found that nonperiodic solutions are ordinarily unstable with respect to small modifications, so that slightly differing initial states can evolve into considerably different states. Systems with bounded solutions are shown to possess bounded numerical solutions. A simple system representing cellular convection is solved numerically. All of the solutions are found to be unstable, and almost all of them are nonperiodic.

*The feasibility of very-long-range weather prediction is examined in the light of these results.*

## Attractors

### Four types:

- fixed points
- limit cycles (*aka* periodic orbits)
- quasiperiodic orbits
- chaotic attractors

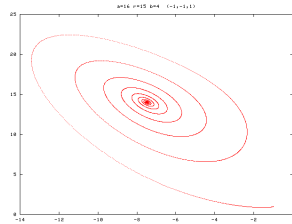
A nonlinear system can have any number of attractors, of all types, sprinkled around its state space

Their basins of attraction (plus the basin boundaries) *partition* the state space

And there's no way, *a priori*, to know where they are, how many there are, what types, etc.

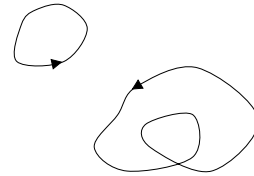
### Attractors

- Fixed point



### Attractors

- Limit cycle

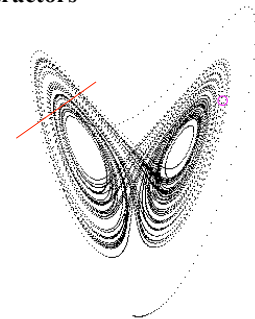


### Attractors

- Quasi-periodic orbit...

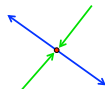
### “Strange” or chaotic attractors

- often fractal
- covered densely by trajectories
- exponential divergence of neighboring trajectories...



### Lyapunov exponents

- nonlinear analogs of eigenvalues: one  $\lambda$  for each dimension

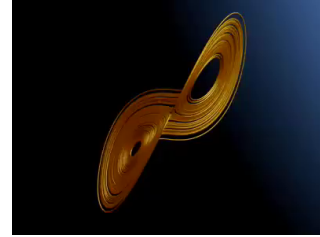
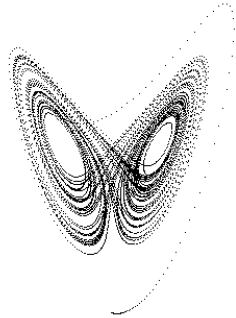


### Lyapunov exponents: summary

- nonlinear analogs of eigenvalues: one  $\lambda$  for each dimension
- negative  $\lambda_i$  compress state space; positive  $\lambda_i$  stretch it
- $\sum \lambda_i < 0$  for dissipative systems
- $\lambda_i$  are same for all ICs in one basin
- long-term average in definition; biggest one ( $\lambda_1$ ) dominates as  $t \rightarrow \infty$
- positive  $\lambda_1$  is a signature of chaos

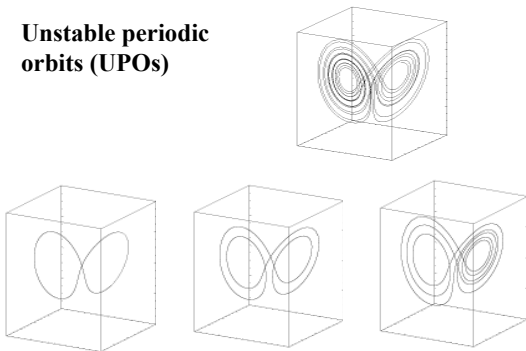
### “Strange” or chaotic attractors:

- exponential divergence of neighboring trajectories
- *often* fractal
- covered densely by trajectories
- contain an infinite number of “unstable periodic orbits”...



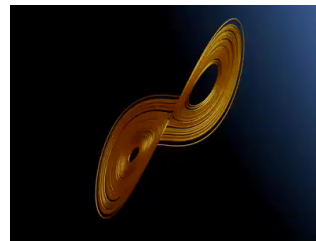
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### Unstable periodic orbits (UPOs)

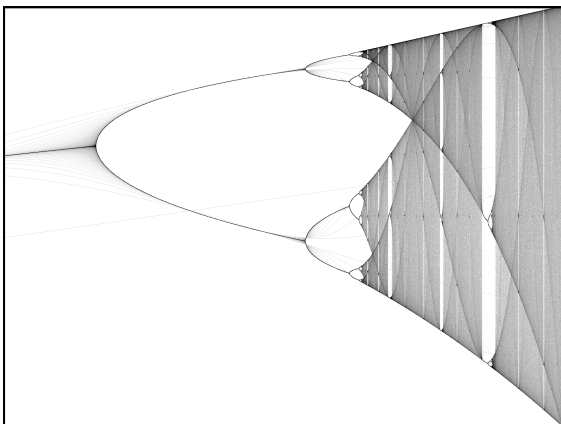


Bradley/Mantilla, *Chaos* 12:596

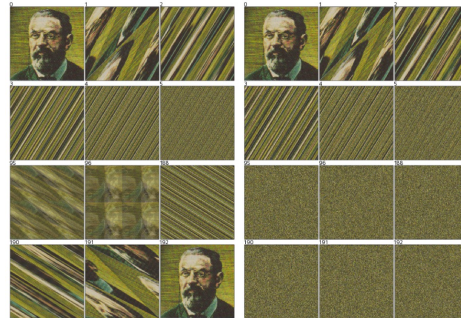
### “Attractor bones”...



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### Poincare recurrence



Representative point chosen at corner of pixel

Representative point chosen at a random place in pixel

Crutchfield *et al.* *Chaos* 25:46 and  
<http://www.mpi-pks-dresden.mpg.de/mpi-doc/kantgruppe/wiki/projects/Recurrence.html>