

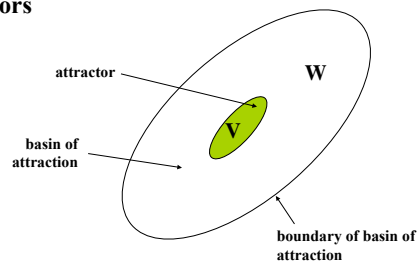
So far: mostly about *maps*.

- discrete time systems:
 - time proceeds in clicks
 - “maps”
 - modeling tool: *difference* equation

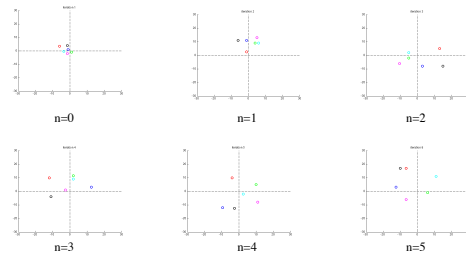
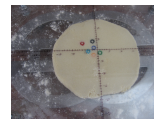
Next up: *flows*

- continuous time systems:
 - time proceeds smoothly
 - “flows”
 - modeling tool: *differential* equations

Attractors



- Attractors exist only in dissipative systems!
- Dissipation \iff contraction of state space under the influence of the dynamics
- Can still have chaos if no dissipation...just not chaotic *attractors*



Conditions for chaos in continuous-time systems

Necessary:

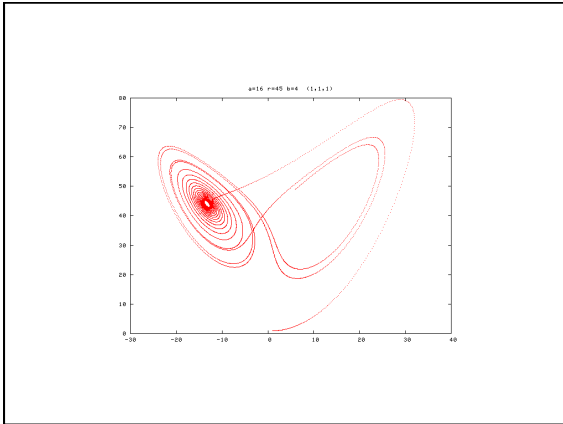
- Nonlinear
- At least three state-space dimensions (NB: only one needed in maps)

Necessary and sufficient:

- “Nonintegrable”
i.e., cannot be solved in closed form

Concepts: review

- State variable
- State space
- Initial condition
- Trajectory
- Attractor
- Basin of attraction
- Transient
- Fixed point (un/stable)
- Bifurcation
- Parameter



A cool Lorenz applet:

www.exploratorium.edu/complexity/java/lorenz.html

(Note: by Jim Crutchfield, another SFI person, who will be here at the end of next week)

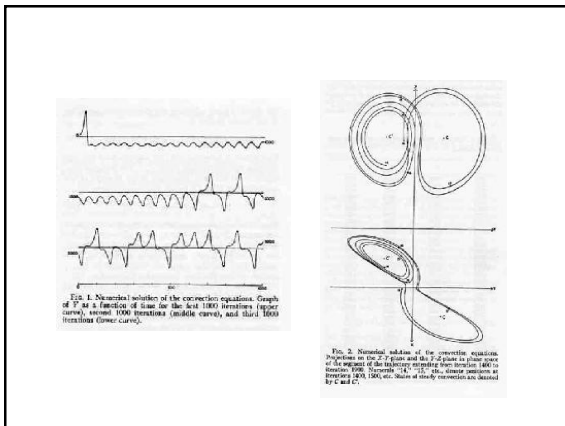
Deterministic Nonperiodic Flow¹

EDWARD N. LORENZ
Massachusetts Institute of Technology
(Manuscript received 18 November 1962, in revised form 7 January 1963)

ABSTRACT

Finite systems of deterministic ordinary nonlinear differential equations may be designed to represent forced dissipative hydrodynamic flow. Solutions of these equations can be identified with trajectories in phase space. For those systems with bounded solutions, it is found that nonperiodic solutions are ordinarily unstable with respect to small modifications, so that slightly differing initial states can evolve into considerably different states. Systems with bounded solutions are shown to possess bounded numerical solutions. A simple system representing cellular convection is solved numerically. All of the solutions are found to be unstable, and almost all of them are nonperiodic. The feasibility of very-long-range weather prediction is examined in the light of these results.

J. Atm. Sci. **20**:130



- Equations:

$$x' = a(y-x)$$


$$y' = rx - y - xz$$

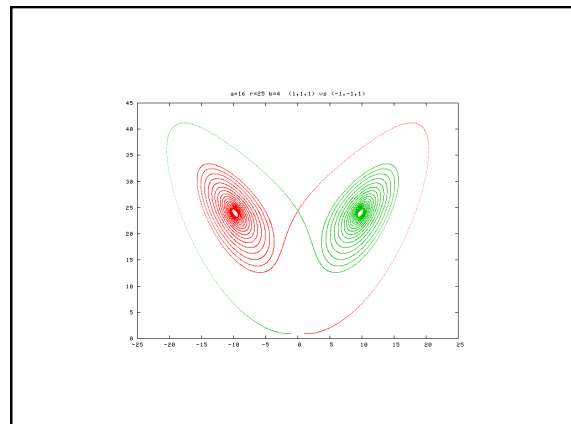
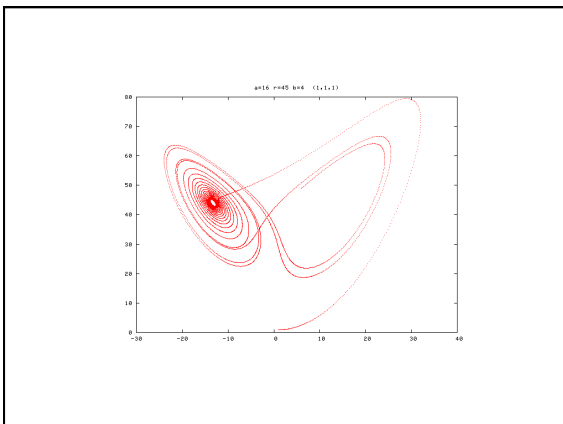
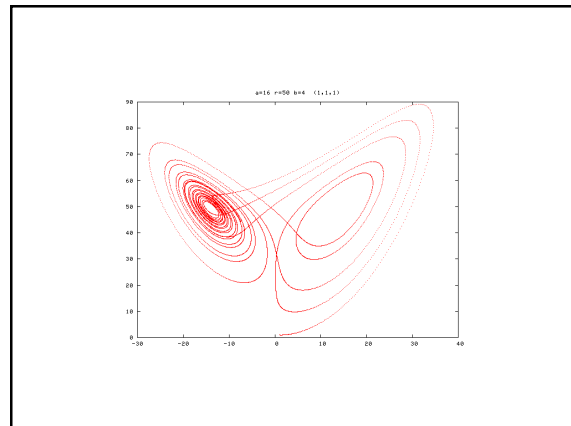
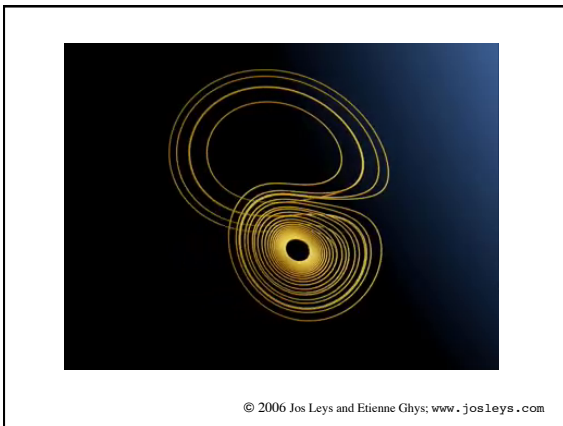
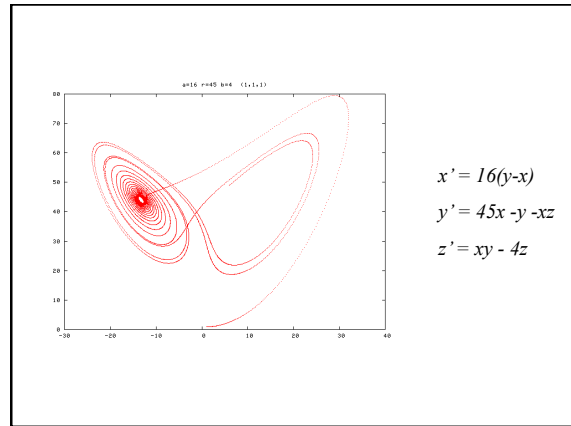
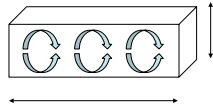
$$z' = xy - bz$$

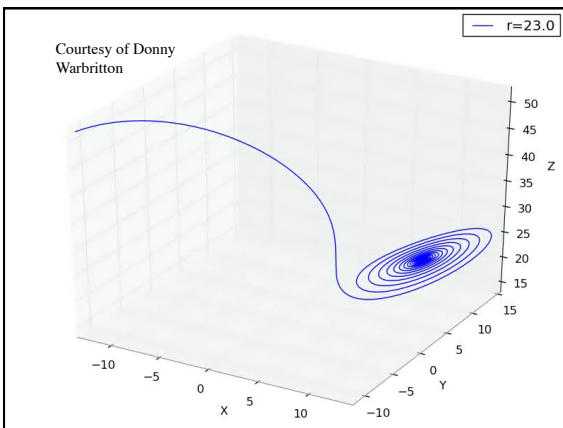
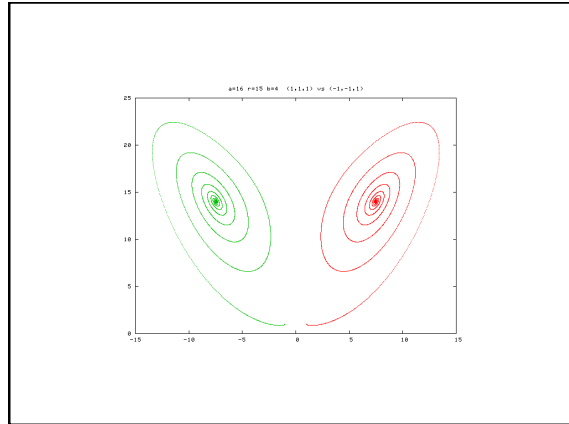
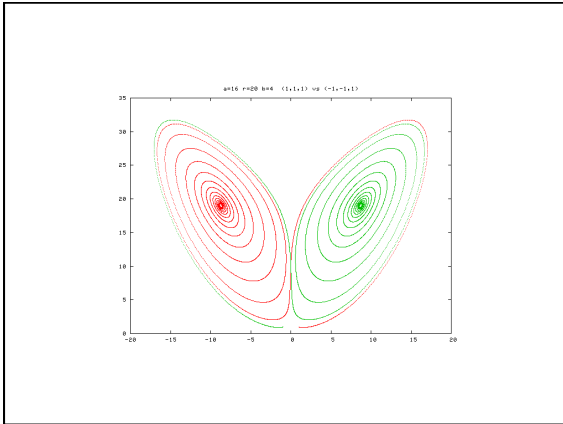
(first three terms of a Fourier expansion of the Navier-Stokes eqns)

- State variables:
 - x convective intensity
 - y temperature
 - z deviation from linearity in the vertical convection profile

Parameters:

- a Prandtl number - fluids property
- r Rayleigh number - related to ΔT 
- b aspect ratio of the fluid sheet





Attractors

Four types:

- fixed points
- limit cycles (*aka* periodic orbits)
- quasiperiodic orbits
- chaotic attractors

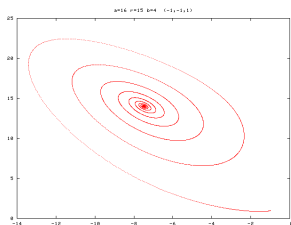
A nonlinear system can have any number of attractors, of all types, sprinkled around its state space

Their basins of attraction (plus the basin boundaries) *partition* the state space

And there's no way, *a priori*, to know where they are, how many there are, what types, etc.

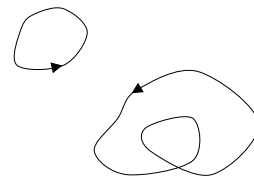
Attractors

- Fixed point



Attractors

- Limit cycle

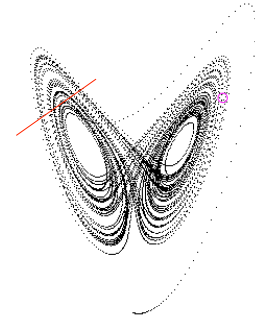


Attractors

- Quasi-periodic orbit...

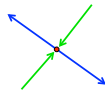
“Strange” or chaotic attractors

- *often* fractal
- covered densely by trajectories
- exponential divergence of neighboring trajectories...



Lyapunov exponents

- nonlinear analogs of eigenvalues: one λ for each dimension

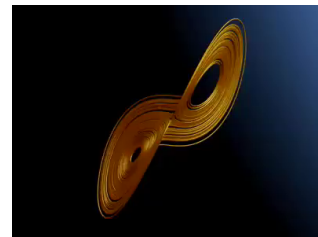
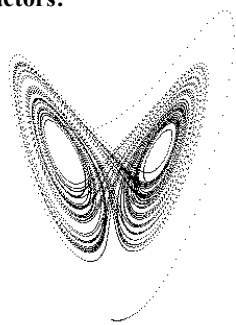


Lyapunov exponents: summary

- nonlinear analogs of eigenvalues: one λ for each dimension
- negative λ_i , compress state space; positive λ_i , stretch it
- $\sum \lambda_i < 0$ for dissipative systems
- long-term average in definition; biggest one dominates as $t \rightarrow \infty$
- *positive λ is a signature of chaos*
- λ_i are same for all ICs in one basin

“Strange” or chaotic attractors:

- exponential divergence of neighboring trajectories
- *often* fractal
- covered densely by trajectories
- contain an infinite number of “unstable periodic orbits”...



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