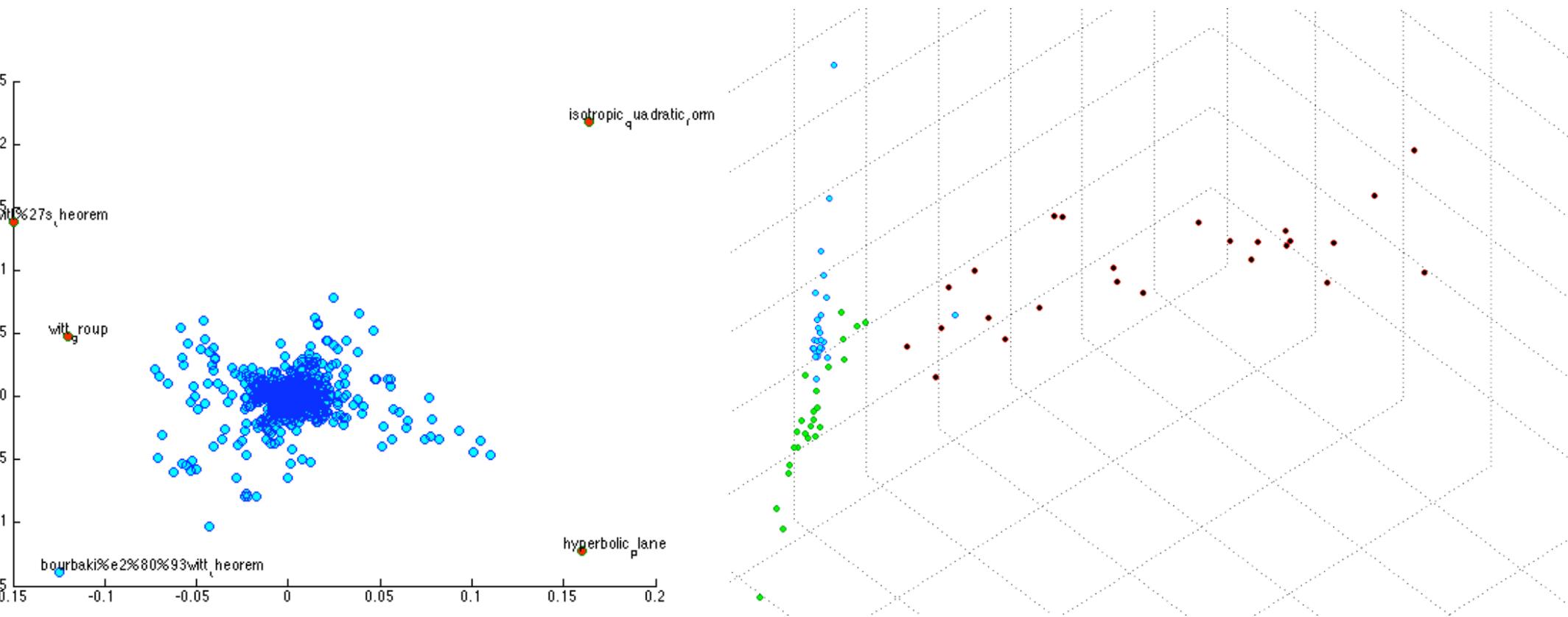


Fraud Detection, Quantum Mechanics, and Complex Systems

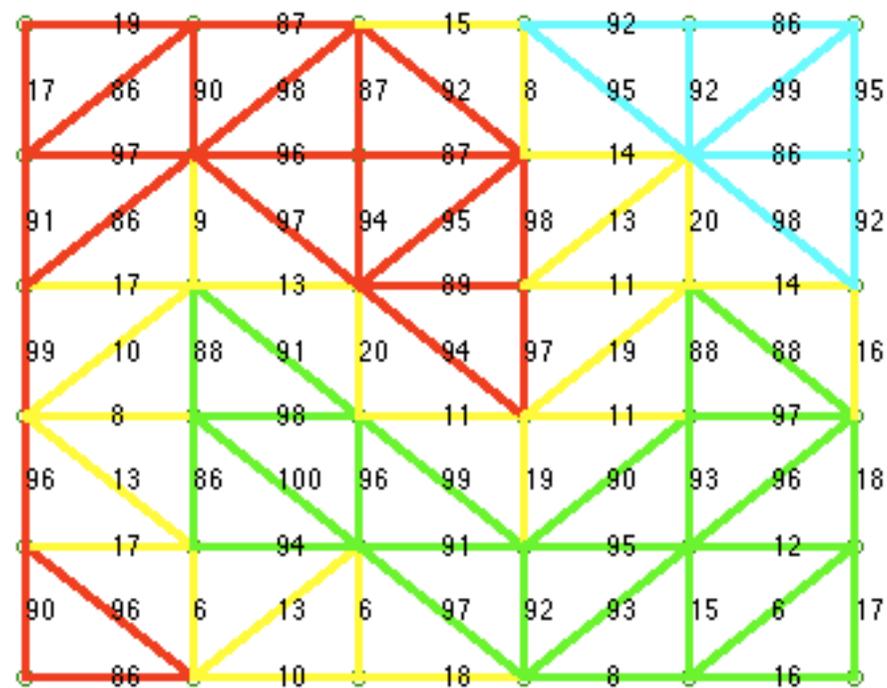
Lecture 3, CSSS10



Greg Leibon
Memento, Inc
Dartmouth College

The Vagabond Problem: Finding a partition $\{R_i\}_{i=1}^K$ of our states into regions that minimizes

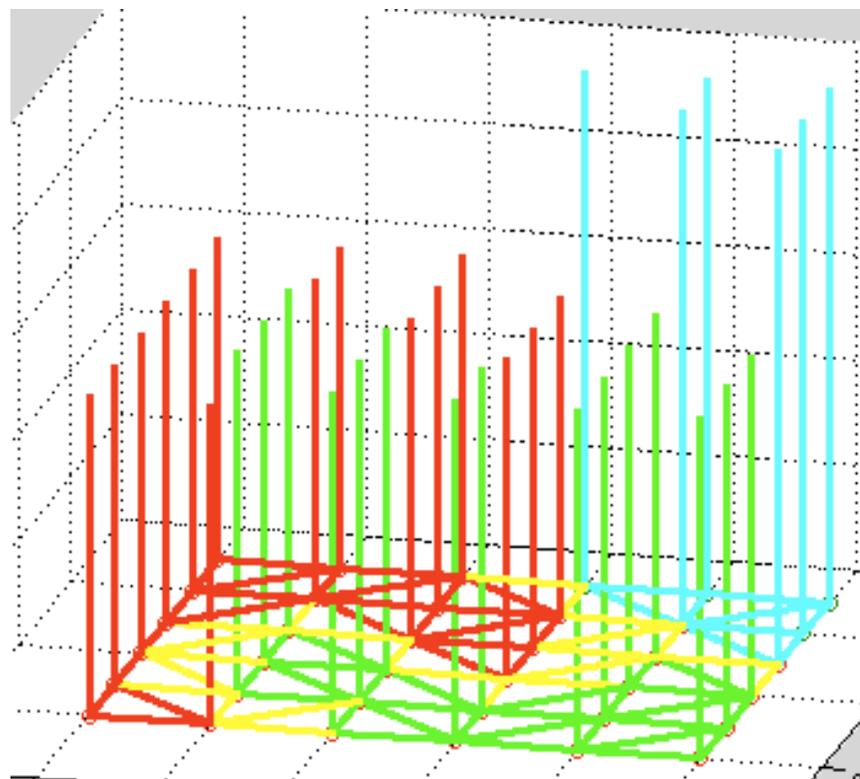
$$\sum_i P(X_{t+dt} \in \bar{R}_i \mid X_t \in R_i) :$$



$N=200$; $K=10$; `[states]=PlotTheState(N,K);`

$$\chi_R(i) = \begin{cases} \frac{1}{\sqrt{\omega(R)}} & \text{if } i \in R \\ 0 & \text{if } i \in \bar{R} \end{cases}$$

The
solution
(I think!)



NP hard

if R_1 and R_2 are disjoint $\langle \chi_{R_1}, \chi_{R_2} \rangle_{\omega} = \delta_1^2$

Recall: $\langle f, g \rangle_{\omega} = f^{tr}[\omega]g$

How we really solve this?

Vagabond Theorem:

$$P(X_{t+dt} \in \bar{R} \mid X_t \in R) \approx \|\chi_R\|_{Dir}^2 dt$$

proof:

$$\begin{aligned}
 P(X_{t+dt} \in \bar{R} \mid X_t \in R) &= \frac{P(X_{t+dt} \in \bar{R}, X_t \in R)}{P(X_t \in R)} && \text{(def. conditional prob.)} \\
 &= \frac{\sum_{i \in \bar{R}, j \in R} P(X_{t+dt} = i, X_t = j)}{\omega(R)} && \text{(def.)} \\
 &= \frac{\sum_{i \in \bar{R}, j \in R} P(X_{t+dt} = i \mid X_t = j) P(X_t = i)}{\omega(R)} && \text{(def. conditional prob.)} \\
 &= \frac{\sum_{i \in \bar{R}, j \in R} (1/\tau_j) p_{ij} dt \omega_i}{\omega(R)} && \text{(instantaneous transition rate)} \\
 &= \frac{\sum_{i \in \bar{R}, j \in S} p_{ij} \pi_i}{\omega(R)} dt && (\omega_i = \tau_j \pi_i) \\
 &= \sum_{i \in \bar{R}, j \in R} \pi_i p_{ij} (\chi_R(i) - \chi_R(j))^2 dt && \text{def.} \\
 &= (1/2) \sum_{i, j} \pi_i p_{ij} (\chi_R(i) - \chi_R(j))^2 dt && \text{(summand 0 if } i, j \in S \text{ or } i, j \in \bar{R}) \\
 &= \langle \chi_R, \Delta \pi \chi_R \rangle_\pi dt && \text{Green's Identity} \\
 &= \|\chi_R\|_{Dir}^2 dt && \text{conformal invariance}
 \end{aligned}$$

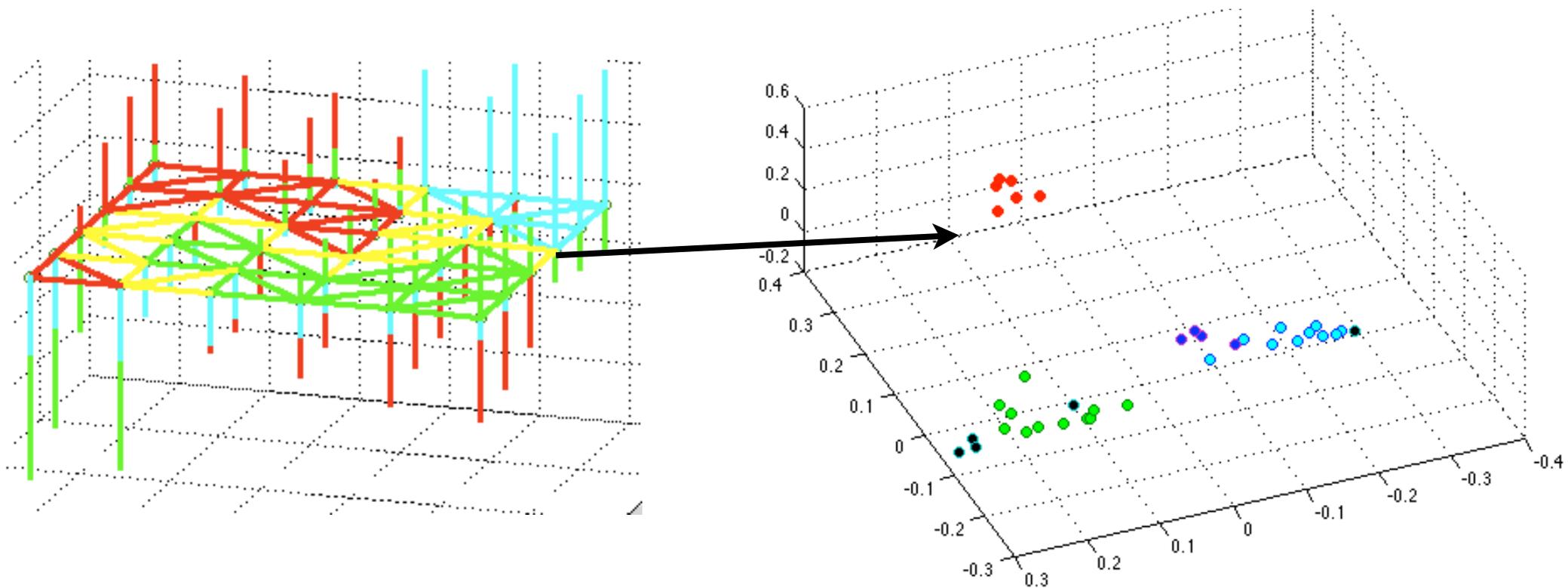
Q.E.D

$$\sum_i P(X_{t+dt} \in \bar{R}_i | X_t \in R_i) = \sum_i \|\chi_{R_i}\|_{Dir}^2.$$

Now we can relax the Vagabond Problem and search for a ω -orthonormal set $\{f_i\}_{i=1}^L$ such that

$$\sum_i \|f_i\|_{Dir}^2$$

is minimized.



Now we cluster in Euclidean space...

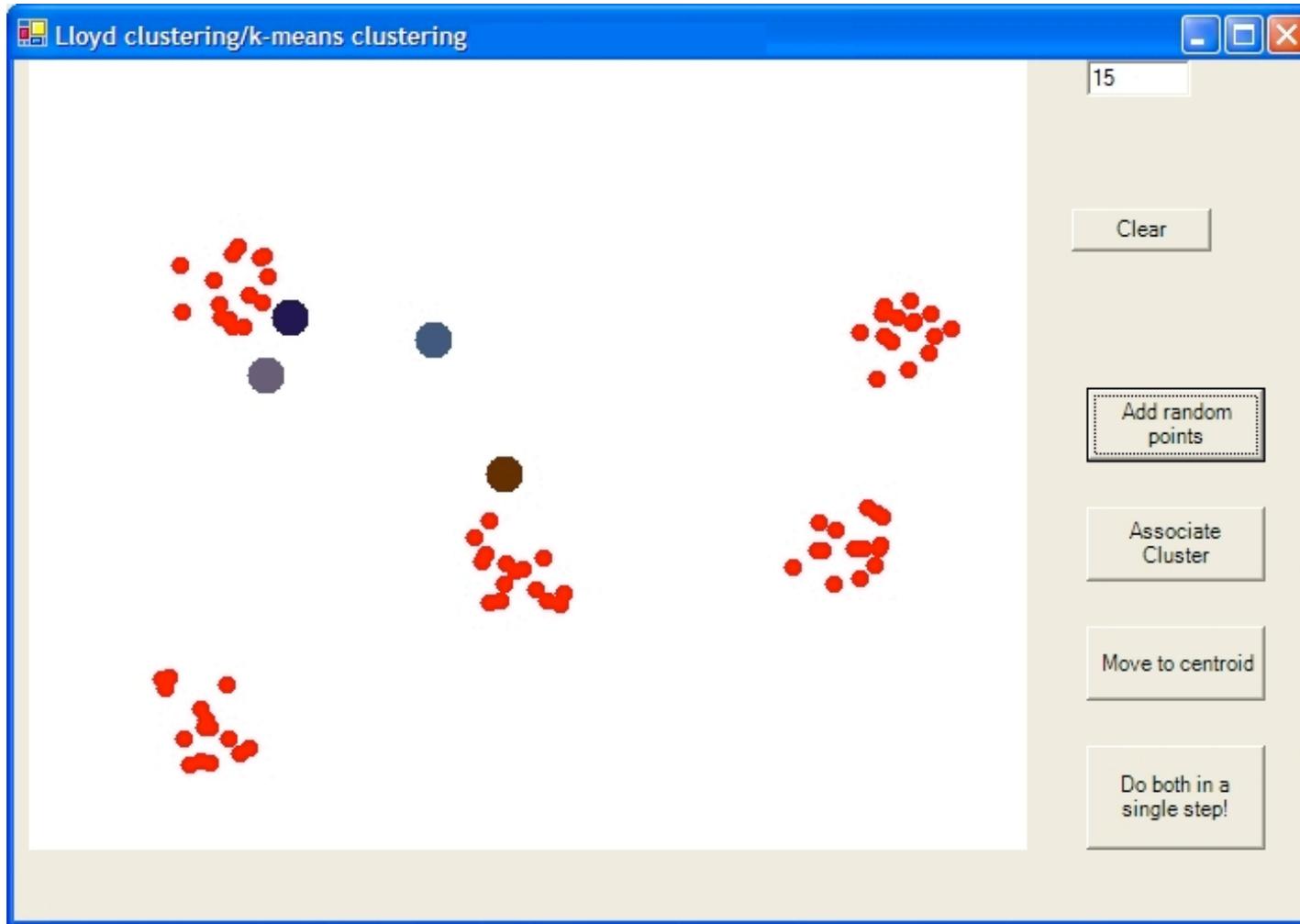
K-means

- Simplest clustering algorithm is **k-means**
- To run requires fixing $K = \#(\text{Clusters})$
- Requires an Euclidean type embedding
- We are attempting to minimizing a loss function:

$$L = \sum_{k=1}^K \sum_{x_i \in C_k} (x_i - \mu_k)^2$$

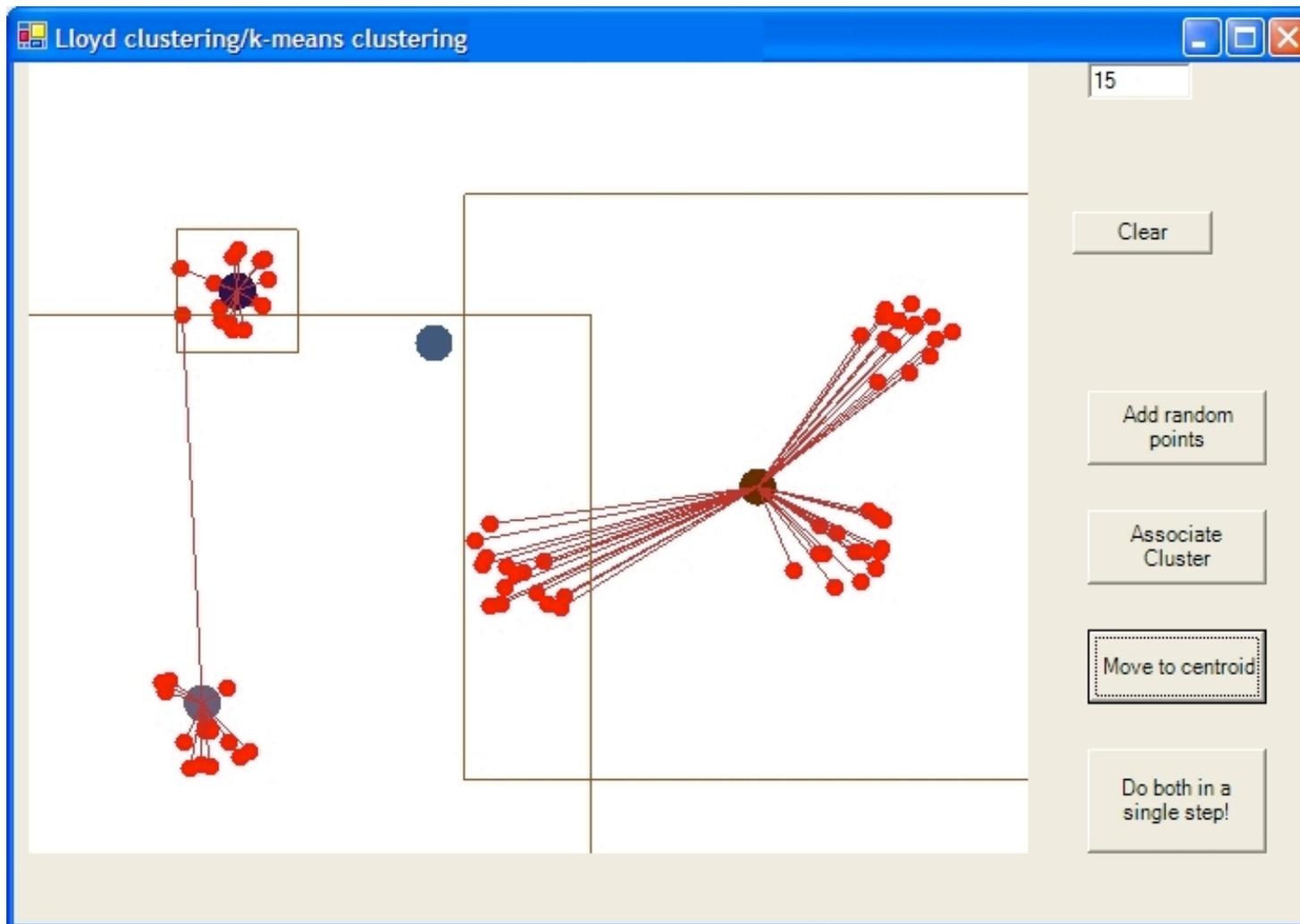
K-means algorithm:

I. Randomly choose points in each cluster and compute centroids.

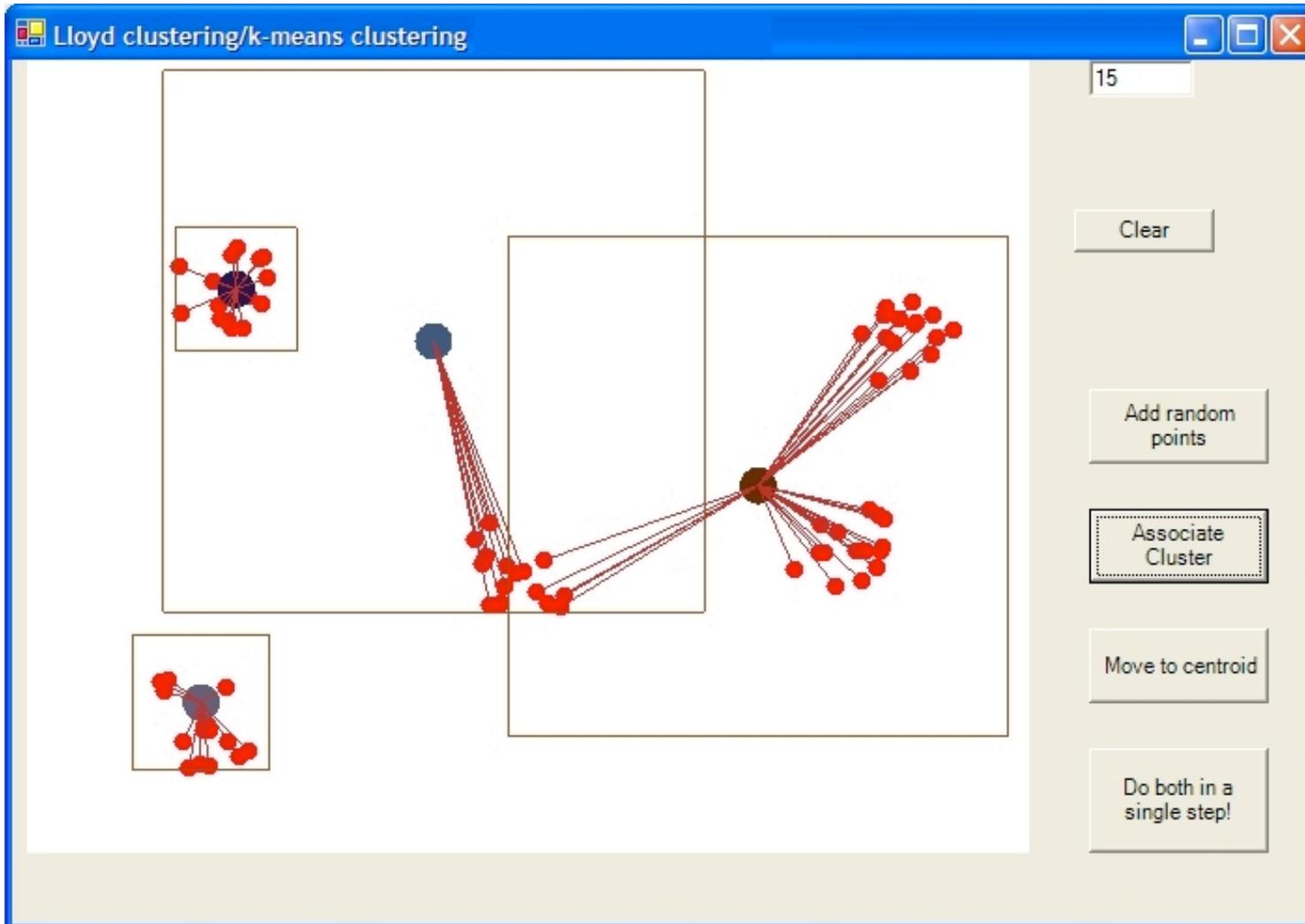


Example from: http://en.wikipedia.org/wiki/K-means_algorithm

2. Organize points by distance to the centroids.
3. Update centroids



4. Repeat...

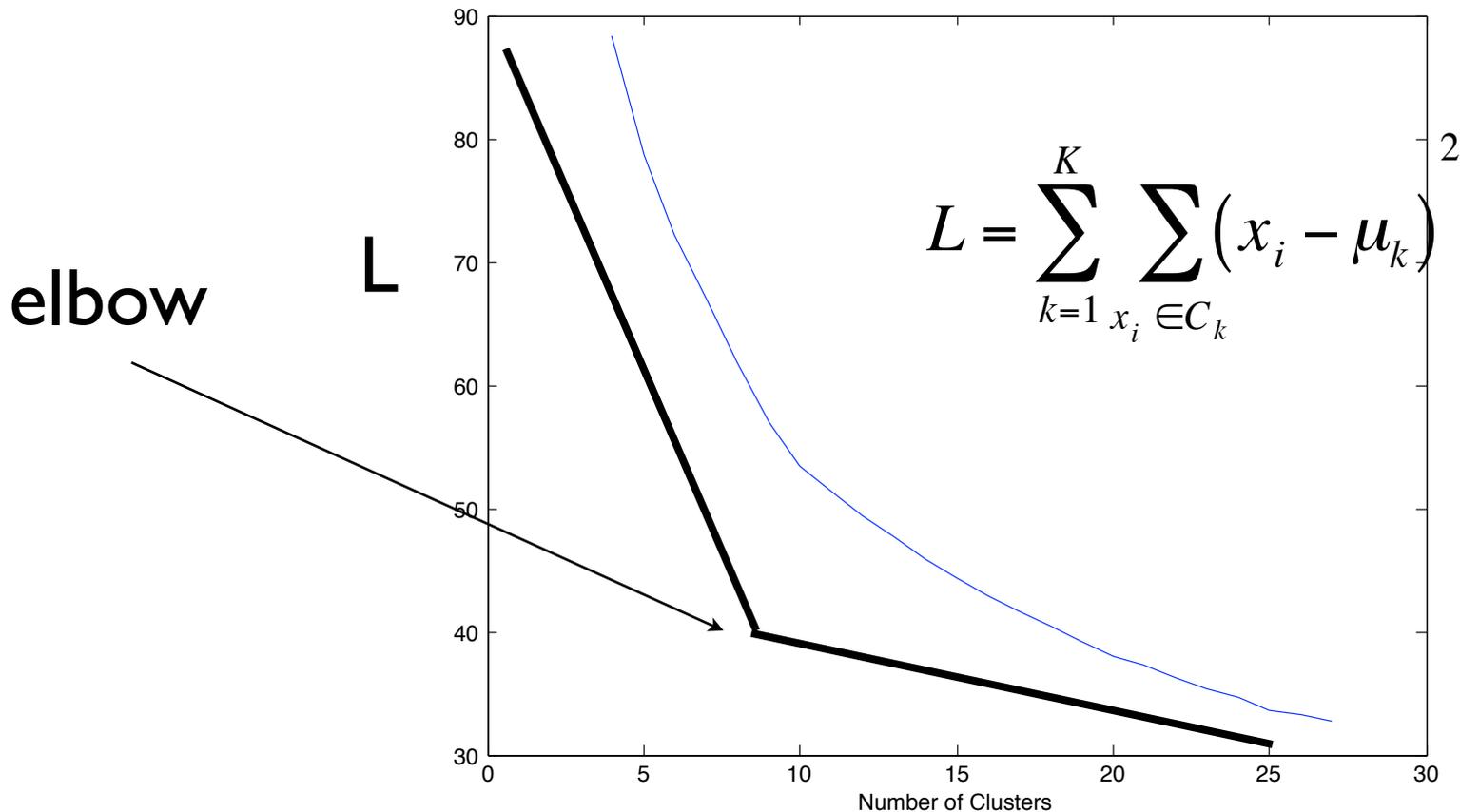


...until stable.

The screenshot shows a software window titled "Lloyd clustering/k-means clustering". The window contains a main canvas with three small diagrams illustrating the clustering process: a cluster of red points with a blue centroid, a cluster of red points with a blue centroid, and a cluster of red points with a blue centroid. To the right of the canvas is a control panel with a text input field containing the number "15". Below the input field are several buttons: "Clear", "Add random points", "Associate Cluster", "Move to centroid" (which is highlighted with a dotted border), and "Do both in a single step!".

A hard part is choosing $K = \#(\text{Clusters})$

Elbowology:



Though in practice this rarely works in a complex multi-scalar system

Now we can relax the Vagabond Problem and search for a ω -orthonormal set $\{f_i\}_{i=1}^L$ such that

$$\sum_i^L \|f_i\|_{Dir}^2$$

is minimized.

How?

$\langle f_i, \Delta_\omega f_i \rangle_\omega = \|f_i\|_{Dir}^2$ is being used as a quadratic form, hence each term in the sum remains the same if we view Δ_ω as the Hermitian operator

$$C_\omega = \frac{\Delta_\omega + \Delta_\omega^*}{2}.$$

In other words,

$$\langle f, \Delta_\omega f \rangle_\omega = \langle f, C_\omega f \rangle_\omega$$

and we have replaced the minimization problem with minimizing

$$\sum_i \frac{\langle f_i, C_\omega f_i \rangle_\omega}{\langle f_i, f_i \rangle_\omega}.$$

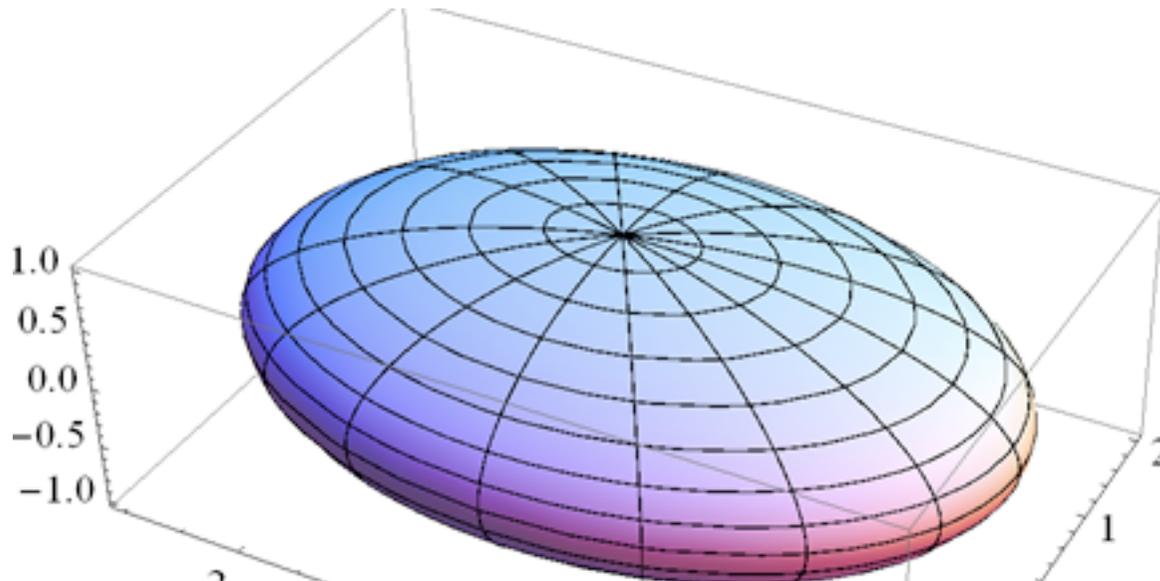
and we have replaced the minimization problem with minimizing

$$\sum_i \frac{\langle f_i, C_\omega f_i \rangle_\omega}{\langle f_i, f_i \rangle_\omega} \quad \text{where} \quad \langle f, C_\omega g \rangle = \langle C_\omega f, g \rangle$$

Why is this good, well it takes two to Tango!

C_ω is our *context operator* and being Hermitian allows us to use the spectral theorem and the Raileigh-Rtitz method to identify our Vagabond Embedding as the ω -unit length eigenfunctions of C_ω with the smallest eigenvalues.

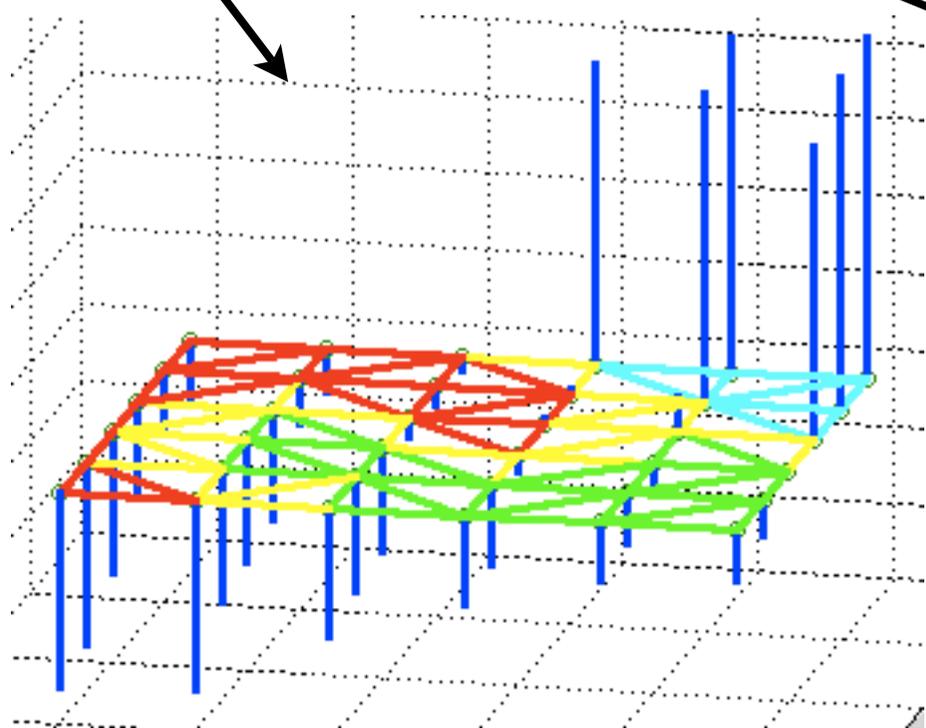
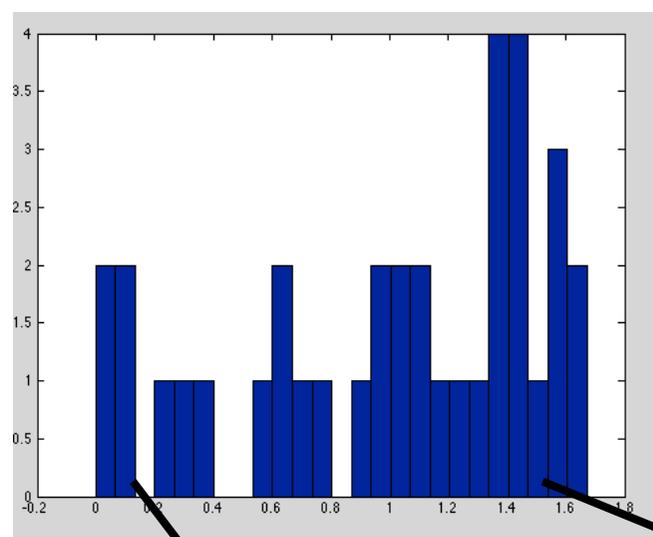
Spectral Theorem: If \langle, \rangle is a Hermitian inner product and $\langle Av, w \rangle = \langle v, Aw \rangle$, then there is an orthonormal basis of A eigenvectors.



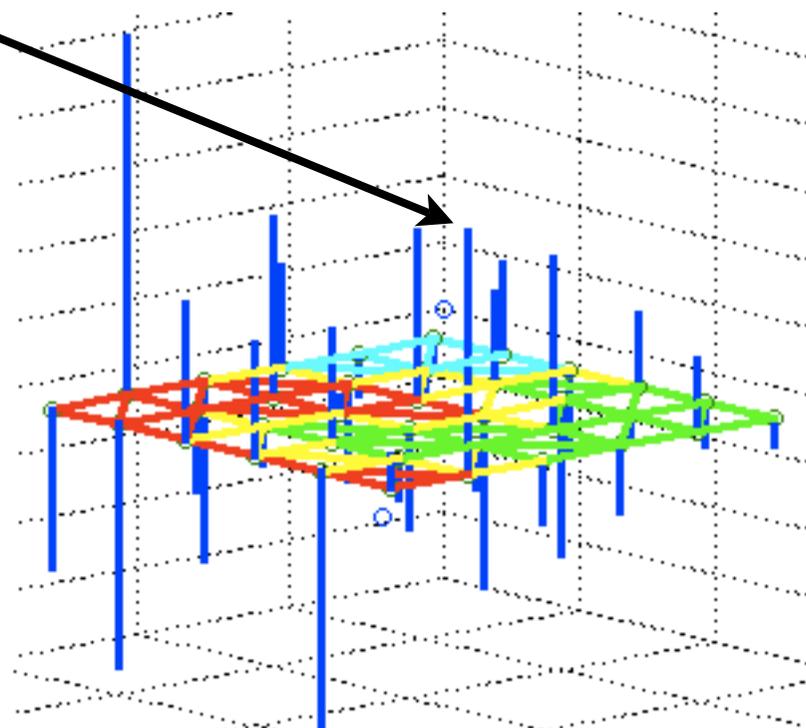
The Context Operator

$$C_{\omega} = \frac{\Delta_{\omega} + \Delta_{\omega}^*}{2}$$

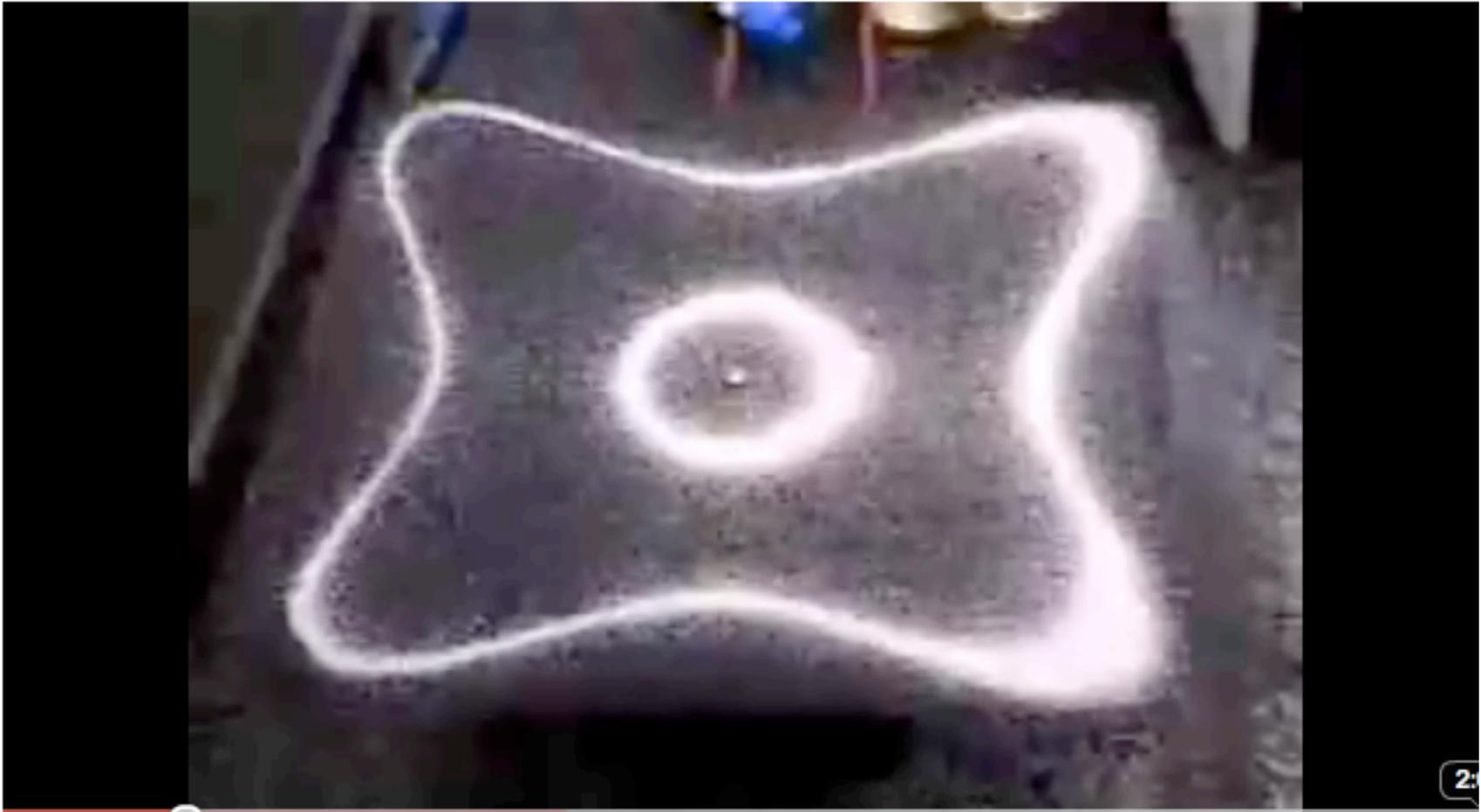
Information in Eigenfunctions
Not Local



λ_2



λ_{30}



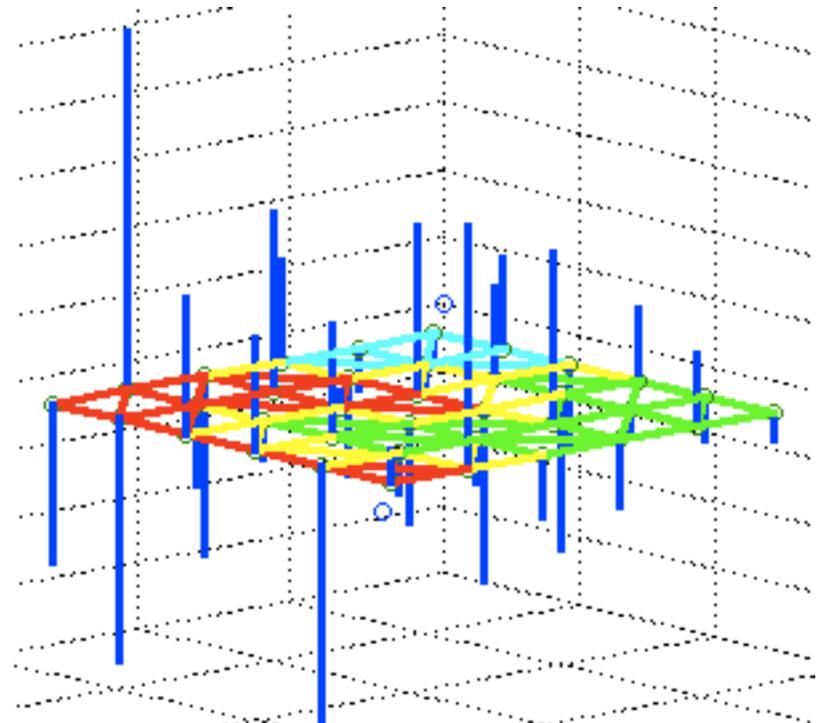
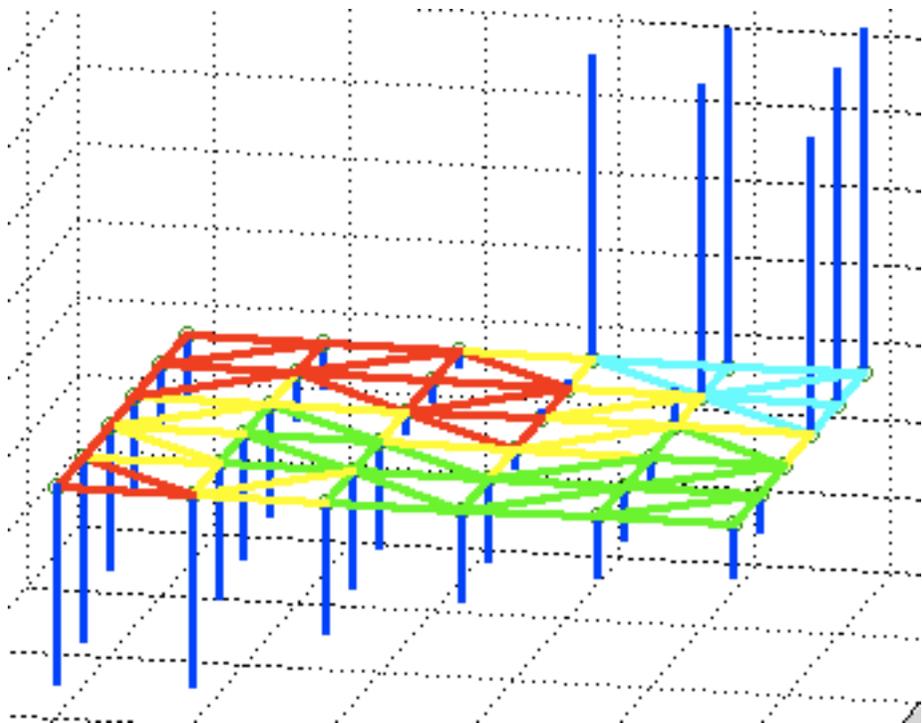
http://www.youtube.com/watch?v=Uu6Ox5LrhJg&feature=response_watch

The Green-Kelvin Identity

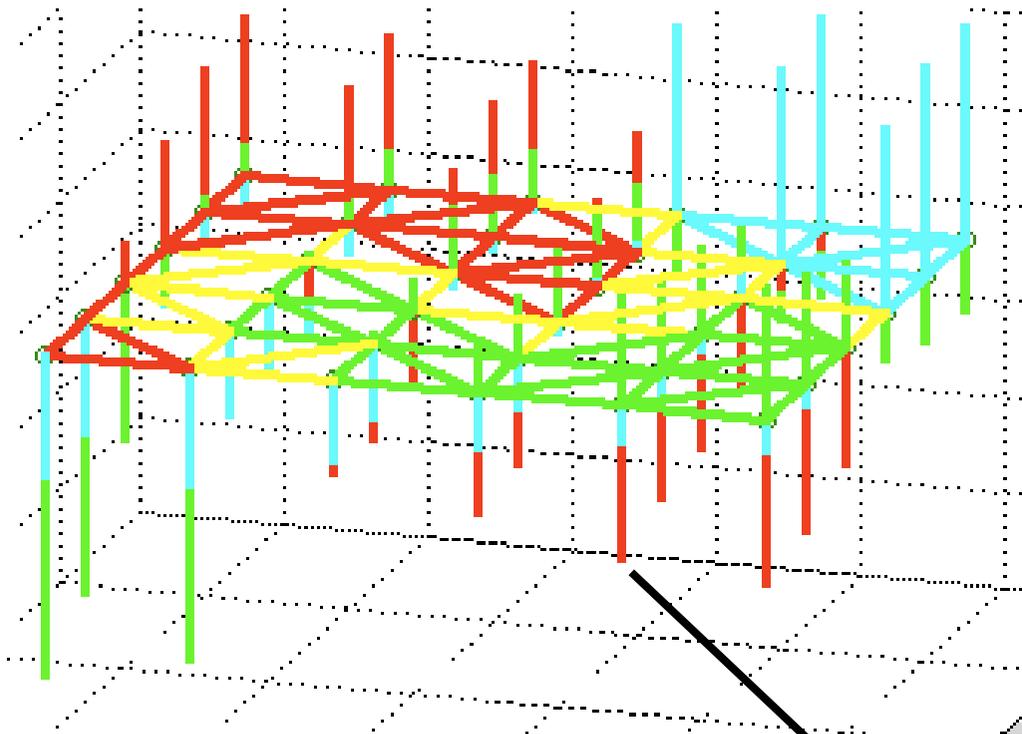
$$\langle f, \Delta f \rangle = \int f \Delta f d\vec{x} = \int |\nabla f|^2 d\vec{x} \quad \langle f, g \rangle = \int_M f(x)g(x)dVol(x)$$

$$\langle f, g \rangle = \sum_i f_i w^i g_i$$

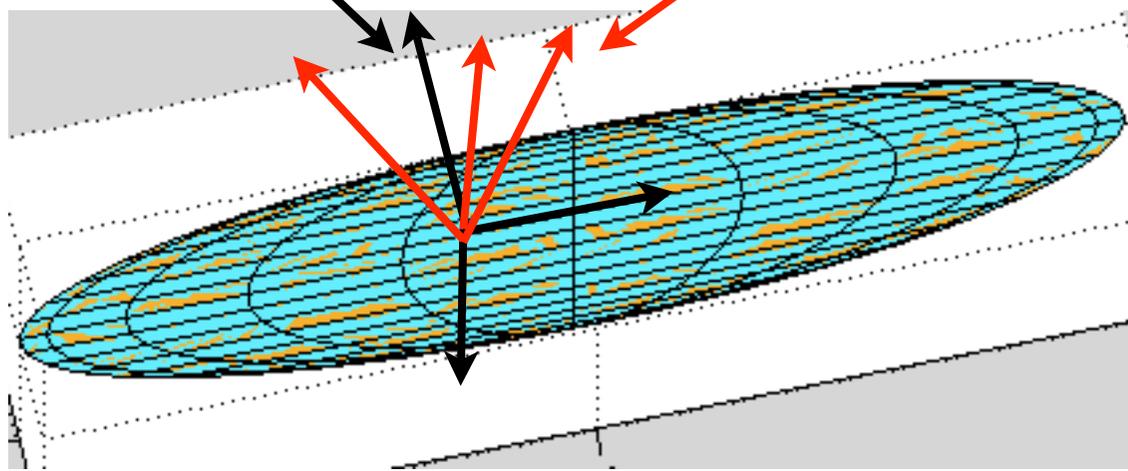
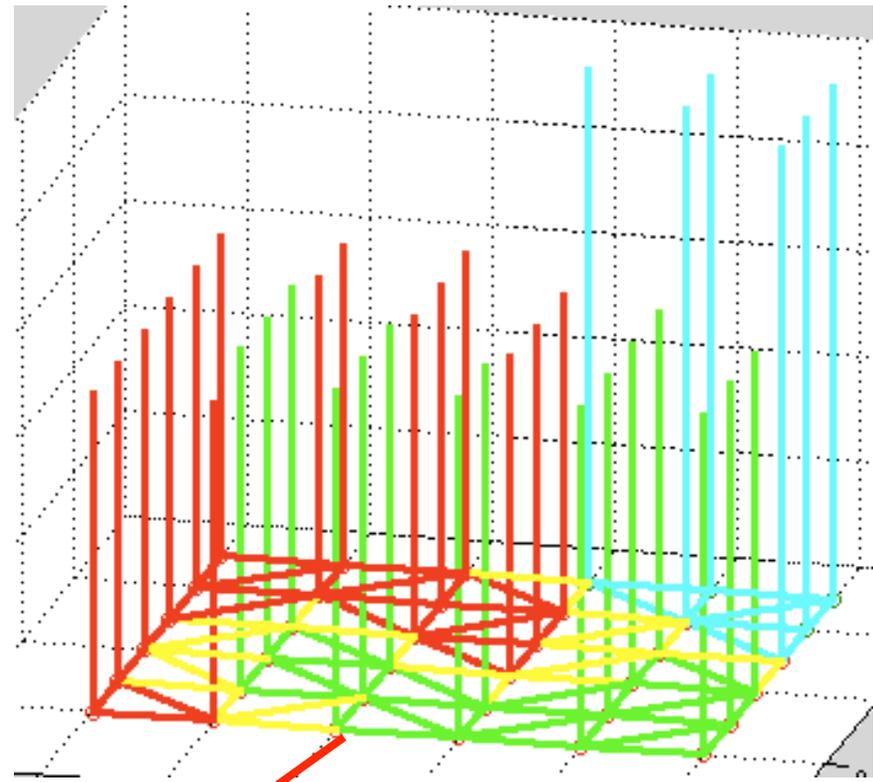
$$\langle f, \Delta f \rangle = \sum_{i,j} f_j w^i (I_i^j - P_i^j) f_j = \frac{1}{2} \sum_{i,j} w^i P_i^j (f_i - f_j)^2$$



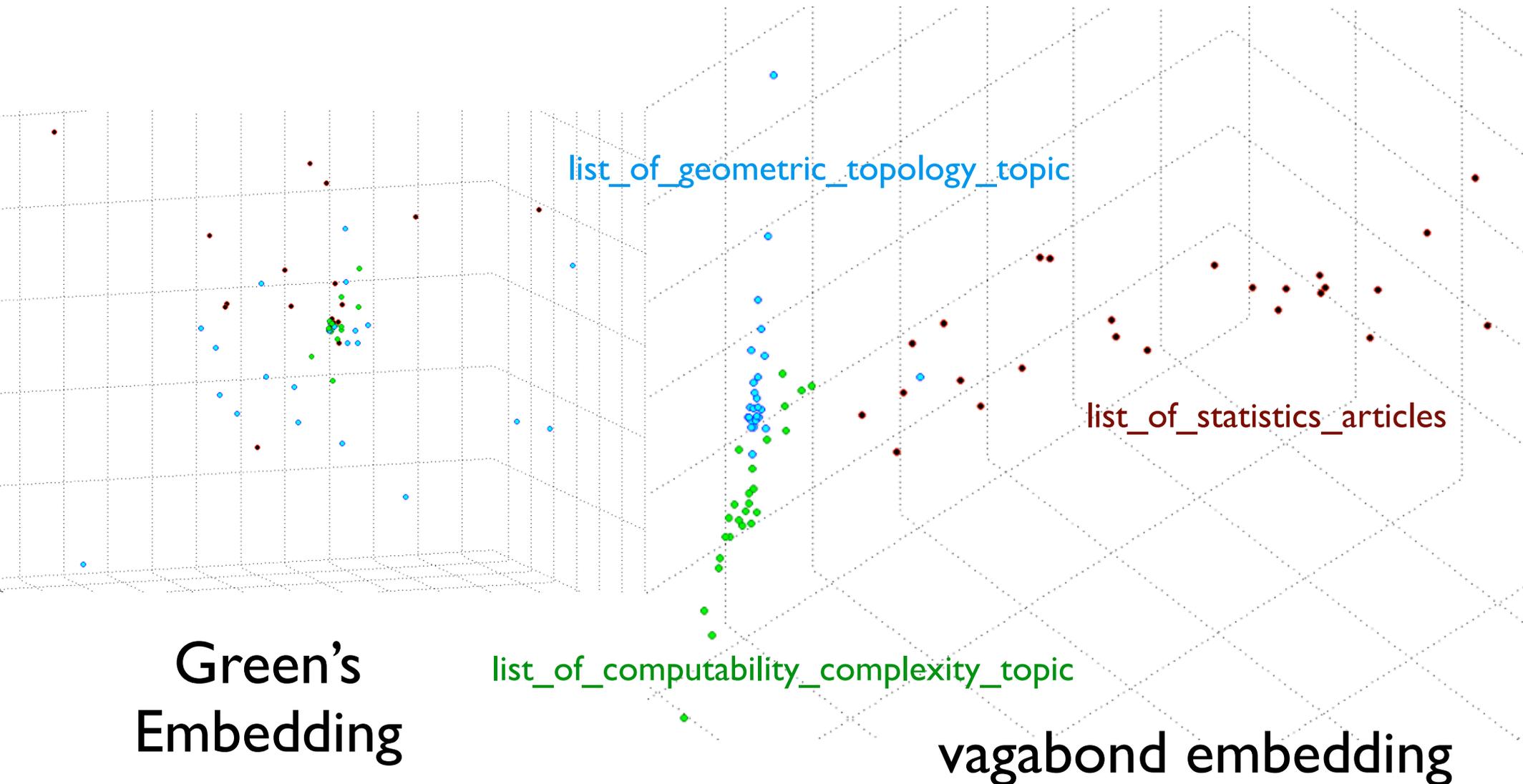
First three (non-trivial) eigenfunctions

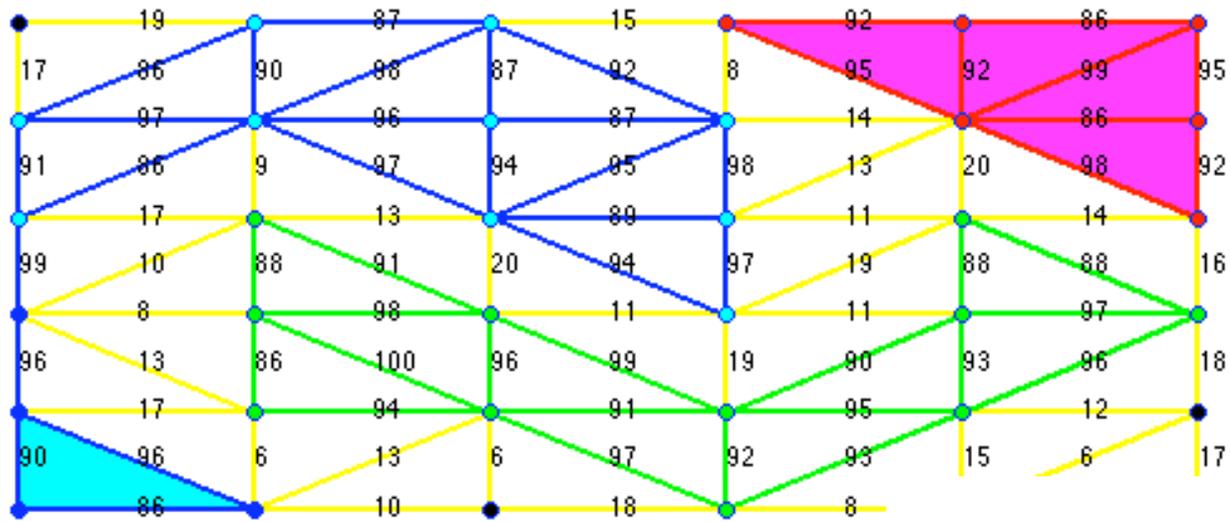


The perfect Vagabond functions



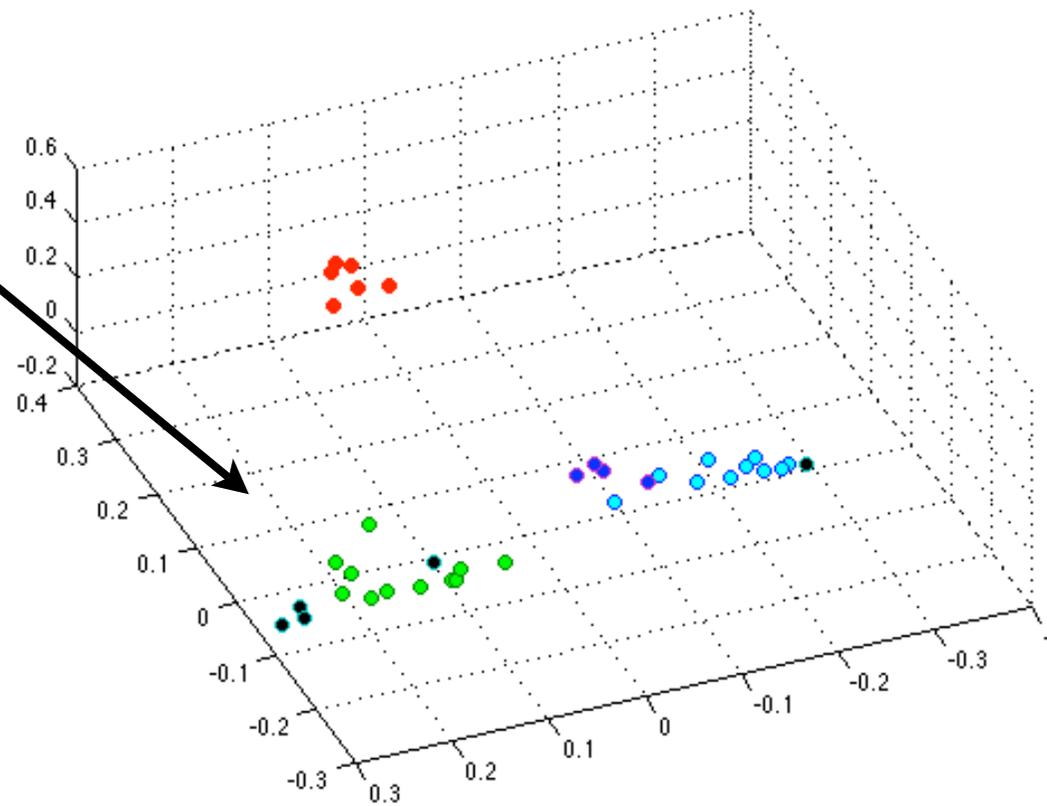
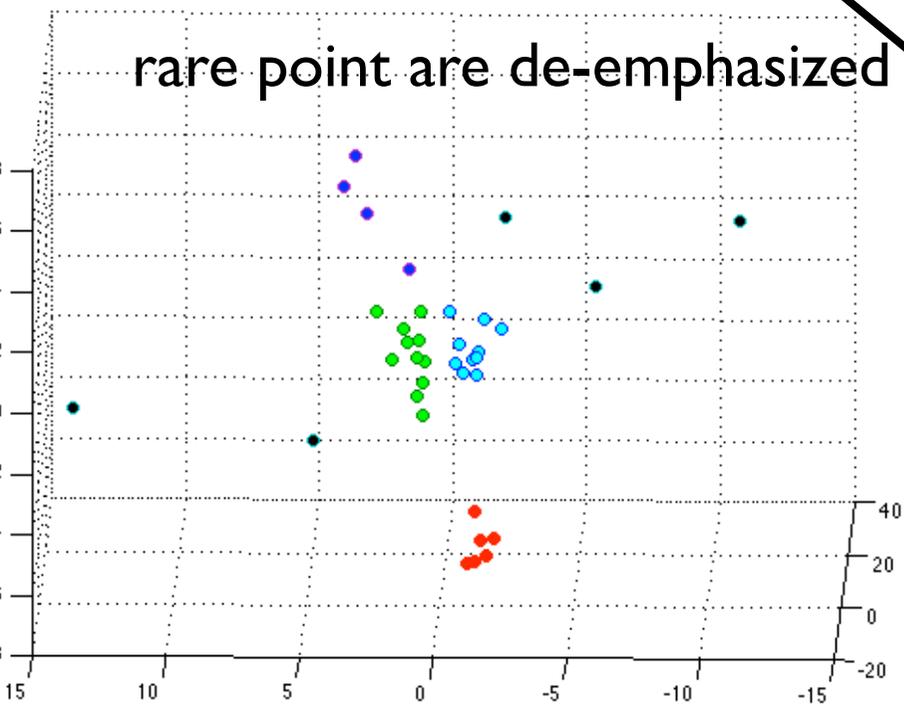
Space of mathematics



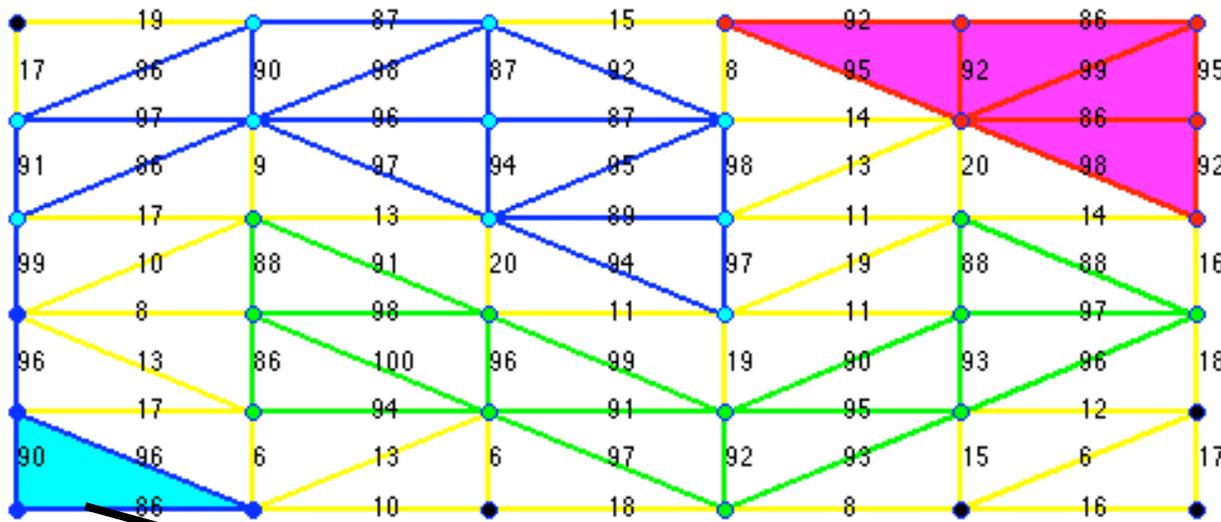


$$\tau_i = 1$$

rare point are de-emphasized

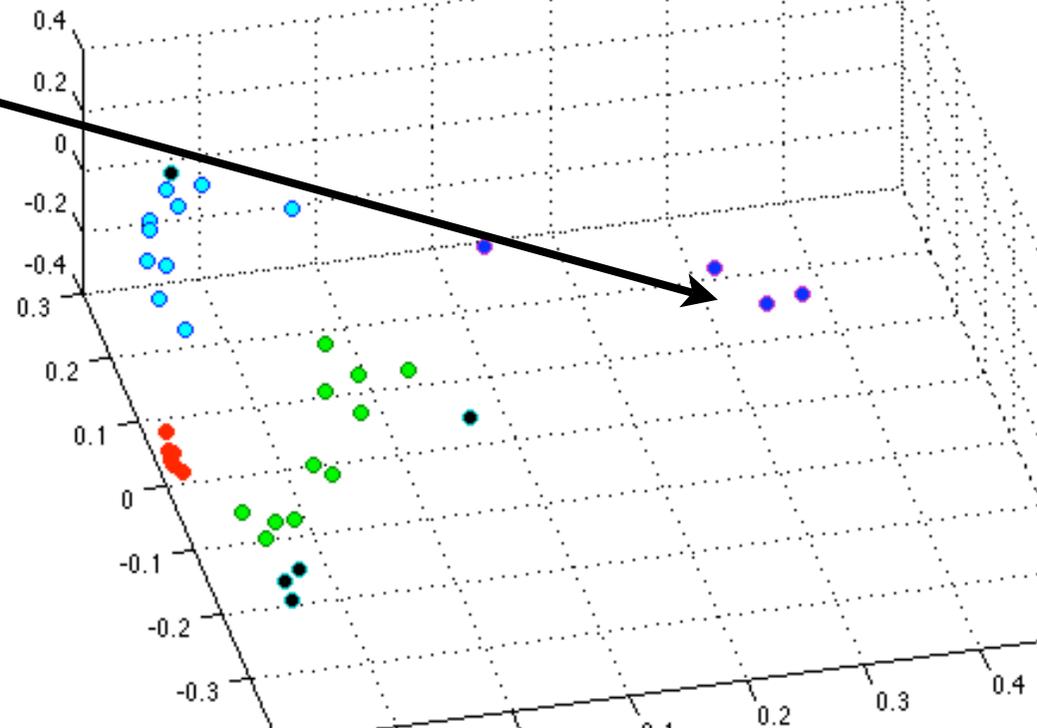
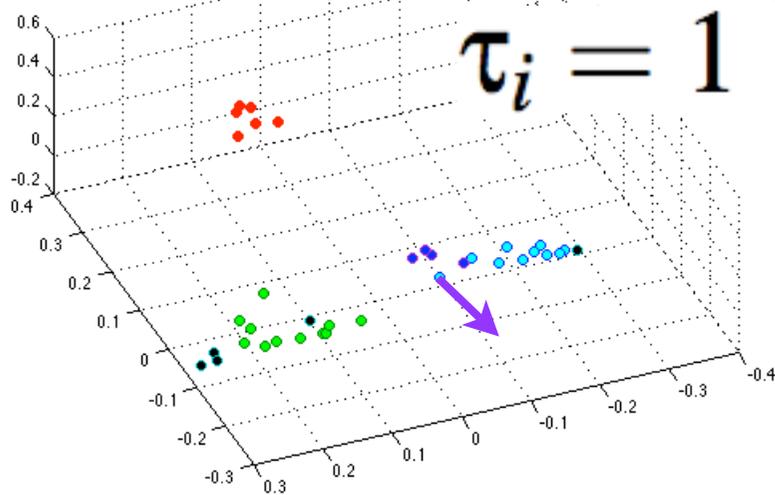


$\tau_0(\text{blue}) = 5, \tau_0(\text{magenta}) = 1/5, \text{ and } \tau_0(\text{other}) = 1$

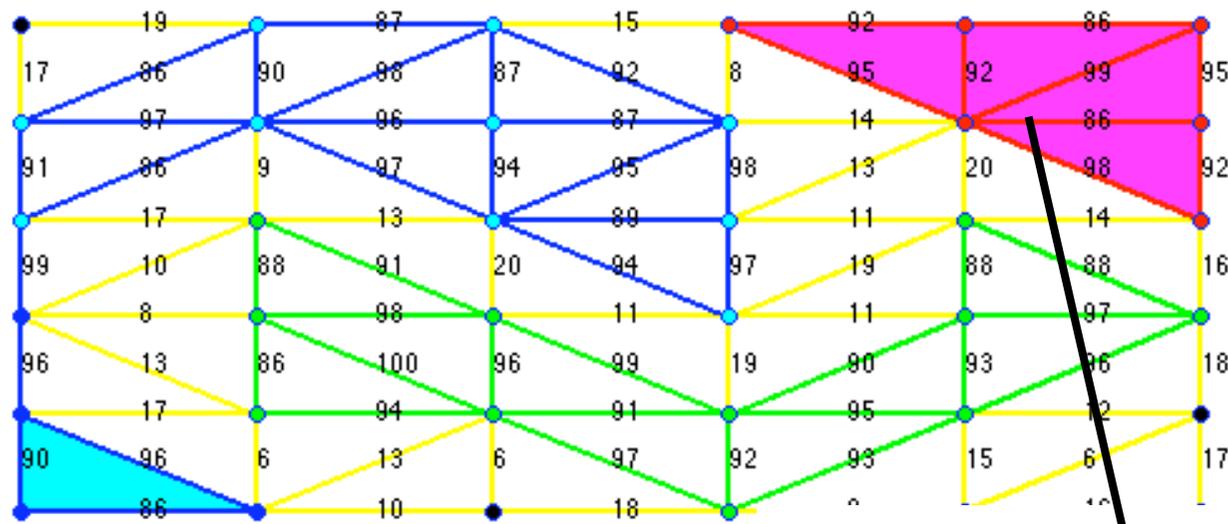


$$\tau = \frac{\tau_0}{\sum \tau_0 \pi}$$

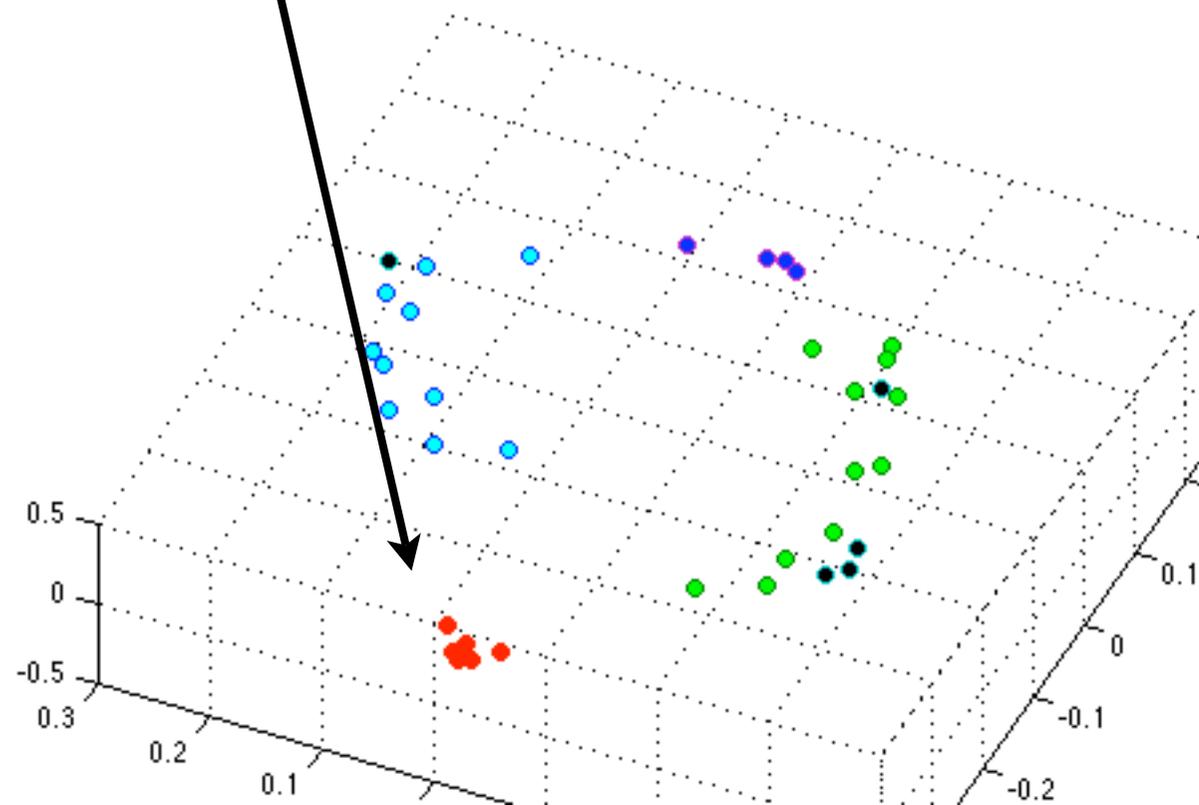
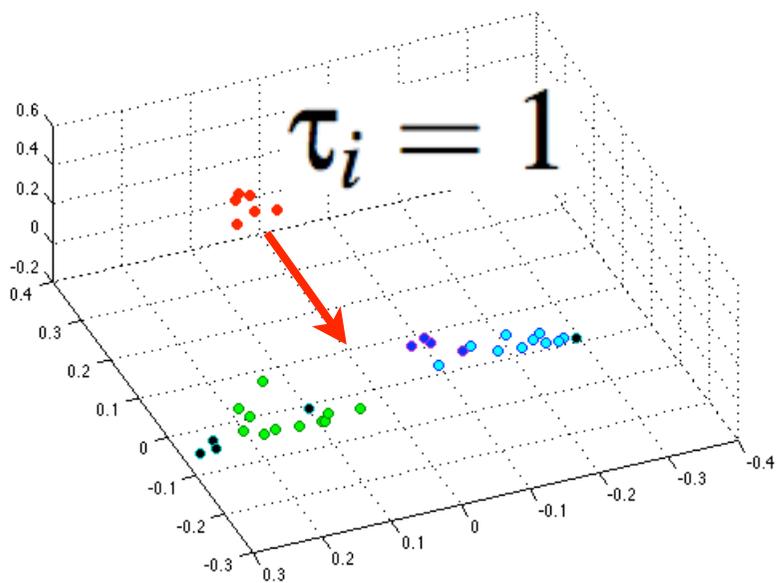
tries to become a cluster



$\tau_0(\text{blue}) = 1/5, \tau_0(\text{magenta}) = 1/5, \text{ and } \tau_0(\text{other}) = 1$



$$\tau = \frac{\tau_0}{\sum \tau_0 \pi}$$



Kakawa's Salted Carmel time! (well, except Kakawa wasn't open)

What is the difference?

figure 1

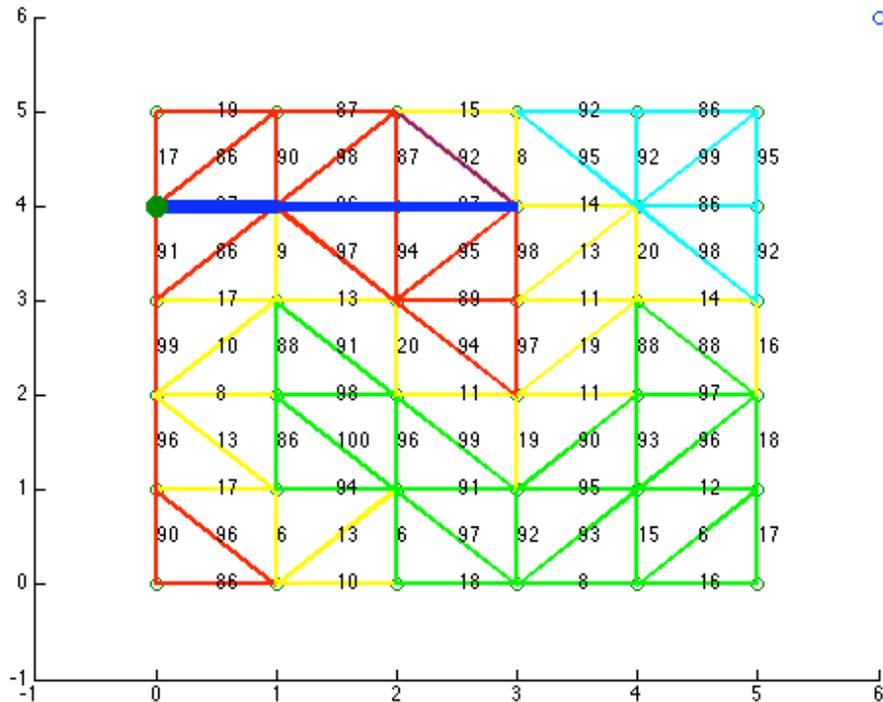
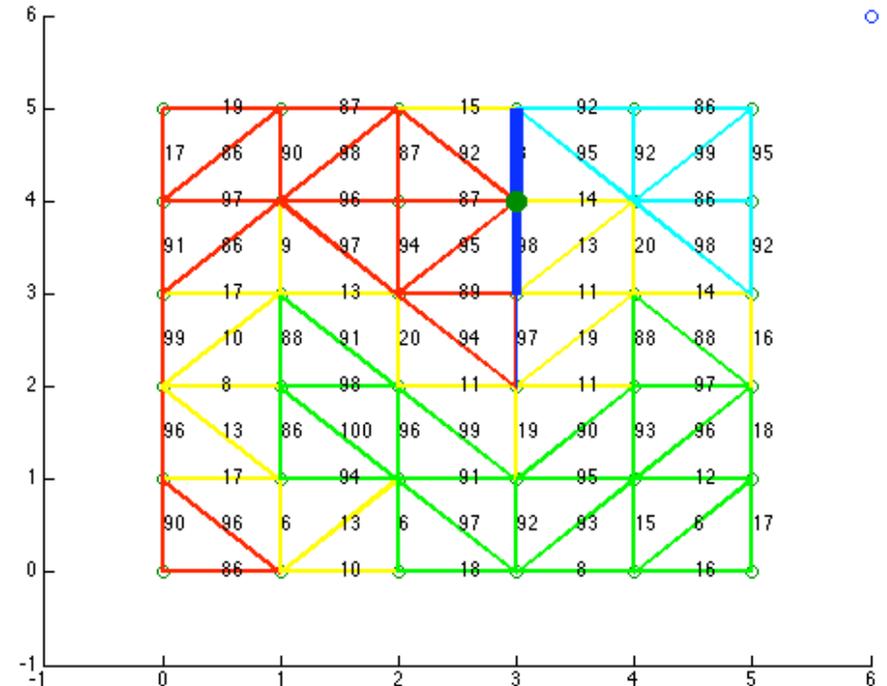


figure 2



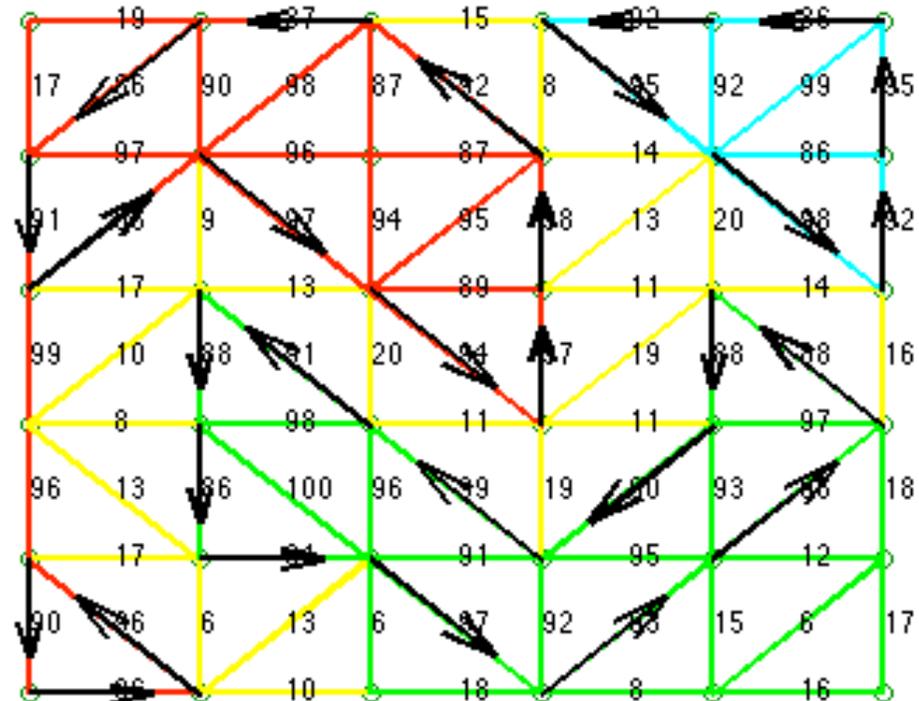
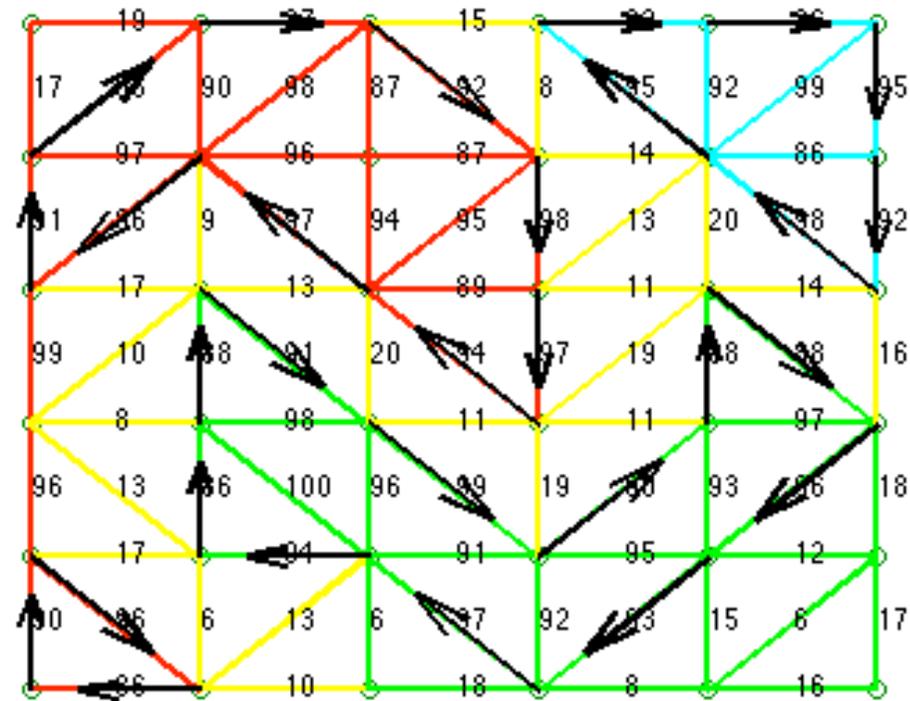
```
close all
Show=0; N=100;
figure(1); [states WR TR]=SnakeRev(N,5, 'Vg', Show, 1, 1, 1, 1, 1, 0, 0, 0, 0, 0, 0, 0);
figure(2); [states WR TR]=SnakeRev(N,5, 'Vg', Show, 0, 1, 1, 1, 1, 1, 0, 0, 0, 0, 0, 0, 0);
```

The answer:

figure 1

o

figure 2



```
close all
Show=1; N=100;
figure(1); [states WR TR]=SnakeRev(N,5,'Vg',Show,1,1,1,1,1,0,0,0,0,0,0);
figure(2); [states WR TR]=SnakeRev(N,5,'Vg',Show,0,1,1,1,1,0,0,0,0,0,0,0);
```

Reversibility

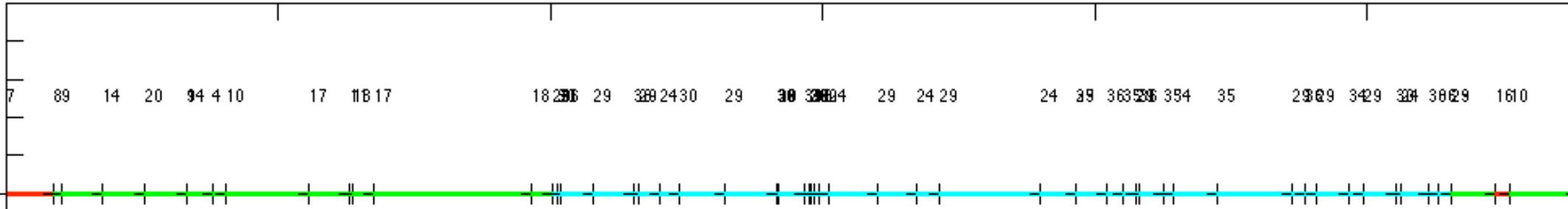
$$P^* = [\pi]^{-1} P^{tr} [\pi]$$

reversible if $P = P^*$

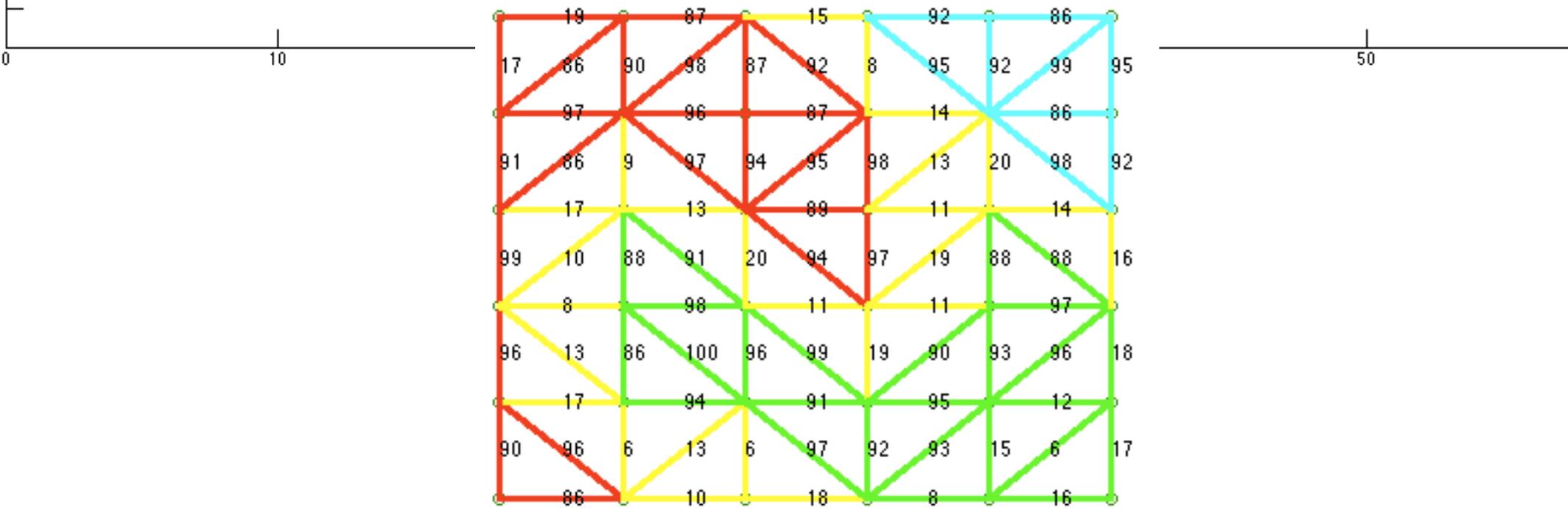
Theorem: Reversible if and only if there is symmetric conductance matrix such that

$$P = [W\mathbf{1}]^{-1}W$$

$$\pi = W\mathbf{1} / (\mathbf{1}^{tr}W\mathbf{1})$$



starting from equilibrium, if I reflected this you'd never know



```
N=200; K=10; [states]=PlotTheState(N,K);
```

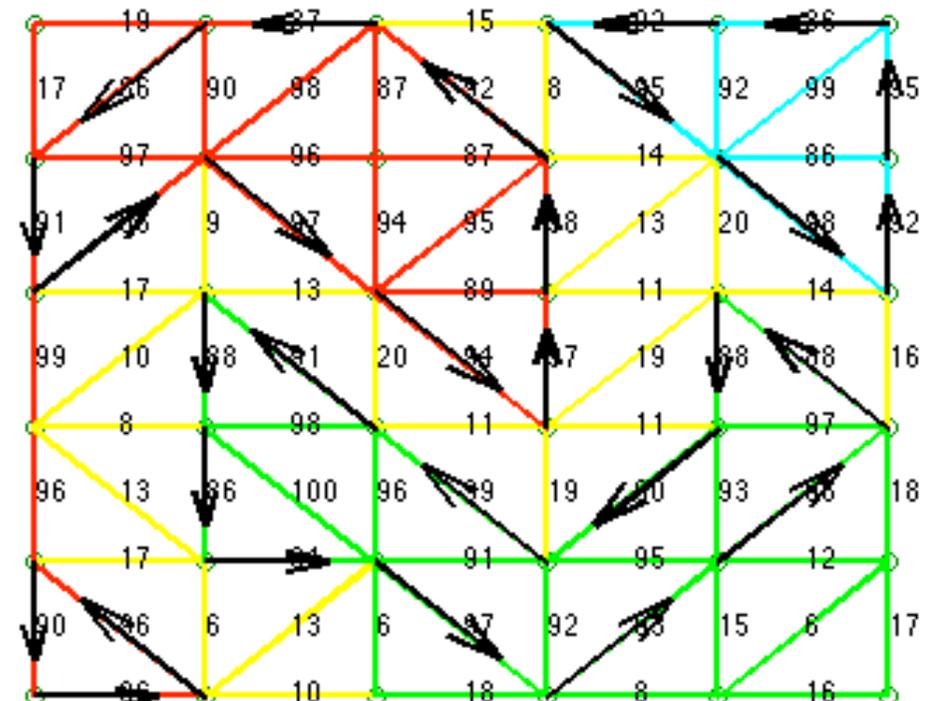
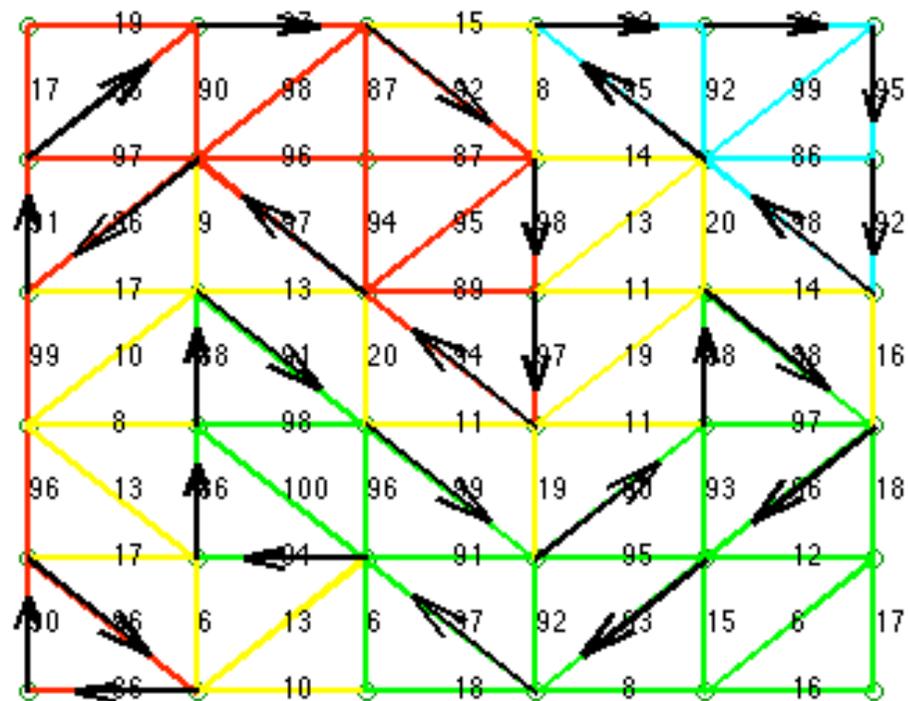
Theorem: For a reversible chain, the vagabond embedding is the PCA of the green's embedding weighted by the equilibrium vector

Theorem: For any chain P , there is a unique (up to a multiplicative constant) conductance W and divergence free, compatible flow F such that

$$P = [W\mathbf{1}]^{-1} (W + F)$$

$$P^* = [W\mathbf{1}]^{-1} (W - F)$$

$$\pi = W\mathbf{1} / (\mathbf{1}^{tr}W\mathbf{1})$$



Flow Independence Corollary: The vagabond embedding and C_ω depend only on the Conductance W and the equilibrium measure ω , and are independent of the flow F .

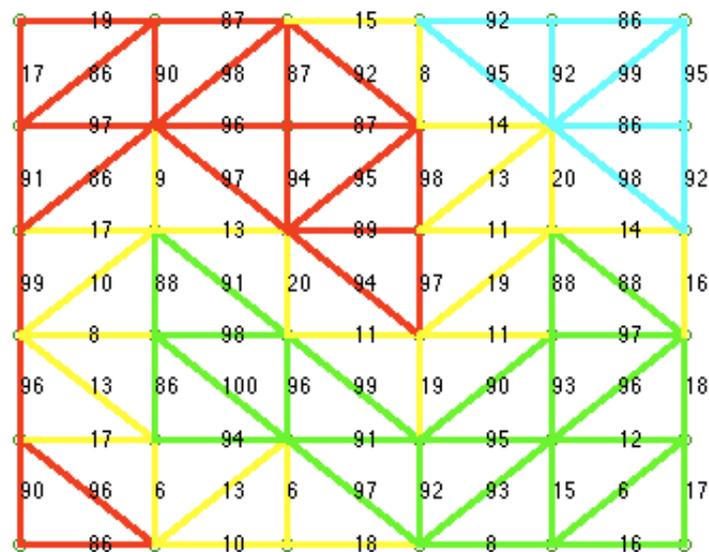
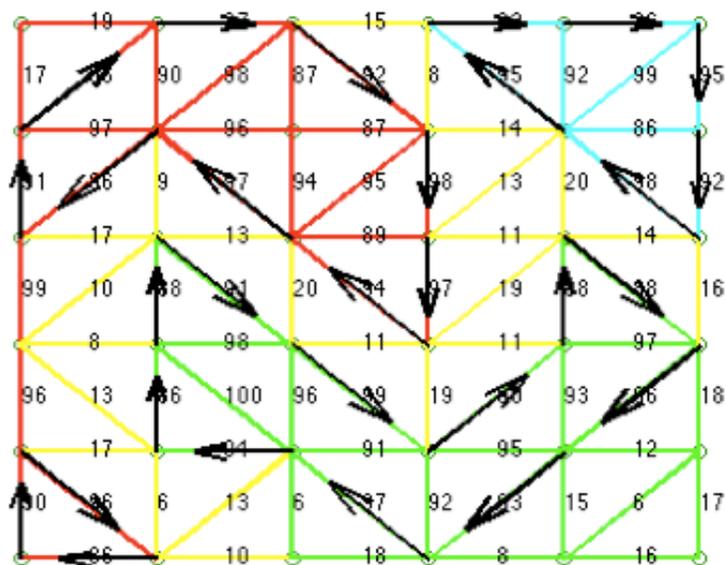
proof: Simply note

$$C_\omega = \frac{\Delta_\omega + \Delta_\omega^*}{2} = [\tau]^{-1} \left(I - \frac{P + P^*}{2} \right) = [\tau]^{-1} (I - [\pi]^{-1}W)$$

we have:

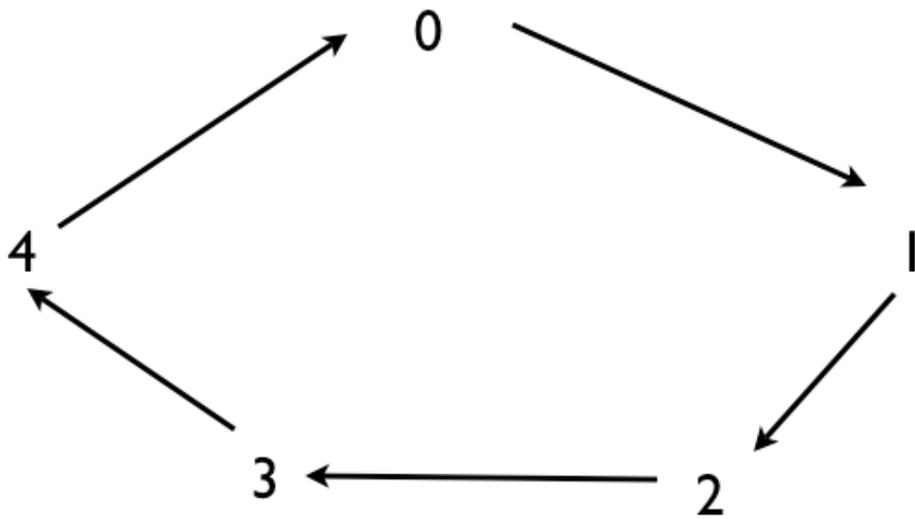
Q.E.D

Clearly the vagabond is missing half the geometry.

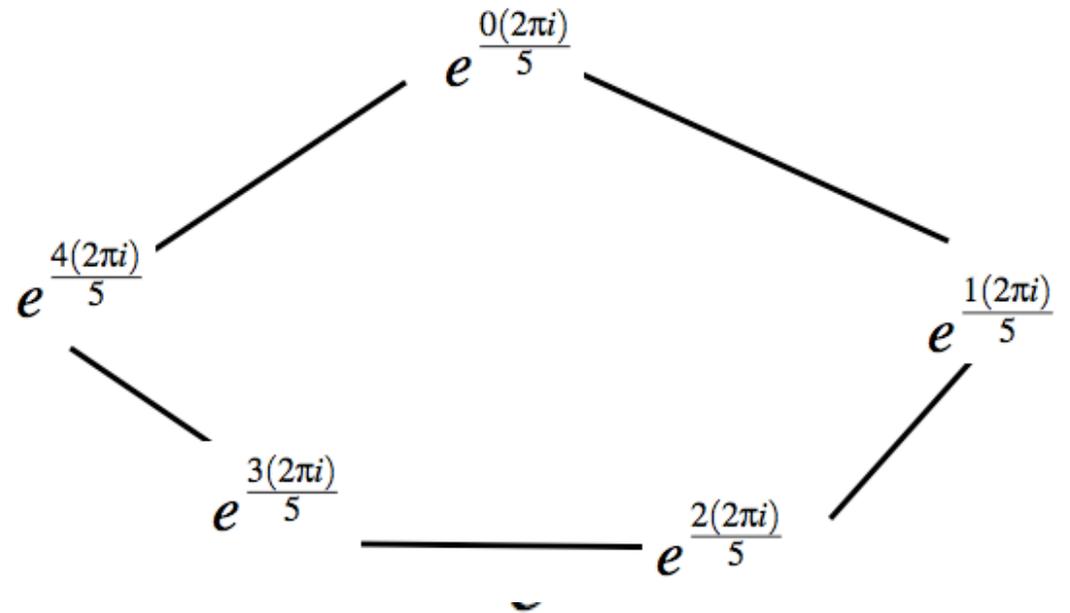


Here we are going to explore a scenario in our space. This scenario is going to be 'cycle like behavior', and hence we must ask: what does a directed cycle look like from the point of view of the functions on the chain?

chain



function



Shall we dance?

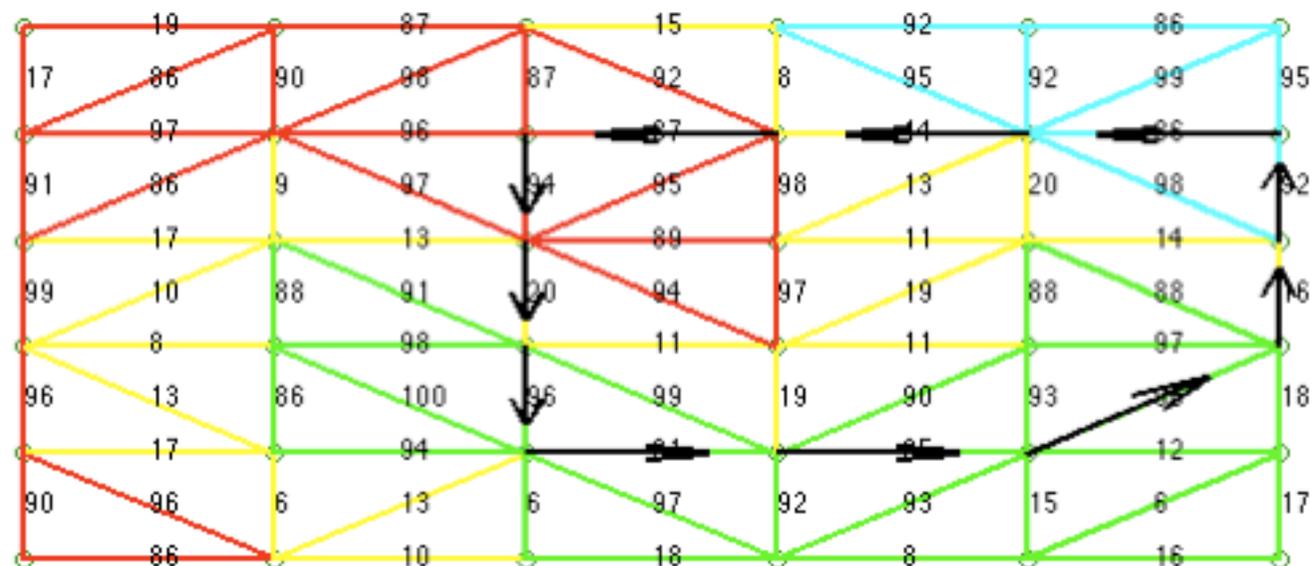
$$e^{-i\left(\frac{K2\pi}{N}\right)n} P^n f = f$$

More generally, cycles will support functions where there is a relatively large frequency κ such that

$$e^{-i\kappa t} e^{-t\Delta} f$$

is pointwise relatively stable. One way to find largish κ such the function f and κ pair locally minimize $\left\| \frac{d}{dt} e^{-(i\kappa+\Delta)t} f \right\|_{\omega}^2$ at $t = 0$. So we are looking for critical points of

$$F(\kappa, f) = \left\| \frac{d}{dt} e^{-(i\kappa+\Delta)t} f \right\|_{\omega}^2 \Big|_{t=0}.$$



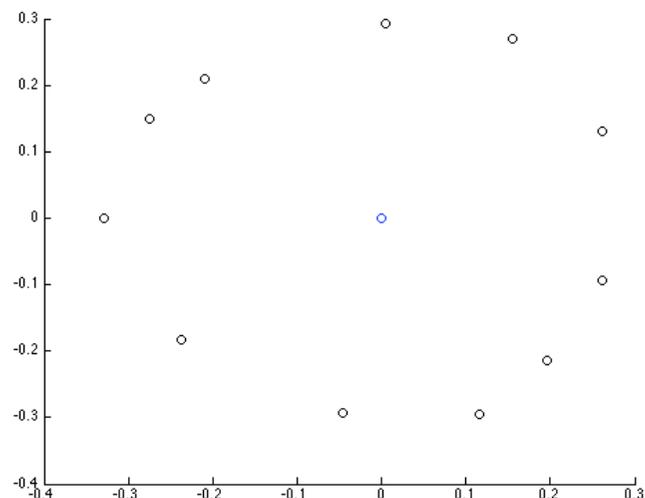
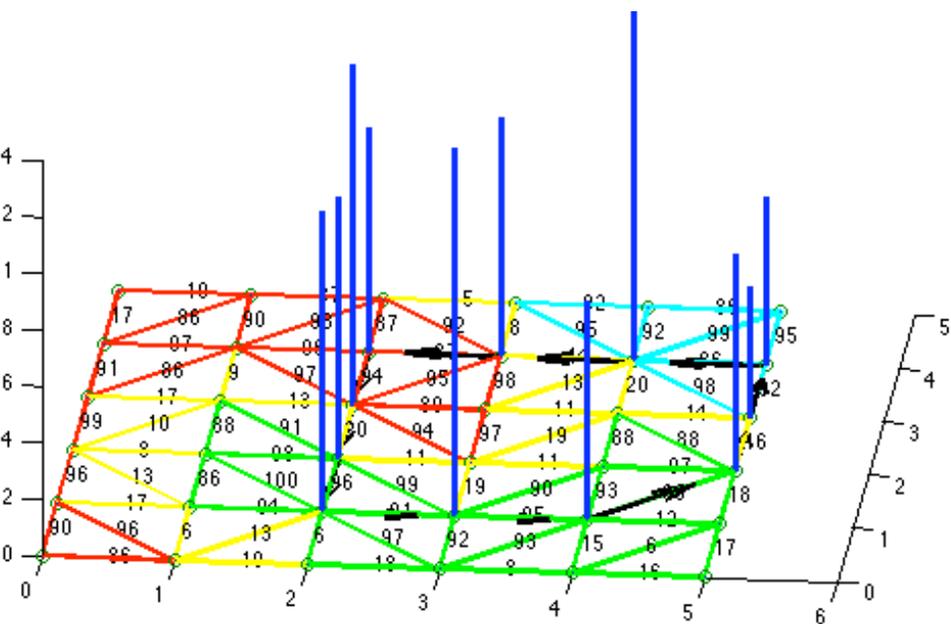
More generally, cycles will support functions where there is a frequency κ such that

$$e^{-i\kappa t} e^{-t\Delta} f$$

is pointwise relatively stable. One way to find such functions would be to minimize $\left\| \frac{d}{dt} e^{-(i\kappa+\Delta)t} f \right\|_{\omega}^2$ at $t=0$.

Hence we are looking for critical points of

$$F(\kappa, f) = \left\| \frac{d}{dt} e^{-(i\kappa+\Delta)t} f \right\|_{\omega}^2 \Big|_{t=0},$$

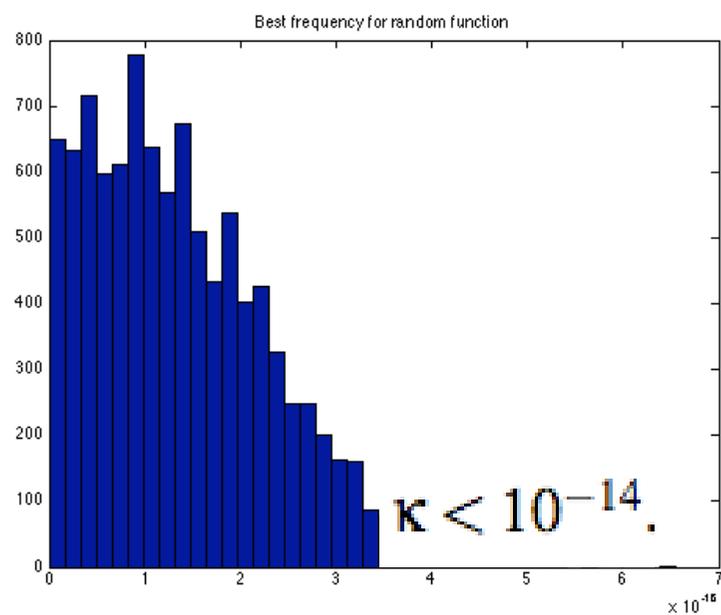
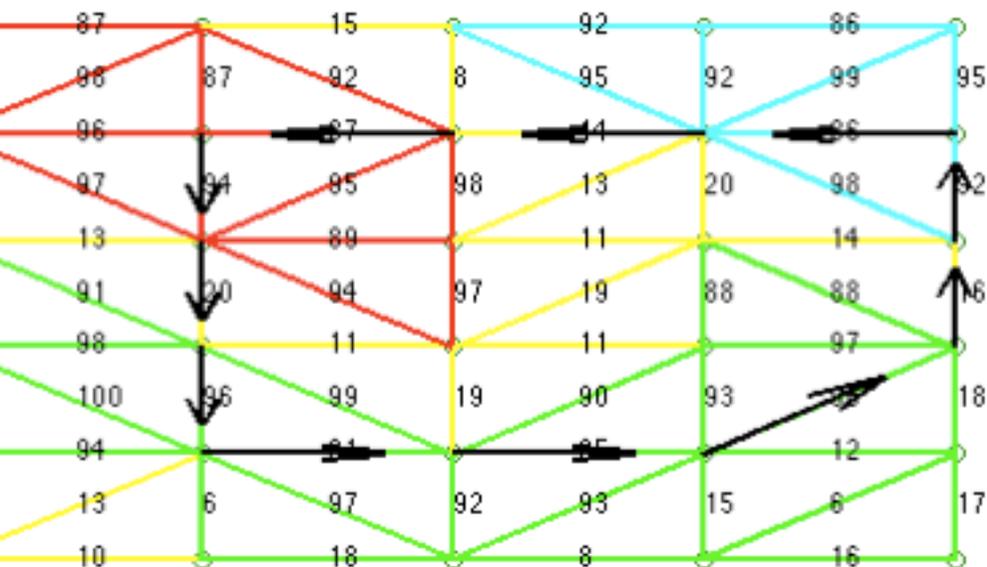


$$\kappa = 0.541126346528966$$

We might call these local minima cycle functions. To get a sense for 'largish', putting a cycle on our toy chain as in figure 1, we can compare the function in the figure 5, which for $\kappa = 0.541126346528966$ has a nearly zero derivative, to a random function (real and imaginary parts random in $[0, 1]$ normalized to have norm 1). For a random function we find its optimal κ for minimizing the derivative and after 10000 runs the largest such $\kappa < 10^{-14}$.

$$\kappa = 0.541126346528966$$

verses



How to find them...

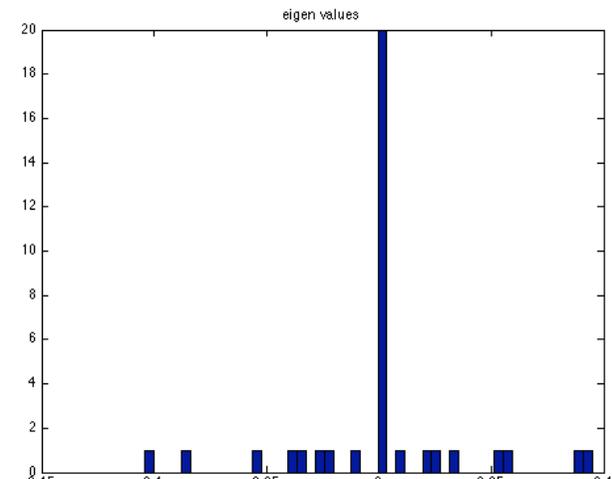
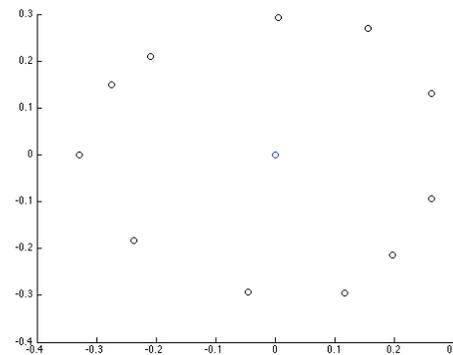
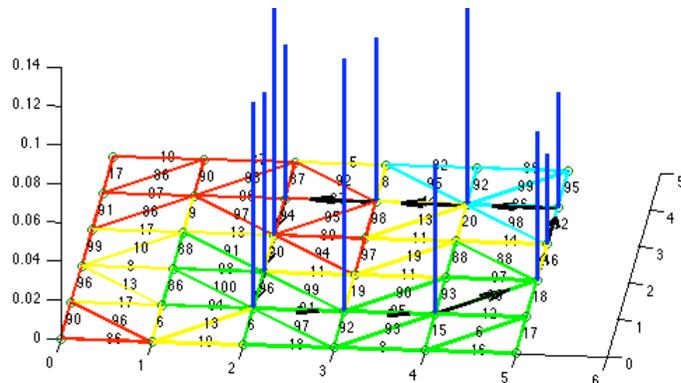
Cycle Detection Lemma: If $B_\omega f = \kappa f$ where $S_\omega = i \frac{\Delta - \Delta^*}{2}$, then $\frac{d}{d\kappa} F = 0$.

proof:

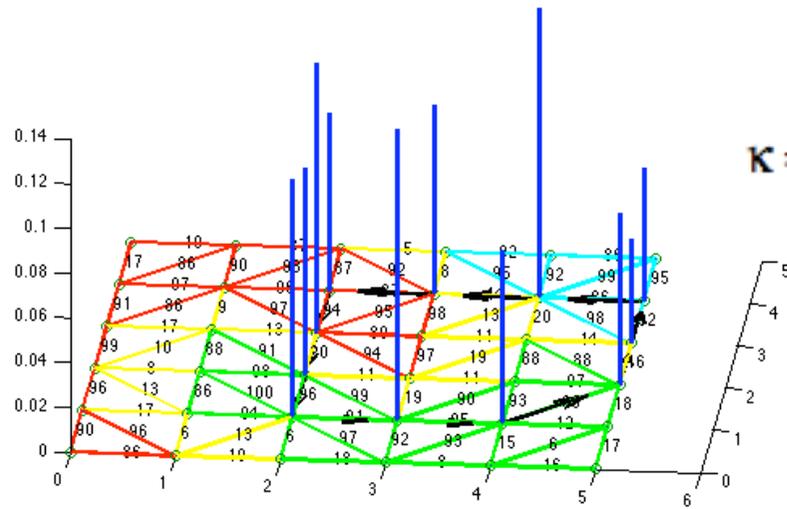
$$\begin{aligned}
 0 &= \frac{d}{d\kappa} \left\| \left. \frac{d}{dt} e^{-(i\kappa + \Delta)t} f \right|_{t=0} \right\|_\omega^2 && \text{(definition)} \\
 &= \frac{d}{d\kappa} \left\| (i\kappa + \Delta) e^{-(i\kappa + \Delta)t} f \right\|_\omega^2 \Big|_{t=0} && \text{(matrix derivative, sec 1.1 ex??)} \\
 &= \frac{d}{d\kappa} \left\| (i\kappa + \Delta) f \right\|_\omega^2 && (e^{At} = I \text{ at } t = 0) \\
 &= \frac{d}{d\kappa} \langle (i\kappa + \Delta) f, (i\kappa + \Delta) f \rangle_\omega && \text{(def. } \langle, \rangle_\omega) \\
 &= \langle if, (i\kappa + \Delta) f \rangle_\omega + \langle (i\kappa + \Delta) f, if \rangle_\omega && \text{(chain rule)} \\
 &= \langle f, (\kappa - i\Delta) f \rangle_\omega + \langle (\kappa - i\Delta) f, f \rangle_\omega && \text{(conjugate linear)} \\
 &= \langle f, ((\kappa - i\Delta) + (\kappa - i\Delta)^*) f \rangle_\omega && \text{(def. adjoint)} \\
 &= \langle f, 2(\kappa - S_\omega) f \rangle_\omega && \text{(compute adjoint)}
 \end{aligned}$$

Q.E.D

$\kappa = 0.541126346528966$

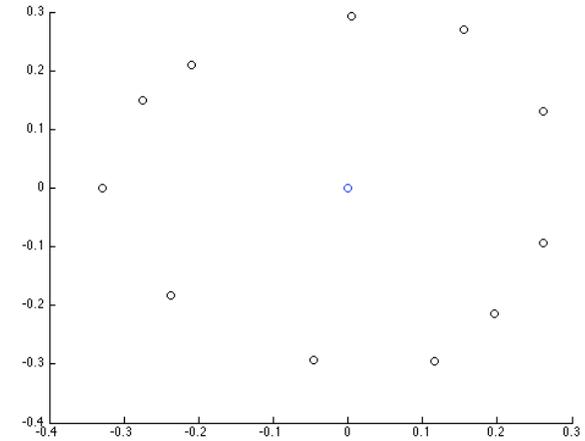


The scenario operator: $S_{\omega} = i \frac{\Delta - \Delta^*}{2}$

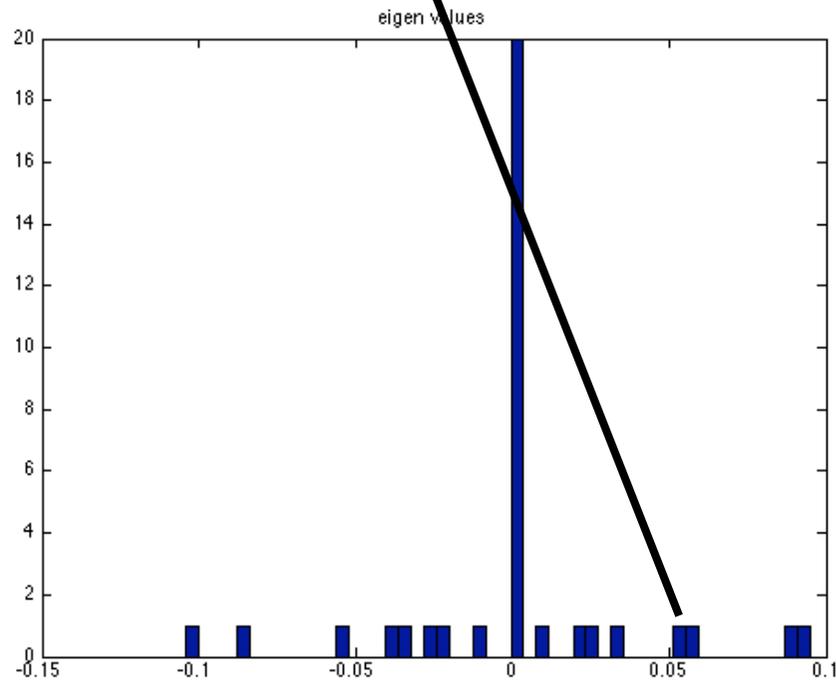


$\kappa = 0.541126346528966$

Text



Scenario operators
have highly localized
eigenfunctions



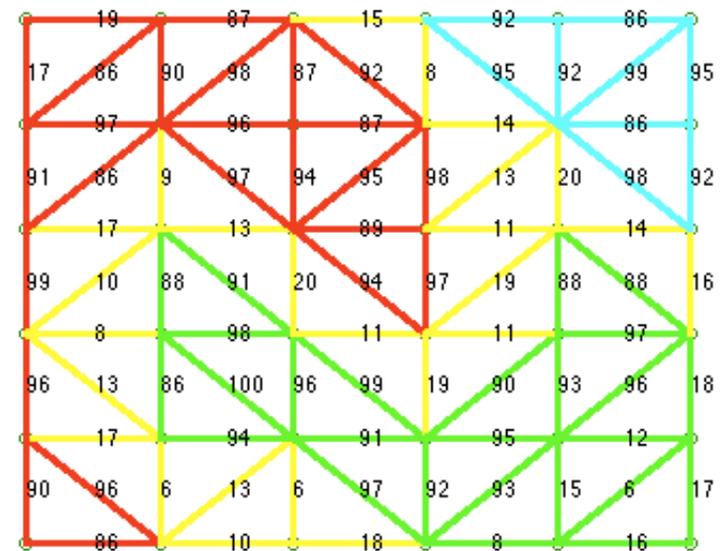
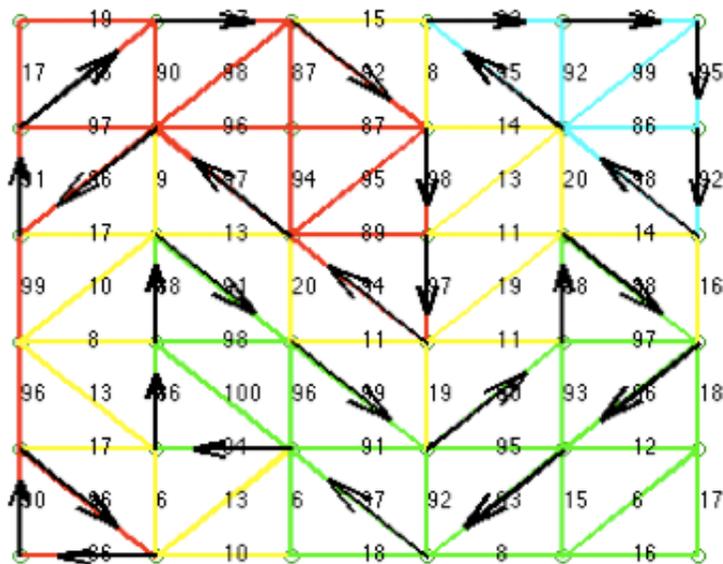
The scenario operator: $S_{\omega} = i \frac{\Delta - \Delta^*}{2}$

Conductance Independence Corollary: S_{ω} depends only on the flow F and equilibrium vector ω and is independent of conductance C . Furthermore the chain is reversible if and only if $S_{\omega} = 0$.

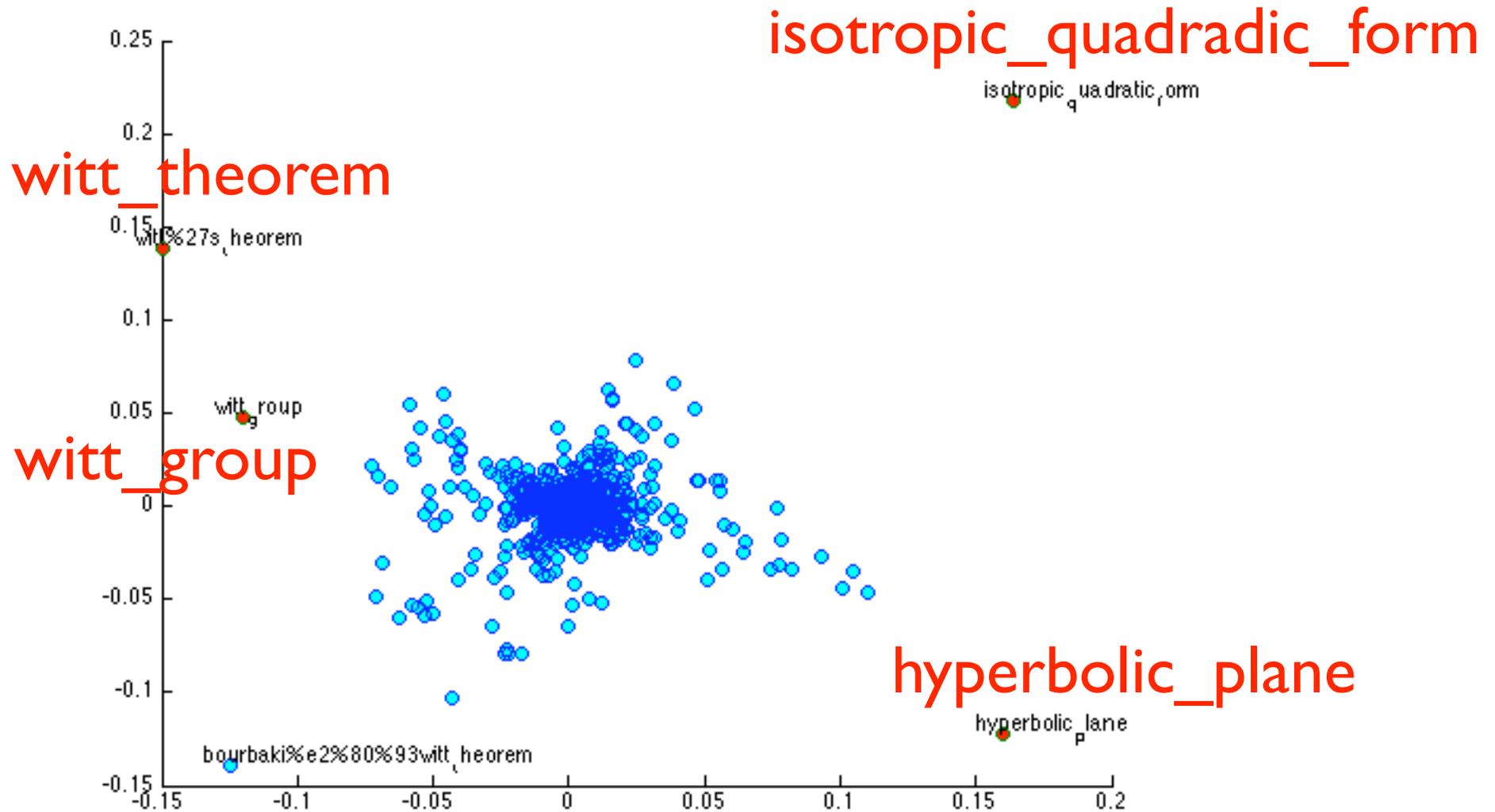
proof:

$$S_{\omega} = i \frac{\Delta_{\omega} - \Delta_{\omega}^*}{2} = i[\tau]^{-1} \left(\frac{P - P^*}{2} \right) = i[\tau]^{-1}[\pi]^{-1}F = i[\omega]^{-1}F$$

Q.E.D

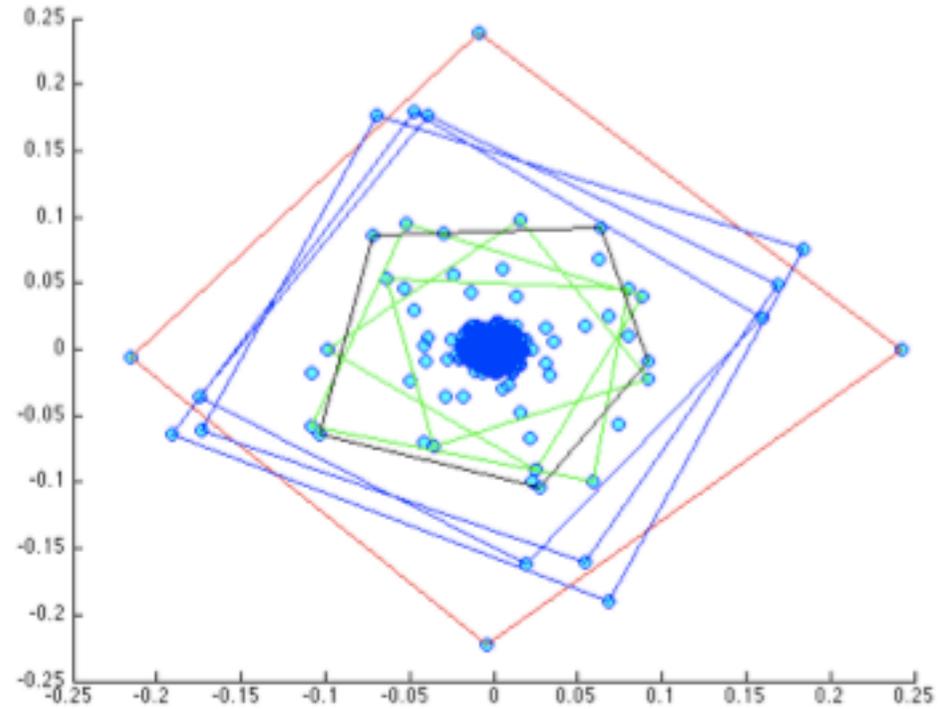
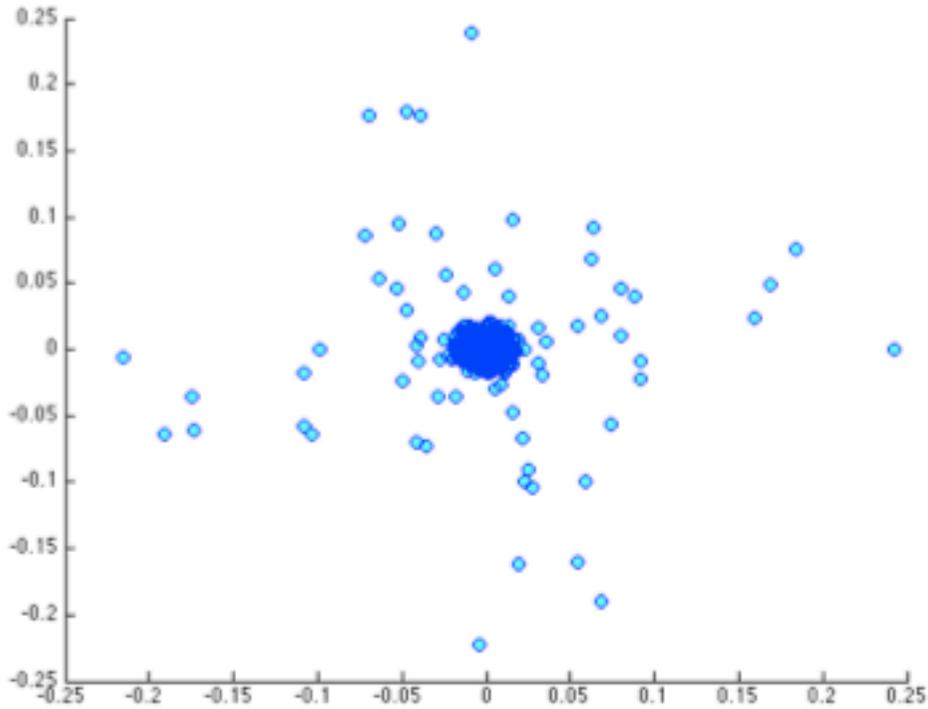


Eigenfunction in the Space of Mathematics



Scenario operators
have highly localized
eigenfunctions

Reality....



We need something like the vagabond clustering for cycles

The Co-conformal Cycle Hunt

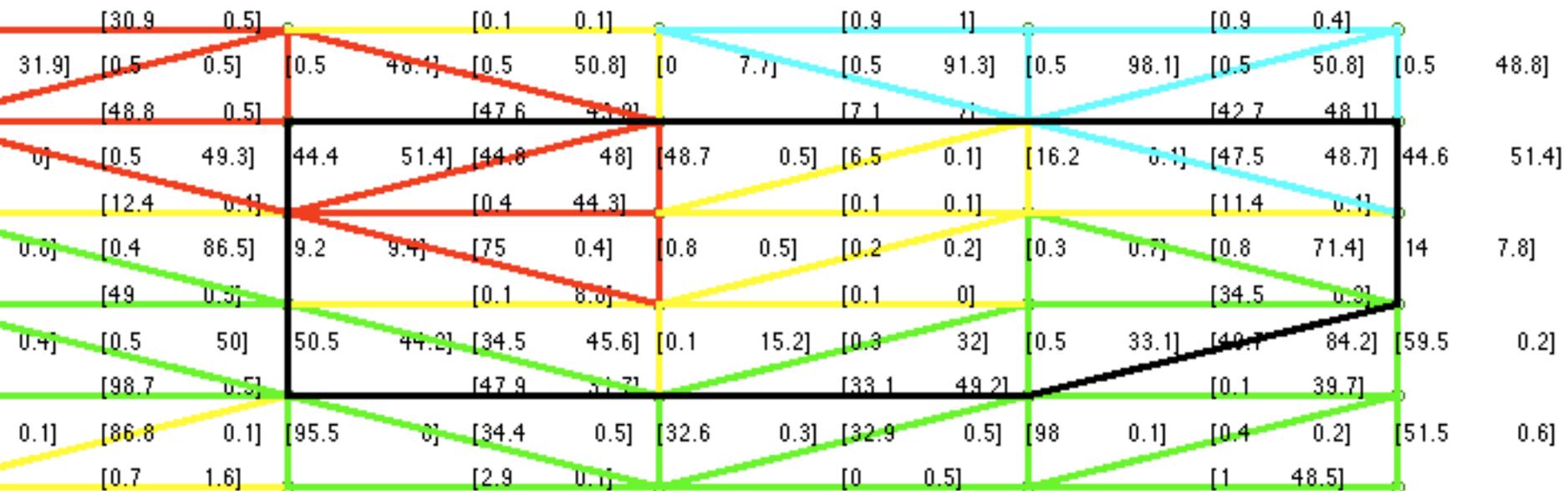
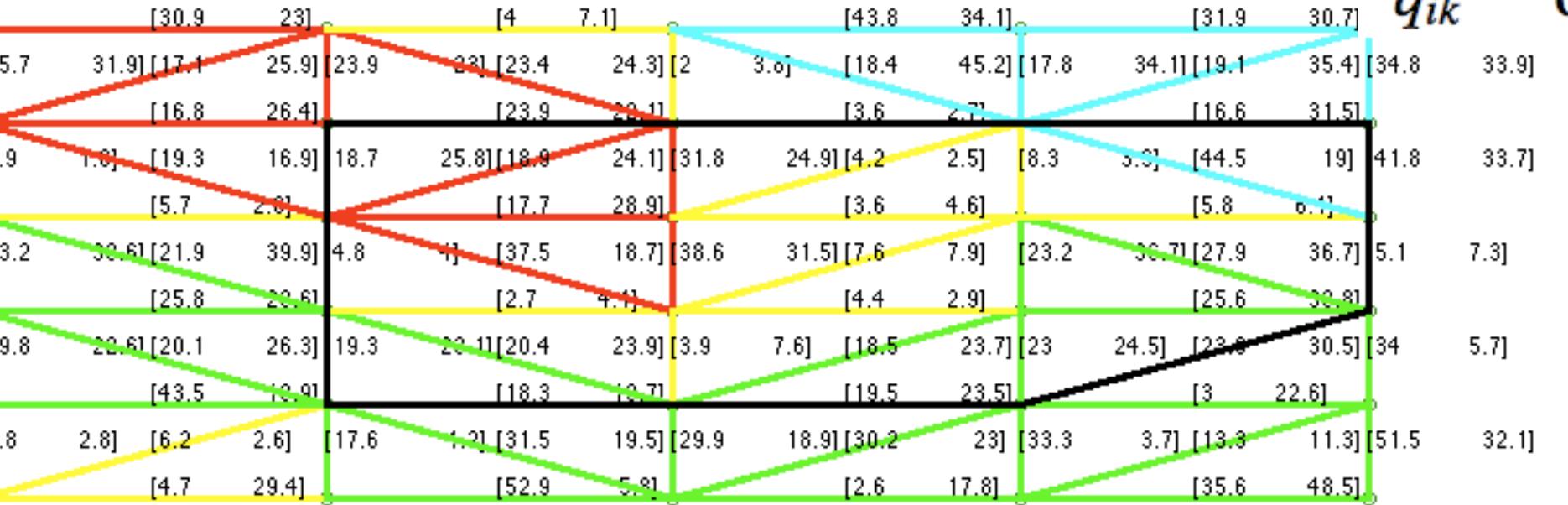


Co-conformal Magnetization

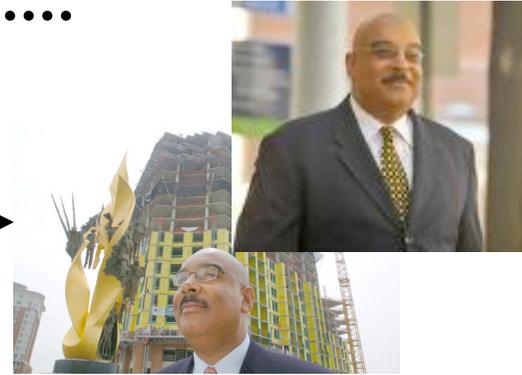
$$Q = [P[\alpha]\mathbf{1}]^{-1}P[\alpha]$$

$$\underline{q}_{ij} = \frac{\alpha_j}{\alpha_k} \underline{p}_{ij}$$

$$\underline{q}_{ik} = \frac{\alpha_j}{\alpha_k} \underline{p}_{ik}$$



Not all cycles are created equally...



Why do we need this fellow



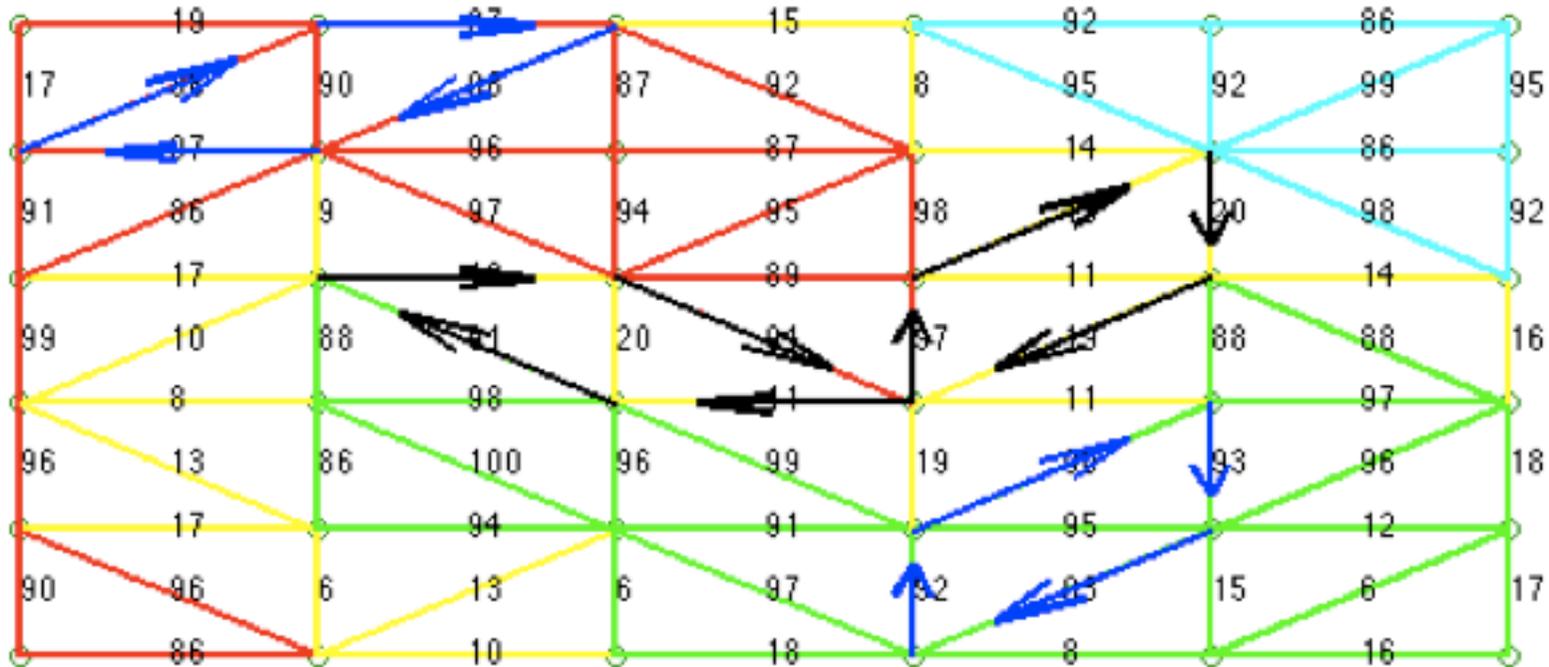
to find a naughty cycle?



Claim: The operator $i[\Delta_\omega, \Delta_\omega^*] = \Delta_\omega \Delta_\omega^* - \Delta_\omega^* \Delta_\omega$

allows us to detect scenarios which are
anomalous relative to the context

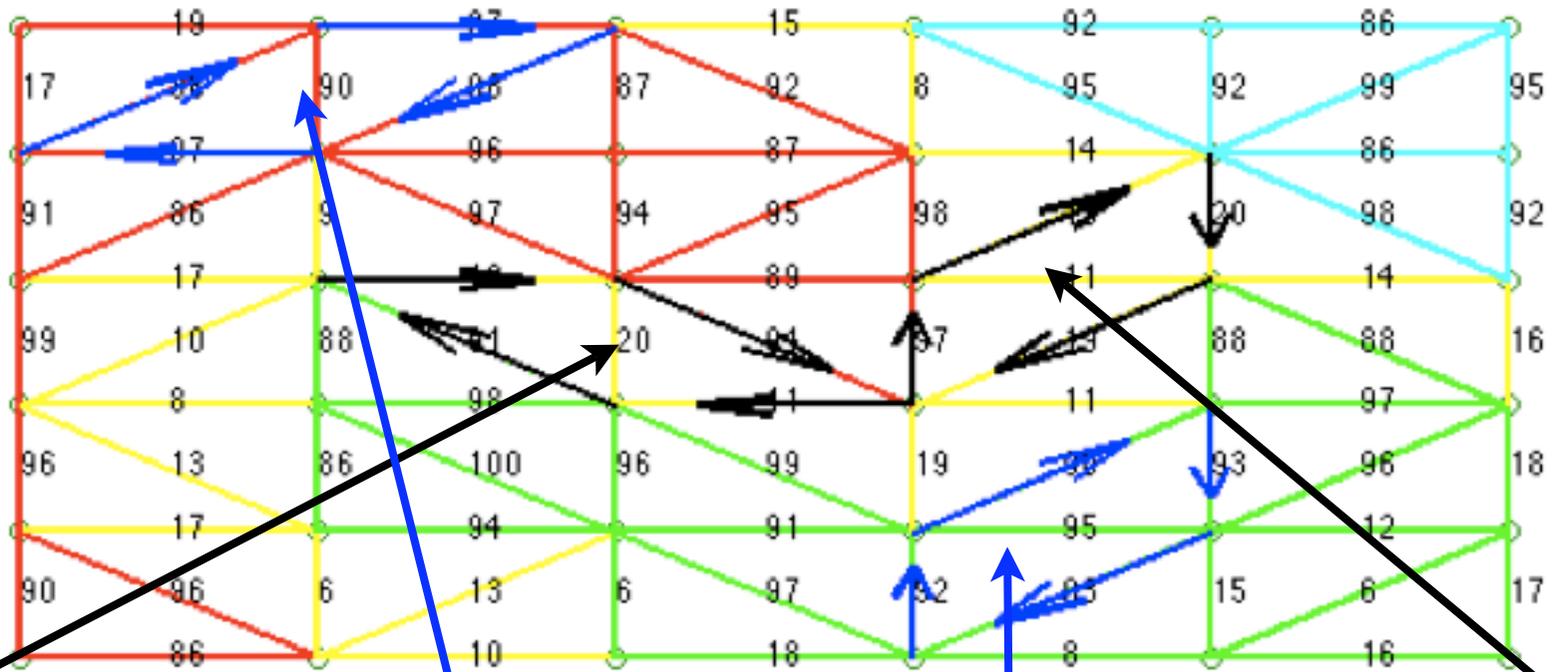
(all cycles weight 10)



Which cycles are in context and which are not?

Claim: The operator $i[\Delta_\omega, \Delta_\omega^*] = \Delta_\omega \Delta_\omega^* - \Delta_\omega^* \Delta_\omega$

allows us to detect scenarios which are anomalous relative to the context



$$\left| \langle [\Delta_\omega, \Delta_\omega^*] \rangle_{\chi_{C_{l,k}}} \right| = 0.5901$$

$$\left| \langle [\Delta_\omega, \Delta_\omega^*] \rangle_{\chi_{C_{r,k}}} \right| = 0.3157$$

$$\left| \langle [\Delta_\omega, \Delta_\omega^*] \rangle_{\chi_{C_{l,b}}} \right| = 0.0039$$

$$\left| \langle [\Delta_\omega, \Delta_\omega^*] \rangle_{\chi_{C_{u,b}}} \right| = 0.0141$$



<http://www.youtube.com/watch?v=KT7xJ0tjB4A>

Mathematics of Quantum Mechanics

$$\langle f, g \rangle_{\omega} = \bar{f}^{tr}[\omega]g$$

Hilbert Space

$$\langle \psi, \psi \rangle = 1$$

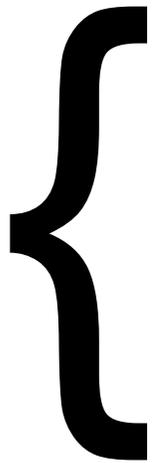
States

$$\langle \psi, A\phi \rangle = \langle A\psi, \phi \rangle$$

Measurements

$$\langle \psi, A\psi \rangle_{\omega} = \langle A \rangle_{\psi}$$

Expectation



$$A_{\omega} - \langle A_{\omega} \rangle_{\psi}$$

$$\langle \hat{A}^2 \rangle_{\psi}$$

$$\sigma_{\psi A} = \sqrt{\langle \hat{A}^2 \rangle_{\psi}}$$

Deviation

Why does this work?

The Uncertainty Principle for Markov Chains Lemma:

$$(\sigma_\psi C_\omega) (\sigma_\psi S_\omega) \geq \left| \frac{1}{2} \langle [\Delta_\omega, \Delta_\omega^*] \rangle_\psi \right|$$

proof: Notice

$$[C_\omega, S_\omega] = C_\omega S_\omega - S_\omega C_\omega = i[\Delta_\omega, \Delta_\omega^*],$$

so

$$\begin{aligned} \left| \frac{1}{2} \langle [\Delta_\omega, \Delta_\omega^*] \rangle_\psi \right| &= \left| \frac{1}{2} \langle [C_\omega, S_\omega] \rangle_\psi \right| && \text{(def. plus foil)} \\ &= \left| \frac{1}{2} \langle [\hat{C}_\omega, \hat{S}_\omega] \rangle_\psi \right| && \text{(differ by constants)} \\ &= \frac{1}{2} |\langle \psi, \hat{C}_\omega \hat{S}_\omega \psi - \hat{C}_\omega \hat{S}_\omega \psi \rangle| && \text{(definition)} \\ &= \frac{1}{2} |\langle \hat{C}_\omega \psi, \hat{S}_\omega \psi \rangle - \langle \hat{S}_\omega \psi, \hat{C}_\omega \psi \rangle| && \text{(self adjoint)} \\ &= |\text{Im} \langle \hat{C}_\omega \psi, \hat{S}_\omega \psi \rangle| && \text{(Hermitian Inner-product)} \\ &\leq |\langle \hat{C}_\omega \psi, \hat{S}_\omega \psi \rangle| && (|a + ib| = \sqrt{a^2 + b^2}) \\ &\leq \sqrt{\langle \hat{C}_\omega \psi, \hat{C}_\omega \psi \rangle} \sqrt{\langle \hat{S}_\omega \psi, \hat{S}_\omega \psi \rangle} && \text{(Cauchy-Schwarz inequality)} \\ &= \sqrt{\langle \psi, \hat{C}_\omega^2 \psi \rangle} \sqrt{\langle \psi, \hat{S}_\omega^2 \psi \rangle} && \text{(Self adjoint)} \\ &= (\sigma_\psi C_\omega) (\sigma_\psi S_\omega) && \text{(definition)} \end{aligned}$$

Q.E.D