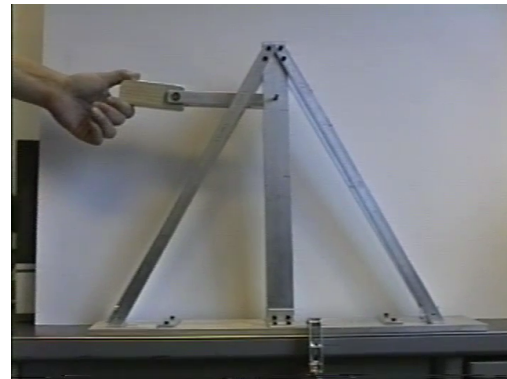


Introduction to Nonlinear Dynamics

Santa Fe Institute
Complex Systems Winter School
8-9 December 2015

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Elizabeth Bradley

Liz Bradley
lizb@cs.colorado.edu

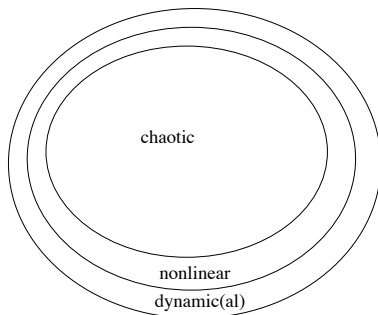


<http://ayresriverblog.com>

Chaos

Complex behavior, arising in a deterministic nonlinear dynamic system, which exhibits two special properties:

- sensitive dependence on initial conditions
- characteristic structure...



Chaos

Complex behavior, arising in a deterministic nonlinear dynamic system, which exhibits two special properties:

- sensitive dependence on initial conditions
- characteristic structure...

Systems that exhibit chaos are ubiquitous; many of them are also simple, well-known, and "well-understood"

Where nonlinear dynamics turns up

- Flows (of fluids, heat, ...)
 - Eddy in creek
 - Weather
 - Vortices around marine invertebrates
 - Air/fuel flow in combustion chambers



Where nonlinear dynamics turns up

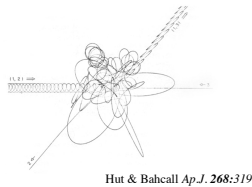
- Driven nonlinear oscillators
 - Pendula
 - Hearts
 - Fireflies



- and lots of other electronic, chemical, & biological systems

Where nonlinear dynamics turns up

- Classical mechanics
 - three-body problem
 - paired black holes
 - pulsar emission
 -
- Protein folding
- Population biology
- And many, many other fields (including yours)



- discrete time systems:
 - time proceeds in clicks
 - “maps”
 - modeling tool: difference equation
- continuous time systems:
 - time proceeds smoothly
 - “flows”
 - modeling tool: differential equations



A useful graphical solution technique

- “cobweb” diagram
- *aka* return map
- *aka* correlation plot

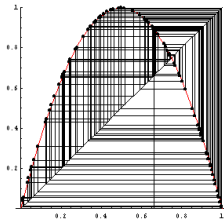
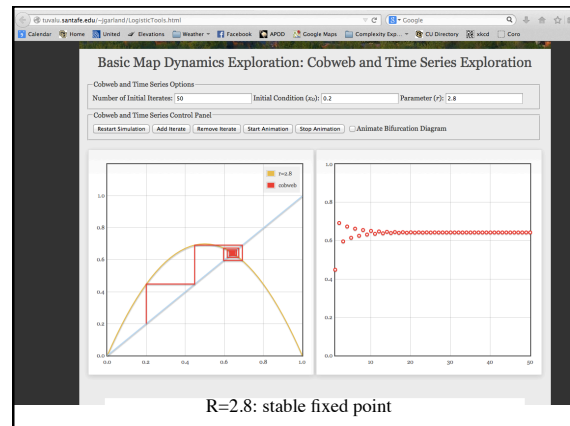


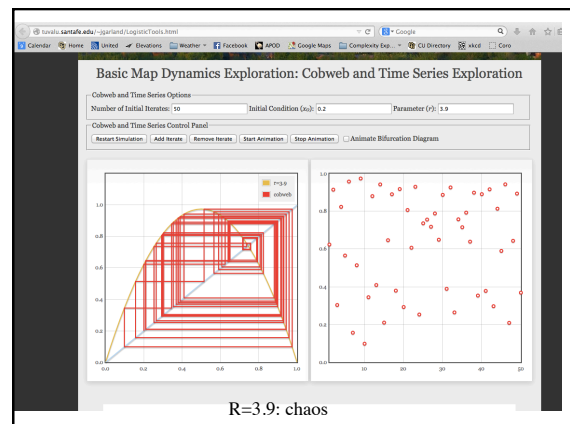
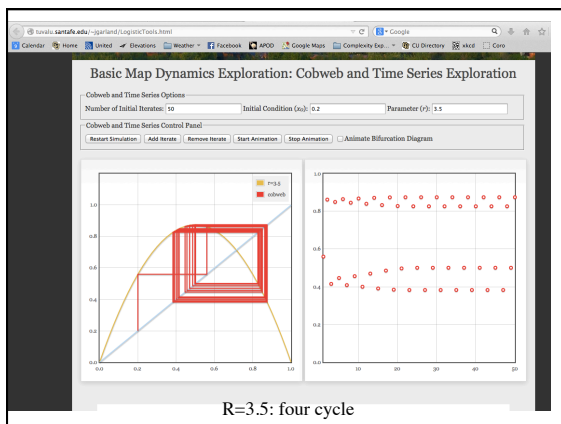
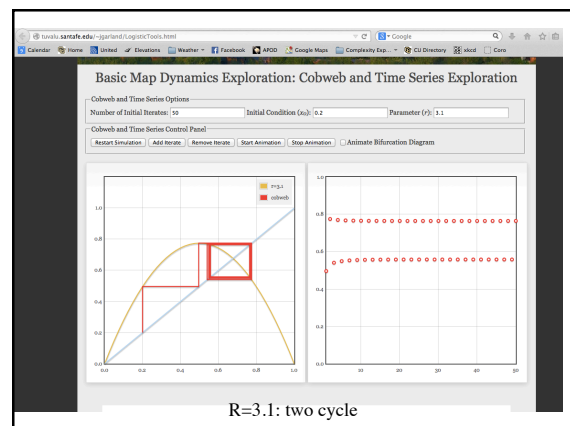
Image from Doug Ravenel's website at URochester

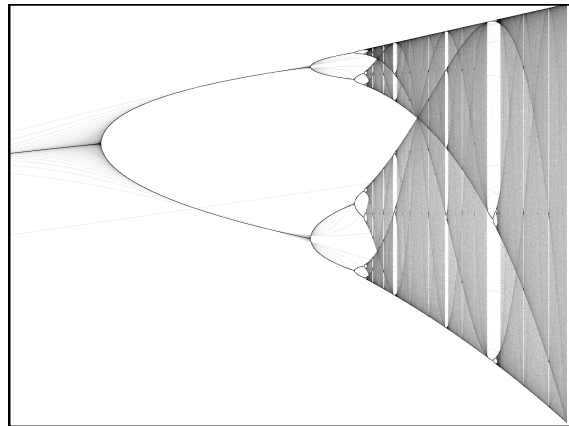
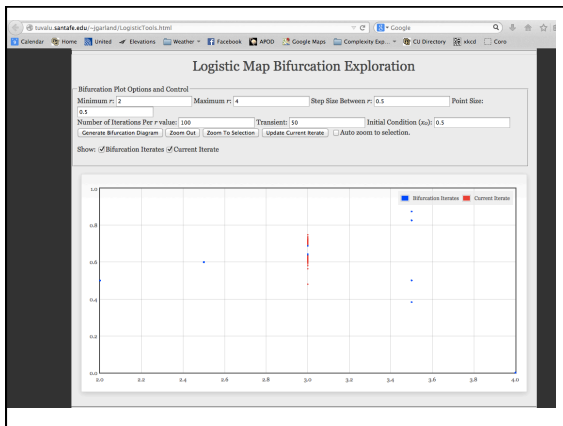


Bifurcations

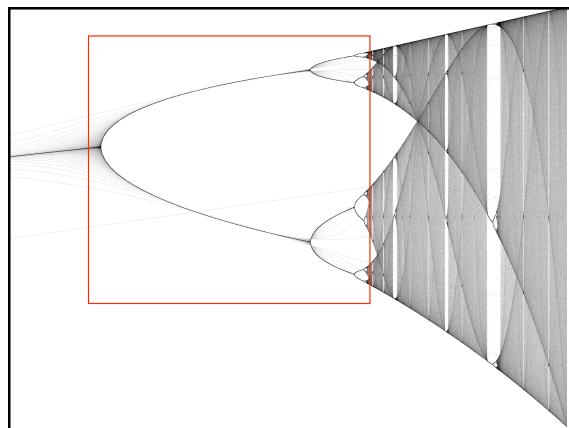
Qualitative changes in the dynamics caused by changes in parameters:

- Heart: pathology
- Eddy in creek: water level
- Olfactory bulb: smell
- Brain: blood chemicals
- Logistic map: R parameter...

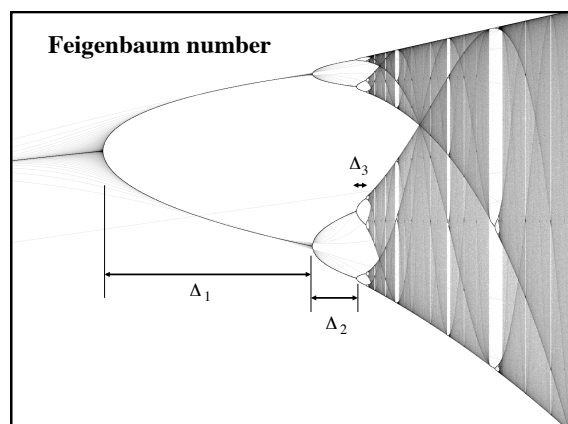




- chaos
- veils/bands: places where chaotic attractor is dense (UPOs)



- chaos
- veils/bands: places where chaotic attractor is dense (UPOs)
- *period-doubling cascade @ low R*

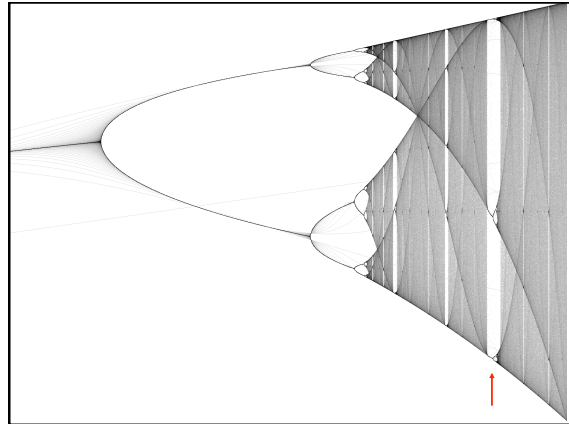


Universality!

Feigenbaum number and many other interesting chaotic/dynamical properties hold *for any 1D map with a quadratic maximum*.

Proof: renormalizations. See Strogatz §10.7

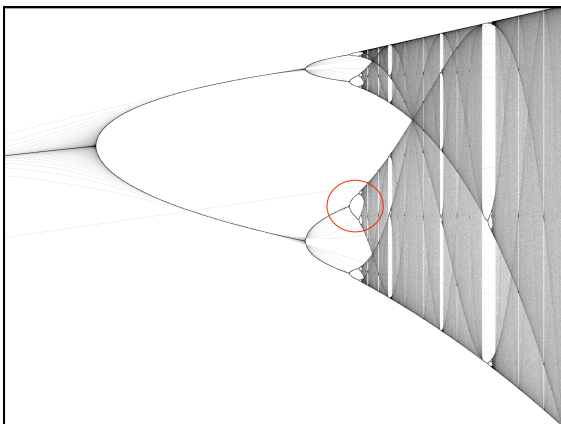
Don't take this too far, though...



- chaos
- veils/bands: places where chaotic attractor is dense (UPOs)
- period-doubling cascade @ low R
- *windows of order within the chaos, complete with their own period-doubling cascades (e.g., 3 to 6 to 12)*

There's something very special about 3...

- Sarkovskii (1964)
3, 5, 7, ..., 3×2 , 5×2 , ..., 3×2^2 , 5×2^2 , ..., 2^2 , 2, 1
- Yorke (1975)
- Metropolis *et al.* (1973)



- chaos
- veils/bands: places where chaotic attractor is dense (UPOs)
- period-doubling cascade @ low R
- windows of order within the chaos, complete with their own period-doubling cascades (e.g., 3 to 6 to 12)
- *small copies of object embedded in it (fractal)*

(lots of other interesting stuff, too — e.g., Misiurewicz points)

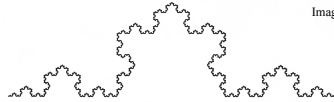


Fractals

- non-integer Hausdorff dimension
- self-similar

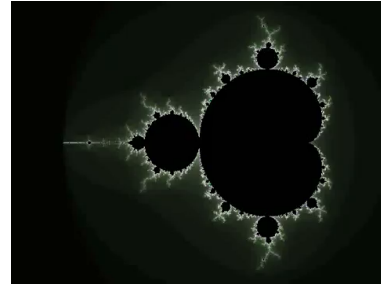


Images from Gleick



Canonical example: the Cantor set!

Another canonical example: the Mandelbrot set



www.youtube.com/watch?v=G_GBwuYu00s

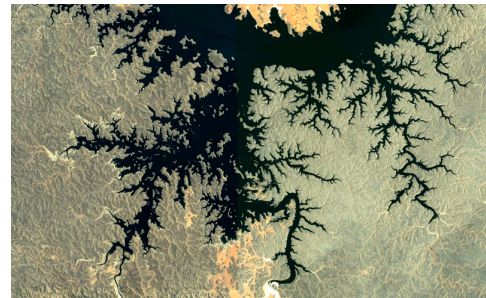
Fractals basins and basin boundaries



Newton's method
on $x^4 - 1 = 0$

From Strogatz

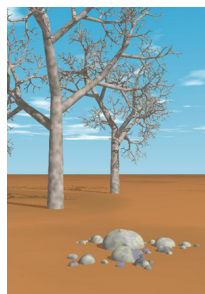
Fractals in the wild



paulbourke.net/fractals/googleearth/

See also: coastlines, trees, lungs, clouds, snowflakes ...

Fractals in computer graphics



Matthew Ward, WPI
davis.wpi.edu/~matt/courses/fractals/trees.html

Fractals and chaos...

The connection: *many (most)* chaotic systems have fractal state-space structure.

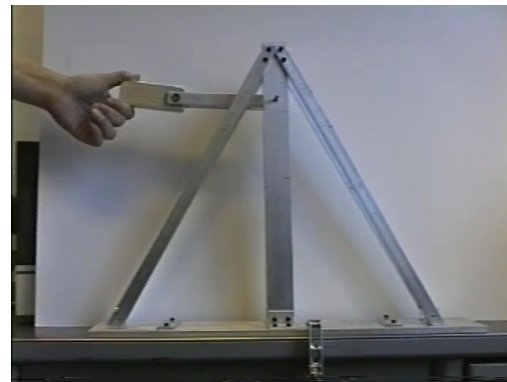
But **not** "all."

So far: mostly about *maps*.

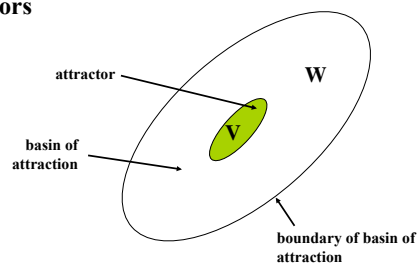
- discrete time systems:
 - time proceeds in clicks
 - “maps”
 - modeling tool: differ~~ence~~*ence* equation

Next up: *flows*

- continuous time systems:
 - time proceeds smoothly
 - “flows”
 - modeling tool: differ~~ential~~*ential* equations



Attractors



- Attractors exist only in dissipative systems!
- Dissipation \iff contraction of state space under the influence of the dynamics
- Can still have chaos if no dissipation...just not chaotic *attractors*



Courtesy of Allison Brown
Best Poster prize, Experimental Chaos Conference, 2012

Conditions for chaos in continuous-time systems

Necessary:

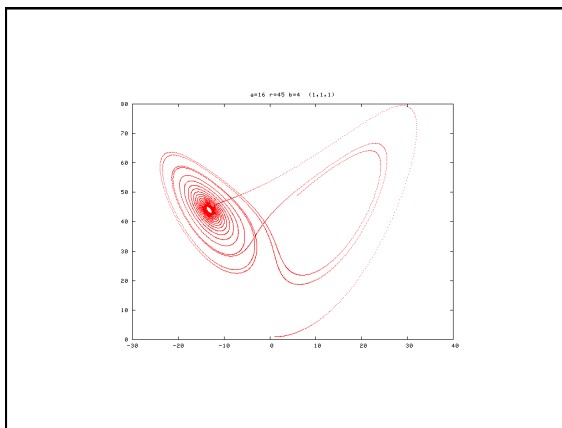
- Nonlinear
- At least three state-space dimensions (NB: only one needed in maps)

Necessary and sufficient:

- “Nonintegrable”
i.e., cannot be solved in closed form

Concepts: review

- State variable
- State space
- Initial condition
- Trajectory
- Attractor
- Basin of attraction
- Transient
- Fixed point (un/stable)
- Bifurcation
- Parameter



Deterministic Nonperiodic Flow¹

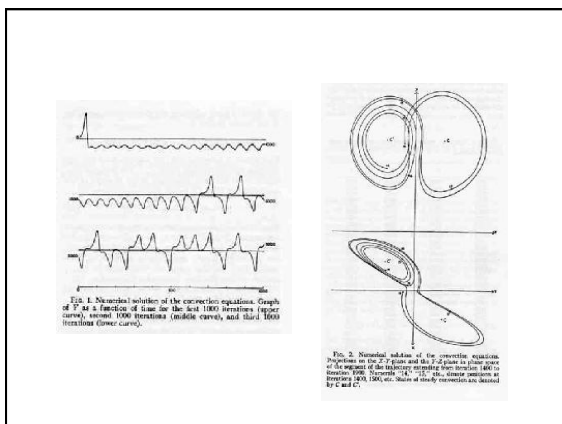
EDWARD N. LORENZ

Massachusetts Institute of Technology

(Manuscript received 18 November 1962, in revised form 7 January 1963)

ABSTRACT

Finite systems of deterministic ordinary nonlinear differential equations may be designed to represent forced dissipative hydrodynamic flow. Solutions of these equations can be identified with trajectories in phase space. For those systems with bounded solutions, it is found that nonperiodic solutions are ordinarily unstable with respect to small modifications, so that slightly differing initial states can evolve into considerably different states. Systems with bounded solutions are shown to possess bounded numerical solutions. A simple system representing cellular convection is solved numerically. All of the solutions are found to be unstable, and almost all of them are nonperiodic. The feasibility of very-long-range weather prediction is examined in the light of these results.

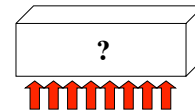


- Equations:

$$x' = a(y - x)$$

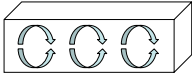
$$y' = rx - y - xz$$


$$z' = xy - bz$$

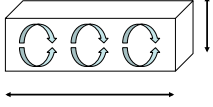


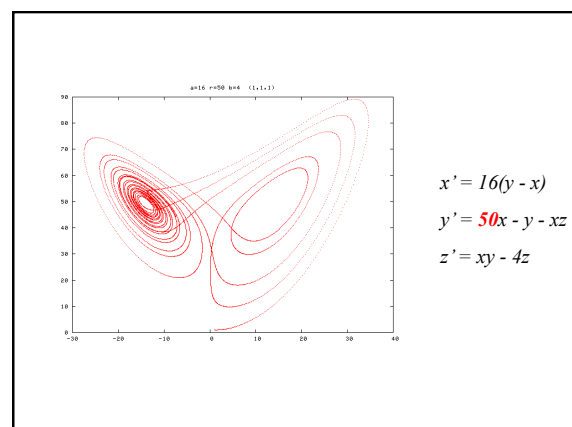
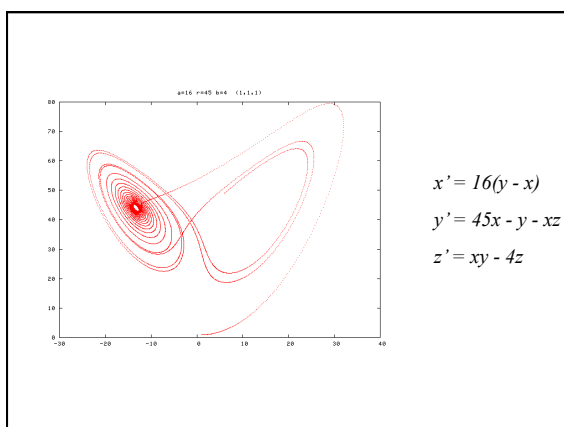
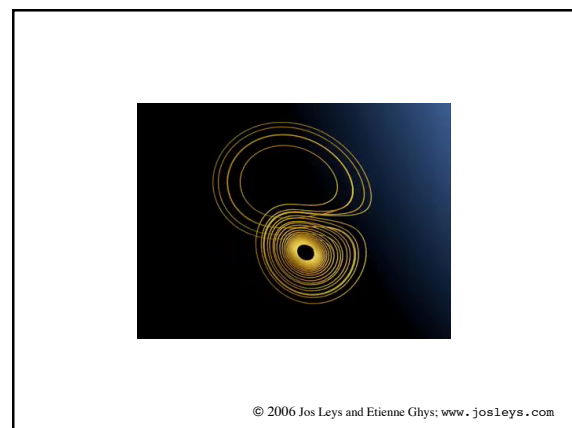
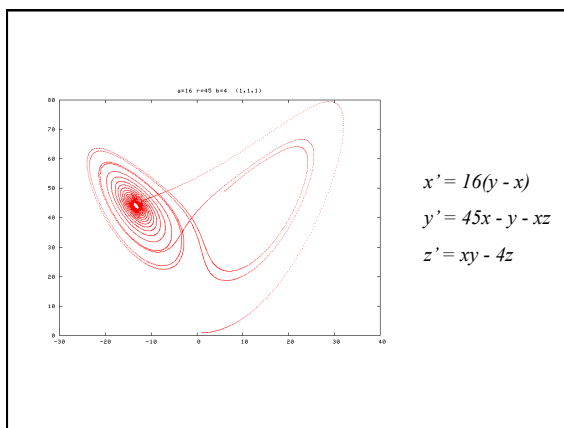
(first three terms of a Fourier expansion of the Navier-Stokes eqns)

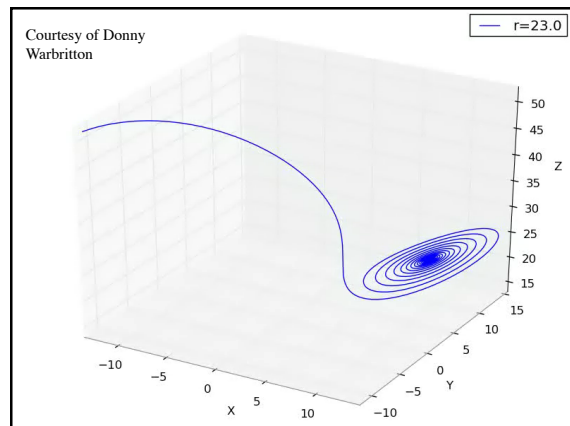
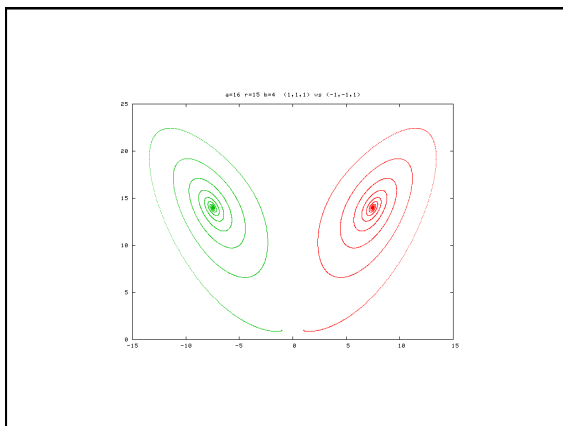
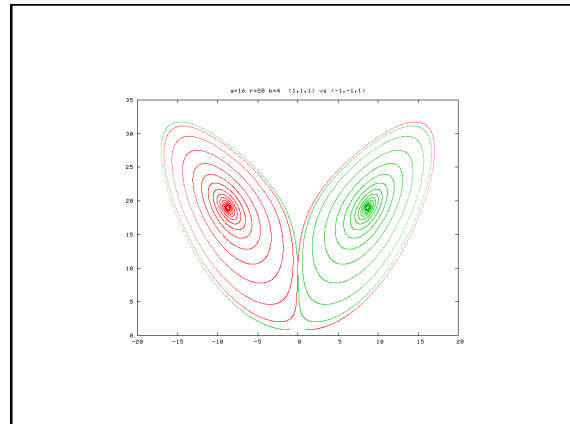
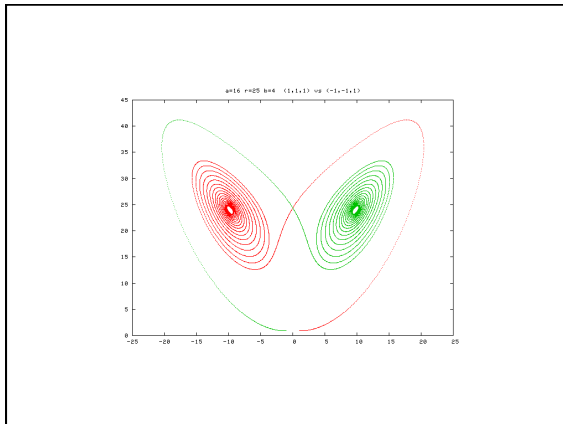
- State variables:
 - x convective intensity
 - y temperature
 - z deviation from linearity in the vertical convection profile



- Parameters:
 - a Prandtl number - fluids property
 - r Rayleigh number - related to ΔT 
 - b aspect ratio of the fluid sheet







Before we leave Lorenz...

Deterministic Nonperiodic Flow¹

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The feasibility of very-long-range weather prediction is examined in the light of these results.

Attractors

Four types:

- fixed points
- limit cycles (*aka* periodic orbits)
- quasiperiodic orbits
- chaotic attractors

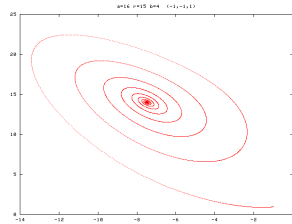
A nonlinear system can have any number of attractors, of all types, sprinkled around its state space

Their basins of attraction (plus the basin boundaries) *partition* the state space

And there's no way, *a priori*, to know where they are, how many there are, what types, etc.

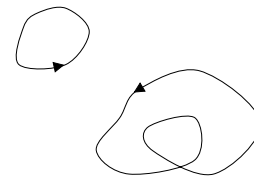
Attractors

- Fixed point



Attractors

- Limit cycle

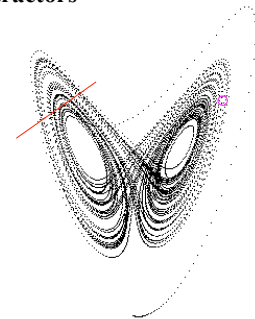


Attractors

- Quasi-periodic orbit...

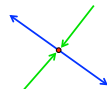
“Strange” or chaotic attractors

- often fractal
- covered densely by trajectories
- exponential divergence of neighboring trajectories...



Lyapunov exponents

- nonlinear analogs of eigenvalues: one λ for each dimension

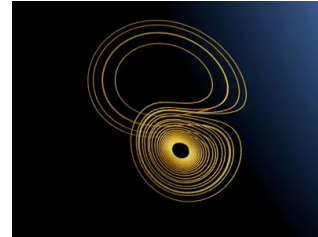
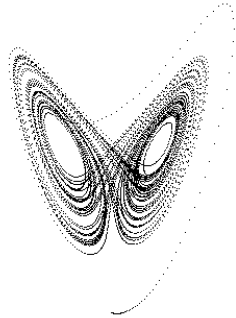


Lyapunov exponents: summary

- nonlinear analogs of eigenvalues: one λ for each dimension
- negative λ_i compress state space; positive λ_i stretch it
- $\sum \lambda_i < 0$ for dissipative systems
- λ_i are same for all ICs in one basin
- long-term average in definition; biggest one (λ_1) dominates as $t \rightarrow \infty$
- positive λ_1 is a signature of chaos

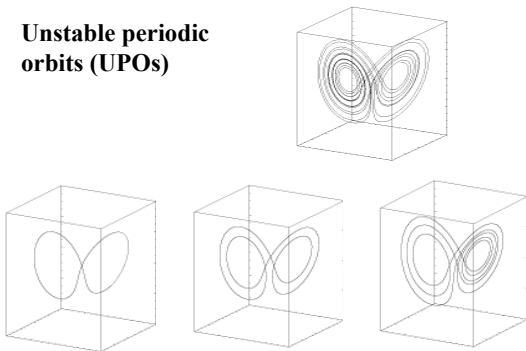
“Strange” or chaotic attractors:

- exponential divergence of neighboring trajectories
- *often* fractal
- covered densely by trajectories
- contain an infinite number of “unstable periodic orbits”...



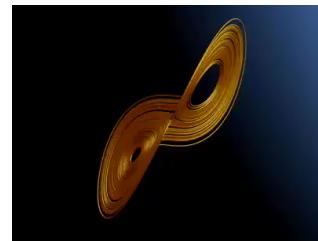
© 2006 Jos Leys and Etienne Ghys; www.josleys.com

Unstable periodic orbits (UPOs)

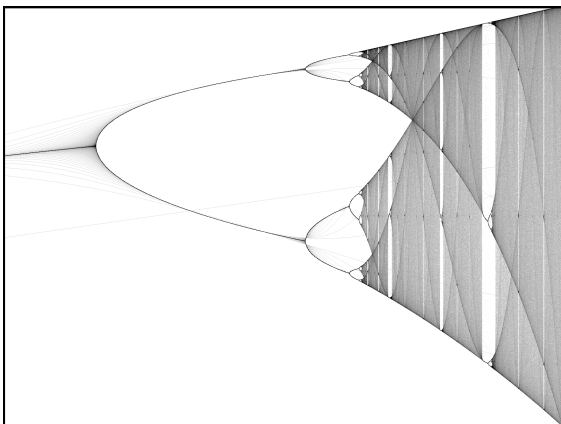


Bradley/Mantilla, *Chaos* 12:596

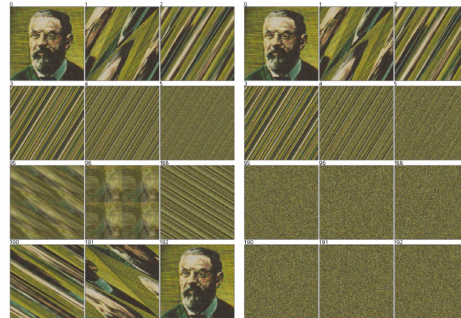
“Attractor bones”...



© 2006 Jos Leys and Etienne Ghys; www.josleys.com



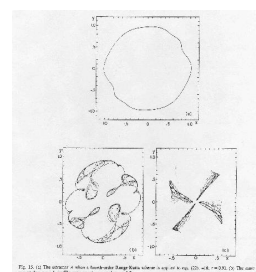
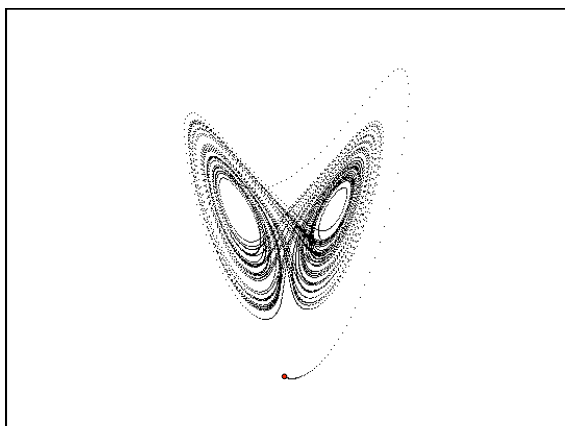
Poincare recurrence



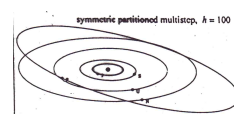
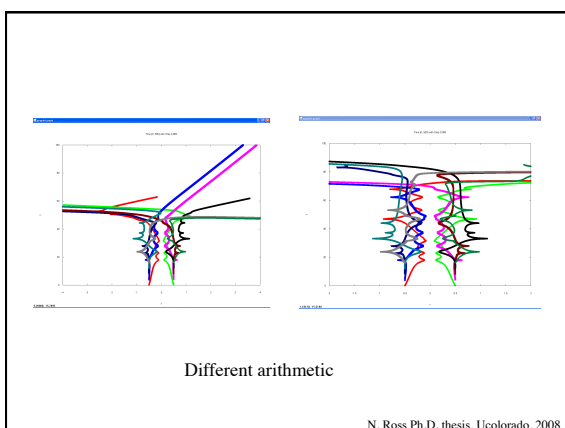
Representative point chosen at corner of pixel

Representative point chosen at a random place in pixel

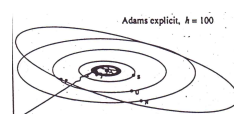
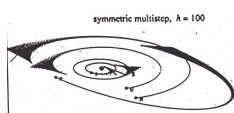
Crutchfield *et al.* *Chaos* 25:46 and
<http://www.mpi.kg-dresden.mpg.de/mpi-doc/kantgruppe/wiki/projects/Recurrence.html>



Different timestep

Lorenz, *Physica D* 35:229

Different solver algorithm...



Moral: numerical methods can run amok in “interesting” ways...

- can cause distortions, bifurcations, etc.
- and these look a lot like *real, physical* dynamics...
- source: algorithms, arithmetic system, timestep, etc.
- Q: what could you do to diagnose whether your results included spurious numerical dynamics?

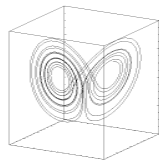
Moral: numerical methods can run amok in “interesting” ways...

- can cause distortions, bifurcations, etc.
- and these look a lot like *real, physical* dynamics...
- source: algorithms, arithmetic system, timestep, etc.
- Q: what could you do to diagnose whether your results included spurious numerical dynamics?
 - *change the timestep*
 - *change the method*
 - *change the arithmetic*

But beware
machine ϵ ...

So ODE solvers make mistakes.

...and chaotic systems are sensitively dependent on initial conditions....



...??!?

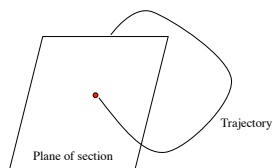
Shadowing lemma

Every* noise-added trajectory on a chaotic attractor is *shadowed* by a true trajectory.

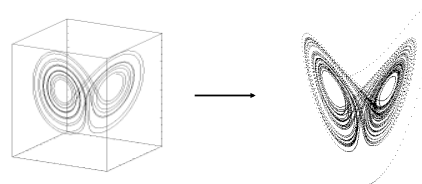
Important: this is for *state* noise, not *parameter* noise.

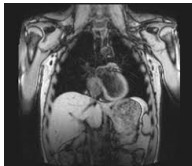
(*) Caveat: not if the noise bumps the trajectory out of the basin

Section



Not the same thing as a projection!



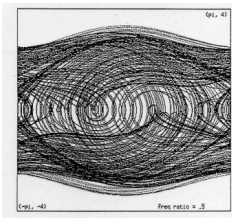
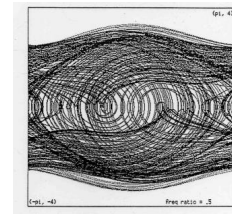


mri.radiology.uiowa.edu

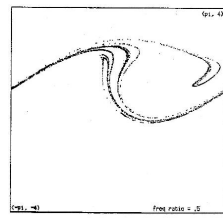


prixray.com

The driven damped pendulum



trajectory

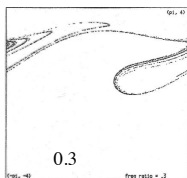


Poincaré section

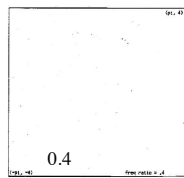
Time-slice sections of periodic orbits: some thought experiments

- pendulum rotating @ 1 Hz and strobe @ 1 Hz?
- pendulum rotating @ 1 Hz and strobe @ 2 Hz?
- pendulum rotating @ 1 Hz and strobe @ 3 Hz?
- pendulum rotating @ 1 Hz and strobe @ 1/2 Hz?
- pendulum rotating @ 1 Hz and strobe @ π Hz? (or some other irrational)

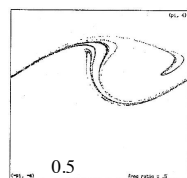
What bifurcations look like on a Poincaré section



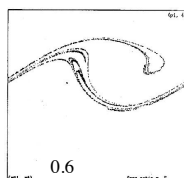
0.3



0.4

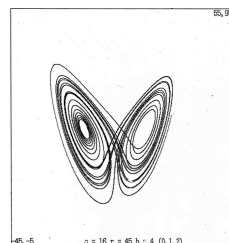


0.5

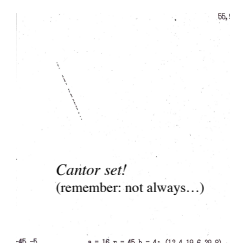


0.6

Spatial sections



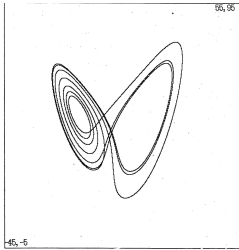
a = 18 r = 45 b = 4 (0.1, 0)



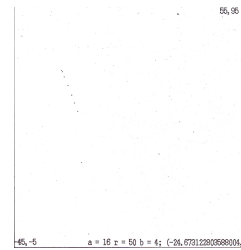
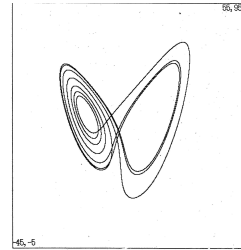
Cantor set!
(remember: not always...)

a = 18 r = 45 b = 4; (12.4, 18.6, 28.8)

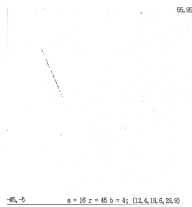
What about a section of a UPO?



?



Aside: finding UPOs

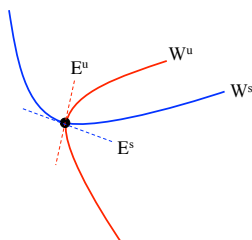


- Section
- Look for close returns
- Cluster
- Average
- See Gunaratne, So papers

Computing sections

- If you're slicing in state space: use the "inside-outside" function
- If you're slicing in *time*: use modulo on the timestamp
- See Parker & Chua for more details

λ_i and the un/stable manifolds (W^u and W^s)

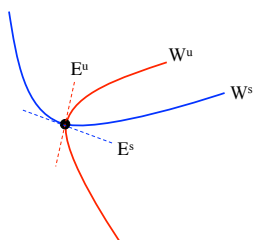


Aside: finding those un/stable manifolds

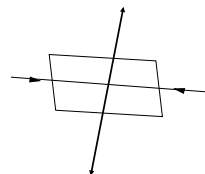
- Linearize the system
- Find the eigenvectors E^s and E^u
- Take a step along E^s ; run time forwards
- Take a step along E^u ; run time backwards
- See Osinga & Krauskopf paper for more details

Note: saddles are not the only possible landscape geometry around fixed points (they're just the most interesting ones!)

These λ_i & manifolds play a critical role in the control of chaos...



Local-linear control* of a hyperbolic point




* e.g., via pole placement

Lyapunov exponents, revisited:

- n -dim system has n λ_i ; $\sum \lambda_i < 0$ for dissipative systems
- λ_i are same for all ICs in one basin
- negative λ_i compress state space along *stable manifolds*
- positive λ_i stretch it along *unstable manifolds*
- biggest one (λ_1) dominates as $t \rightarrow \infty$
- positive λ_1 is a signature of chaos
- calculating them:
 - From equations: eigenvalues of the variational matrix (see variational system notes on CSC15446 course webpage, which you can access from Liz's homepage.)
 - From data: various creative algorithms...

Calculating λ (& other invariants) from data

- The bible: H. Kantz & T. Schreiber, *Nonlinear Time Series Analysis*
- Associated software: TISEAN
www.mpi-pks-dresden.mpg.de/~tisean
- A recent review article: EB & H. Kantz, "Nonlinear Time Series Analysis Revisited," *CHAOS* **25**:097610 (2015)




TISEAN
Nonlinear Time Series Analysis

Rainer Hegger
Holger Kantz
Thomas Schreiber

[Go to Version 3.0.1 \(released March 2007\)](#)

[Go to Version 2.1 \(released December 2000\)](#)



TISEAN 3.0.1

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- [Embedding and Poincaré sections](#)
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- [Unsupported](#)

Generating time series

A few routines are provided to generate test data from simple equations. Since there are powerful packages (for ex. Helena Nusse and Jim Yorke) that can generate chaotic data, we have only included a minimal selection here.

Lyapunov exponents are an important means of quantification for unstable systems. They are however difficult to estimate from a time series. Unless low dimensional, high quality data is at hand, one should not attempt to calculate the full spectrum. Try to compute the maximal exponent first. The two implementations differ slightly. While `lyap_k` implements the formula by Kantz, `lyap_r` uses that by Rosenstein et al. which differs only in the definition of the neighbourhoods. We recommend to use the former version, `lyap_k`.

The estimation of Lyapunov exponents is also discussed in the [introduction](#) paper. A recent addition is a program to compute finite time exponents which are not invariant but contain additional information.

```
Maximal exponent lyap_k lyap_r
Lyapunov spectrum lyap_spec
```

Description of the program: `lyap_k`

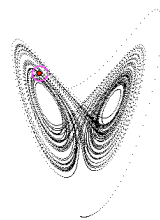
The program estimates the largest Lyapunov exponent of a given scalar data set using the algorithm of Kantz.

Usage:

`lyap_k [Options]`

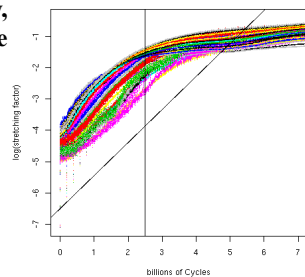
Everything not being a valid option will be interpreted as a potential datafile name. Given no datafile at all, means read stdin. Also - means stdin

Kantz's algorithm:



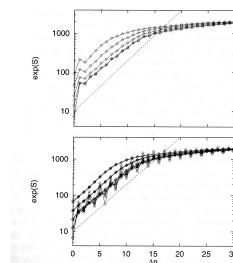
1. Choose point K •
2. Look at the points around it (ϵ neighborhood)
3. Measure how far they are from K
4. Average those distances
5. Watch how that average grows with time (Δn)
6. Take the log, normalize over time $\rightarrow S(\Delta n)$
7. Repeat for lots of points K and average the $S(\Delta n)$

If you're lucky,
things look like
this.



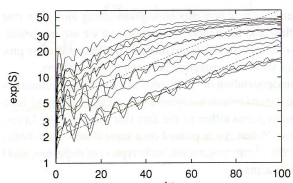
- The slope of the scaling region — iff one exists — is the λ .

Or this:



This is fig 5.3 in Kantz & Schreiber

If you're not lucky:

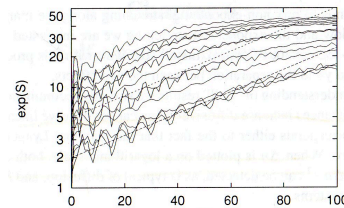


- The slope of the scaling region — iff one exists — is the λ .



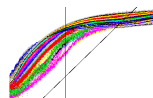
This is fig 5.4 in Kantz & Schreiber

What do you think those oscillations
might be?



Calculating λ (& other invariants) from data

Be careful! TISEAN has lots of knobs and its results are incredibly sensitive to their values!



Different colors on that plot from before = different settings for one of those knobs

Option	Description	Default
-l#	number of data to be used	whole file
-x#	number of lines to be ignored	0
-c#	column to be read	1
-M#	maximal embedding dimension to use	2
-m#	minimal embedding dimension to use	2
-d#	delay to use	1
-r#	minimal length scale to search neighbors	(data interval)/1000
-R#	maximal length scale to search neighbors	(data interval)/100
-l#	number of length scales to use	5
-n#	number of reference points to use	all
-s#	number of iterations in time	50
-t#	'thinner window'	0
-o#	output file name	without file name: 'datafile'.lyap (or stdin.lyap if the data were read from stdin)
-V#	verbosity level 0: only panic messages 1: add input/output messages 2: add statistics for each iteration	3
-h	show these options	none

Description of the Output:

For each embedding dimension and each length scale the file contains a block of data consisting of 3 columns

1. The number of the iteration
2. The logarithm of the stretching factor (the slope is the Lyapunov exponent if it is a straight line)
3. The number of points for which a neighborhood with enough points was found

Calculating λ (& other invariants) from data

• Be careful! TISEAN has lots of knobs and its results are incredibly sensitive to their values!

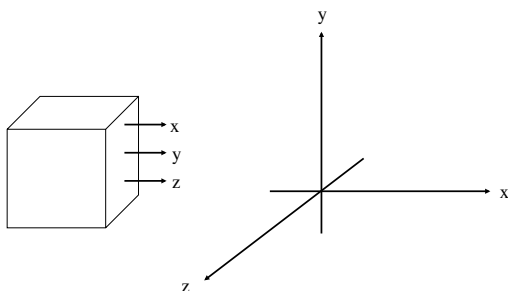
• Use your dynamics knowledge to understand & use those knobs intelligently

• Look at the results plots. For example, do not blindly fit a regression line to something that has no scaling region (this is a good idea in general, of course)

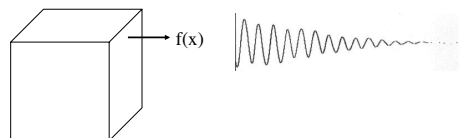
Fractal dimension:

- Capacity
- Box counting
- Correlation (d2 in TISEAN)
- Lots of others:
 - Kth nearest neighbor
 - Similarity
 - Information
 - Lyapunov
 - ...
- See Chapter 6 and §11.3 of Kantz & Schreiber

We've been assuming that we can measure all the state variables...

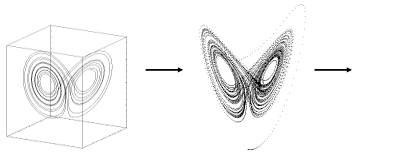


But often you can't.



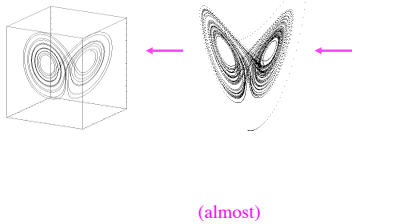
Rarely do you even *know* what they are. And even if you did, you might not be able to measure them directly...

How to undo a projection?

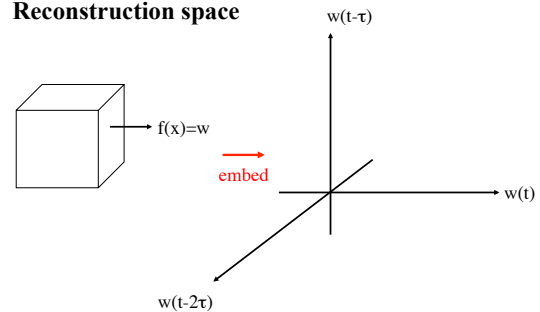


Delay-coordinate embedding

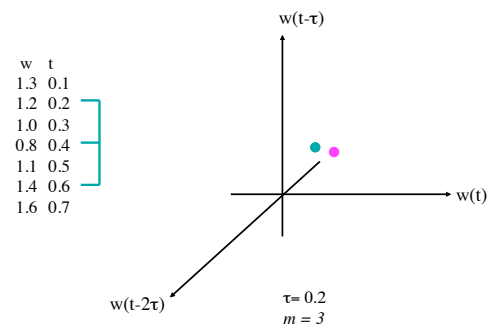
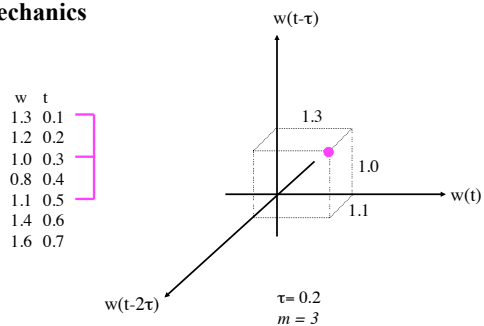
“reinflate” that squashed data to get a *topologically identical* copy of the original thing.



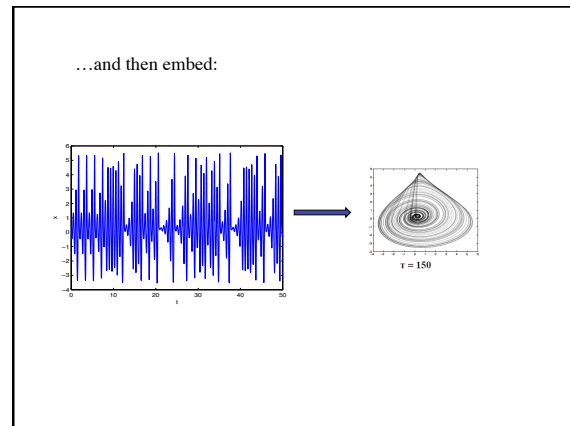
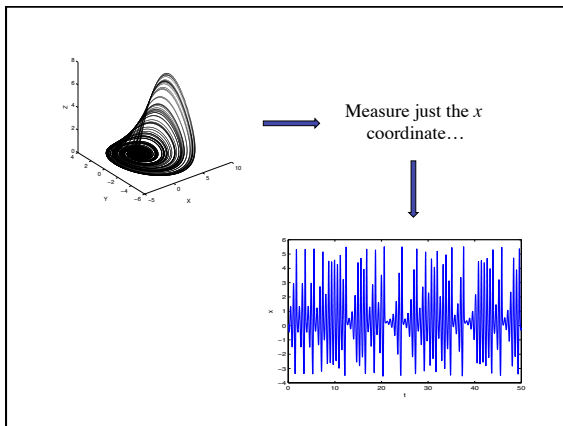
Reconstruction space



Mechanics

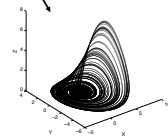
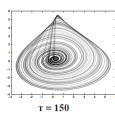


TISEAN's `delay` command does this



Takens* theorem

For the **right τ** and **enough dimensions**, the embedded dynamics are diffeomorphic to (and thus have same topology as) the original state-space dynamics.



* Whitney, Mane, ...

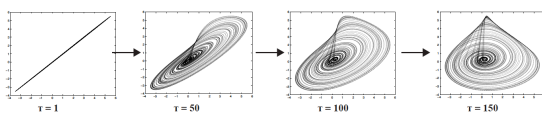
Note: the measured quantity must be a smooth, generic function of at least one state variable, and must be uniformly sampled in time.

Diffeomorphic: mapping from the one to the other is differentiable and has a differentiable inverse.

What that means:

- *qualitatively* the same shape (topology)
- have same dynamical invariants (e.g., λ)

Choosing τ :



TISEAN contains tools that help you do this (e.g., `mutual`)

Choosing m

$m > 2d$: **sufficient** to ensure no crossings in reconstruction space (Takens et al.)...

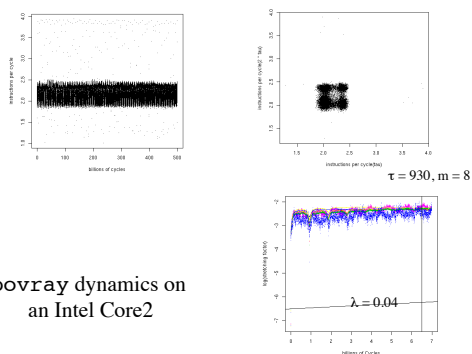
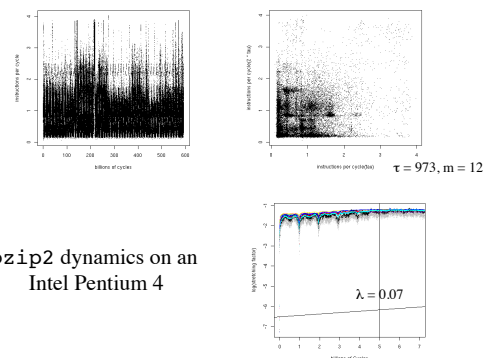
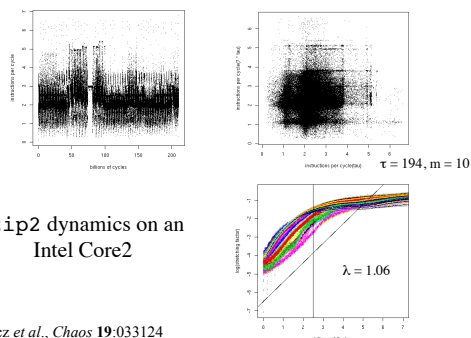
...but that may be overkill, and you rarely know d anyway.

"Embedology" paper: $m > 2 d_{\text{box}}$
(box-counting dimension)

TISEAN contains tools that help you do this (e.g., `false_nearest`)

NLTSA* of computer performance dynamics

* nonlinear time-series analysis



Caveat: need enough data...

If Δt is not uniform

~~Theorem (Takens): for $\tau > 0$ and $m > 2d$, reconstructed trajectory is diffeomorphic to the true trajectory~~

~~Conditions: evenly sampled in time, smooth generic measurement function~~

Interspike interval embedding

idea: lots of systems generate spikes — hearts, nerves, etc.

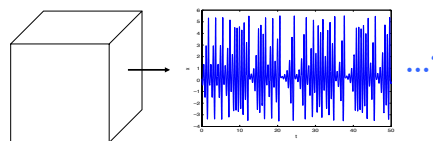
if you assume that the spikes are the result of an integrate-and-fire system, then the Δt has a one-to-one correspondence to some state variable's *integrated* value...

in which case the embedding theorems still hold.

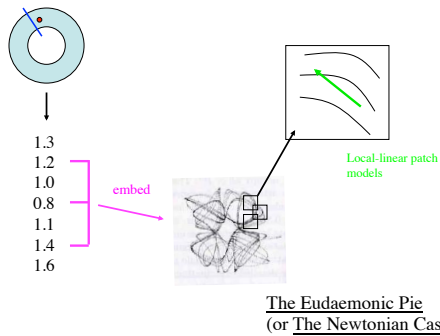
(with the Δt s as state variables)

Sauer *Chaos* 5:127

Prediction

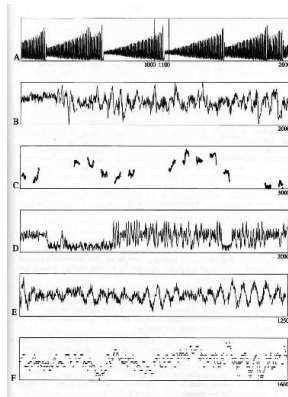


Predicting the path of a roulette ball...



The Santa Fe competition

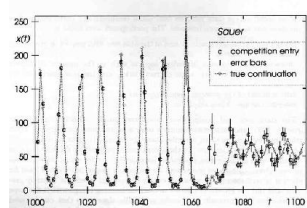
- Weigend & Gershenfeld, 1992
- put a bunch of data sets up on an ftp server
- and invited all comers to predict their future
- chronicled in *Time Series Prediction: Forecasting the Future and Understanding the Past*, Santa Fe Institute, 1993 (from which the images on the following half-dozen slides were reproduced)



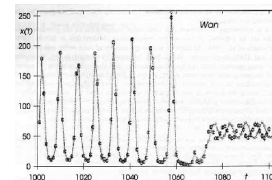
The Santa Fe competition: data

- Laboratory laser
- Medical data (sleep apnea)
- Currency rate exchange
- RK4 on some chaotic ODE
- Intensity of some star
- A Bach fugue

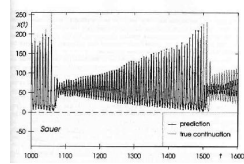
Embedding + patch models: (Sauer)



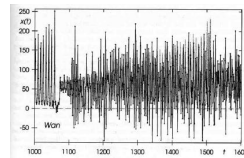
Neural net: (Wan)



Further out:

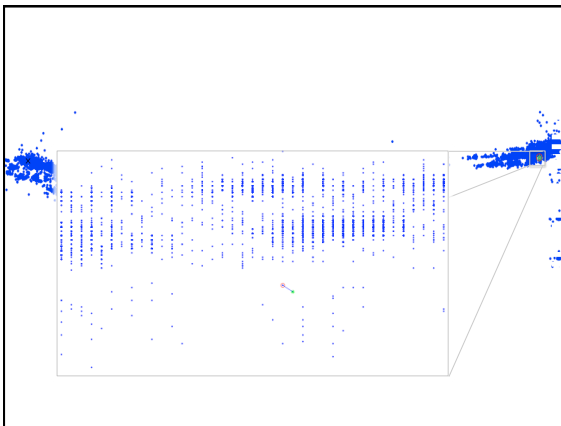
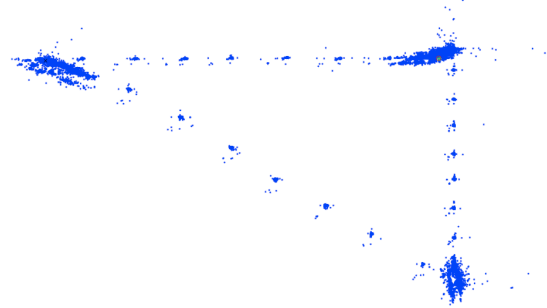


Sauer

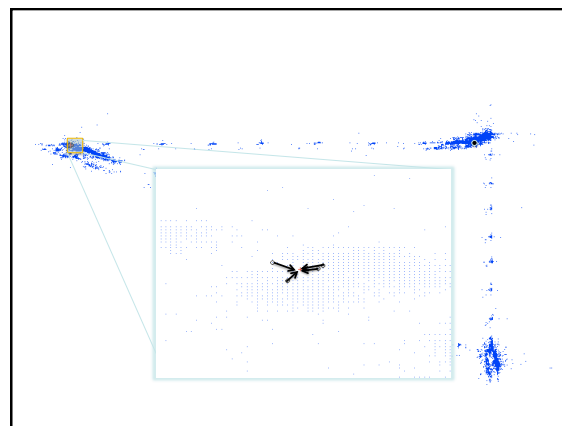
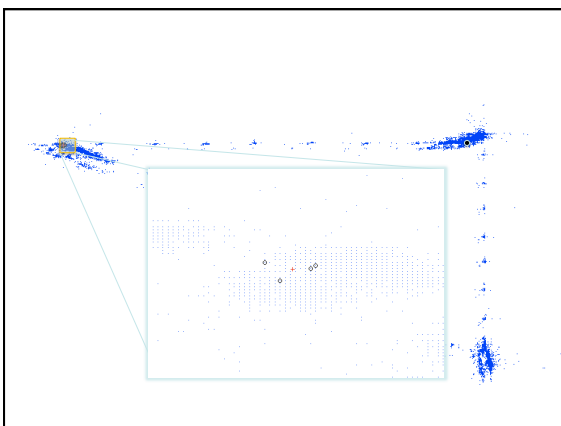
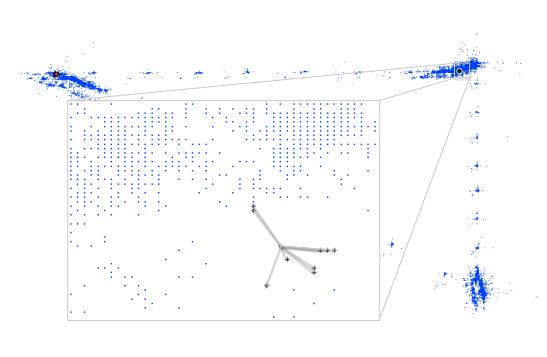


Wan

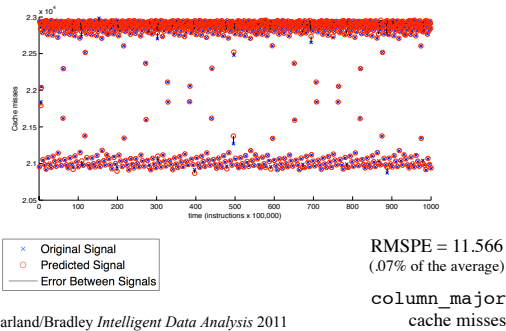
An even simpler prediction method:
Lorenz's method of analogues



A k -nearest neighbor modification of LMA



Using kLMA to predict computer dynamics

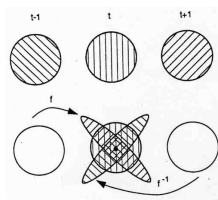


Noise...

Linear filtering: a bad idea if the system is chaotic

Nonlinear alternatives:

- use the stable and unstable manifold structure on a chaotic attractor...

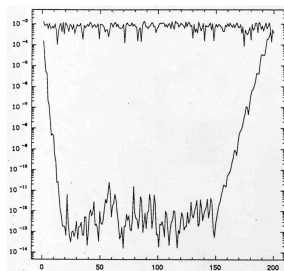


Farmer & Sidorowich, in *Evolution, Learning and Cognition*, World Scientific, 1983

Idea:

- If you have a model of the system, you can simulate what happens to each point in forward *and backward* time
- If your system has transverse stable and unstable manifolds, that does useful things to the noise balls
- Since all three versions of that data should be identical at the middle time, can average them
- ➡ noise reduction!
- Works best if manifolds are perpendicular, but requires only transversality

Results:



Farmer & Sidorowich, in *Evolution, Learning and Cognition*, World Scientific, 1983

Noise...

Linear filtering: a bad idea if the system is chaotic

Nonlinear alternatives:

- use the stable and unstable manifold geometry on a chaotic attractor
- what about using the *topology* of the attractor?

Computational Topology

Why: this is the fundamental mathematics of shape. complements geometry.

What: compute topological properties from finite data

How:

- introduce resolution parameter
- count components and holes at different resolutions
- deduce topology from patterns therein

V. Robins Ph.D. thesis, UColorado, 1999



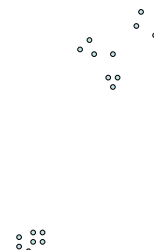
Connectedness: definitions

• how many “lumps” in a data set:

• ϵ -connectedness (after Cantor)

• ϵ -connected components

• ϵ -isolated points:

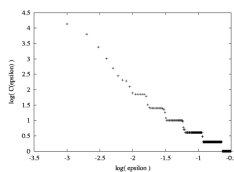


Connectedness: examples

If the data points are samples of a disconnected fractal like this:



The number of connected components looks like this:



(note obvious tie-in to fractal dimension...)

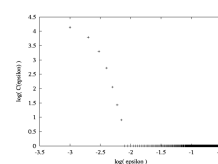
Robins et al., *Physica D* 139:276, *Nonlinearity* 11:913

Connectedness: examples

If the data points are samples of a connected set like this:



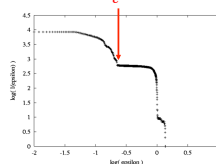
The number of connected components looks like this:



Robins et al., *Physica D* 139:276, *Nonlinearity* 11:913

Connectedness and filtering

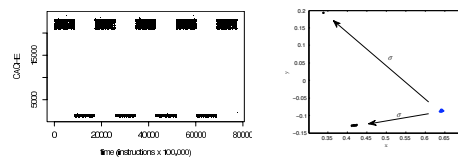
The effect of noise is to add isolated points to the set and a shoulder to the $C(\epsilon)$ curve:



So if you know that the object is connected — like the attractor of a flow — you can reasonably assume that any isolated points are noisy, and remove them by pruning with $\epsilon = \epsilon^*$

Robins et al., *Intelligent Data Analysis* 8:505, *Chaos* 14:305

Continuity and filtering



Idea:

- deterministic, differentiable dynamics (maps & flows) are *continuous*

Conjecture:

- if the image of a connected set is not connected, more than one dynamics is at work

Approach:

- track connectedness over time

Applications:

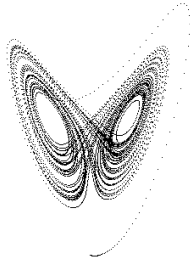
- pulling apart interleaved dynamics, removing noise...

Alexander et al., *CHAOS*, 2012

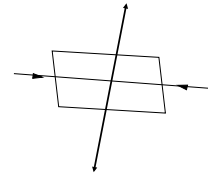
Chaos and control

key concepts:

- dense attractor coverage
- exponential trajectory separation
- un/stable manifold structure
- local-linear control

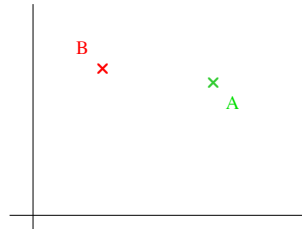


Recall: local-linear control of a saddle point works successfully in a region whose geometry is defined by the λ_i and the W^s/W^u , together with the sensor & actuator capabilities...



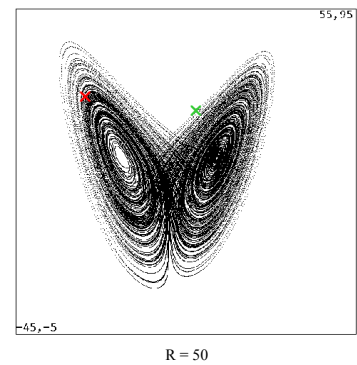
Control:

getting from A to B, minimizing some cost functional

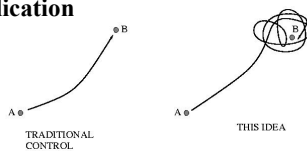


Lorenz System:

denseness, reachability, and control



Denseness & reachability in a real engineering application

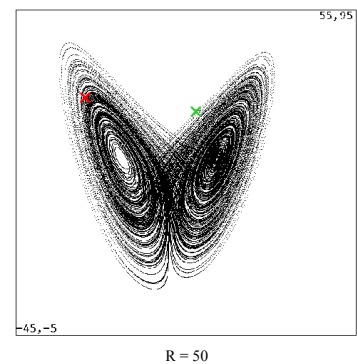


- can control position/volume/density of attractor — *within limits*
- possibly not reachable any other way
- not for time-critical applications (that “eventually”)

Using Chaos to Broaden the Capture Range of a Phase-Locked Loop

Elizabeth Bradley, Member, IEEE

Now what?

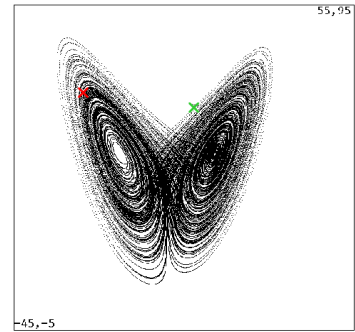


OGY control

- dense attractor coverage \rightarrow reachability
- un/stable manifold structure + UPO denseness + local-linear control \rightarrow controllability

Ott et al., PRL 64:1196

Use local-linear control, designed using the eigenvalues and eigenvectors at that point \times to balance a chaotic system on a UPO passing through that point.



But you're relying on denseness to get you into the controllable region, and that may take a while...

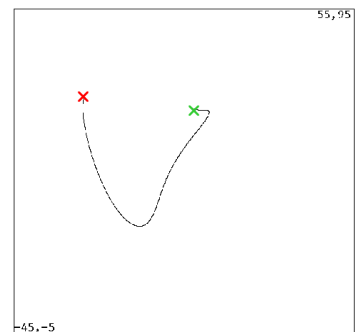
- dense attractor coverage \rightarrow reachability
- un/stable manifold structure + UPO denseness + local-linear control \rightarrow controllability
- exploit sensitive dependence, too???
 \Rightarrow "targeting"

Lorenz System:

SDOIC-based targeting

OGY & co. have been used in *tons* of systems; see Shimbro review paper.

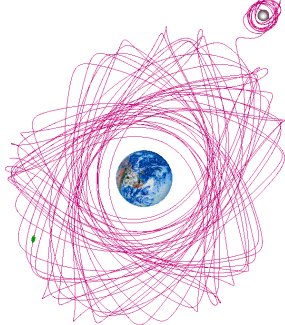
Alfred Hubler has done a lot of cool stuff in this area as well.



Four R switches; 240X faster

Bradley, Cybernetics & Systems 26:299

Program in
Applied
Mathematics



Erik Bollt

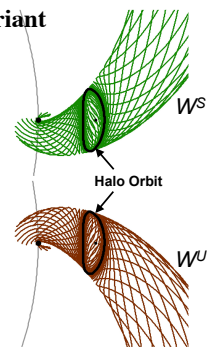
University of Colorado at Boulder
Boulder CO 80309-0526
(303) 492-4668

Other cool ways to use invariant manifolds

Want to get a spacecraft onto a "halo orbit," which is a UPO of the dynamics.

Unstable Periodic Orbits (UPOs) have invariant manifolds:

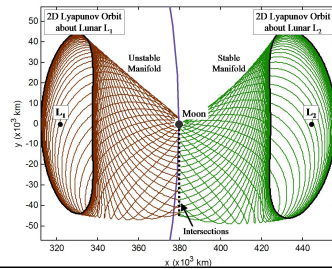
- Stable Invariant Manifold (W^s)
 – The set of all trajectories a particle could use to arrive onto the UPO.
- Unstable Invariant Manifold (W^u)
 – The set of all trajectories a particle could take after a small perturbation from the UPO.



Jeff Parker, PhD thesis, UColorado 2008

Low-energy (cheap) orbit transfers

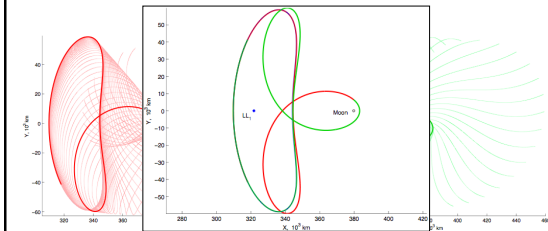
- Depart along W_{L1}^U & arrive on W_{L2}^S



Jeff Parker, PhD thesis, UColorado 2008

Homoclinic orbits - The best case

- If a trajectory in Stable and Unstable intersect ("homoclinic connection")



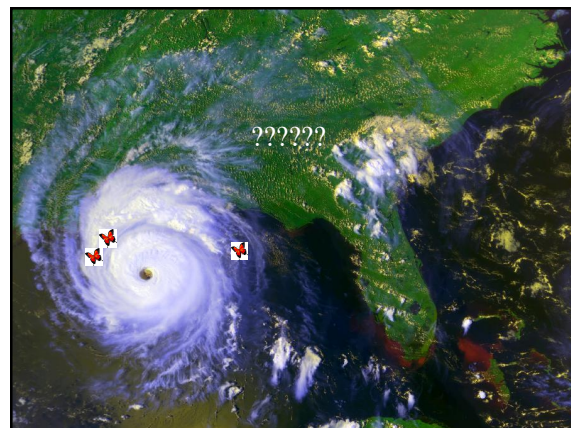
Unstable Manifold of an LL_1 Lyapunov Orbit

Stable Manifold of an LL_1 Lyapunov Orbit

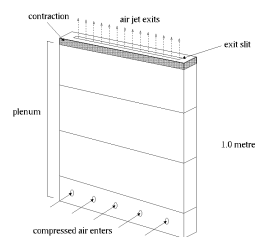
Jeff Parker, PhD thesis, UColorado 2008

Can we do any of that in spatially extended systems?

(i.e. harness the butterfly effect, exploit un/stable manifold geometry?)

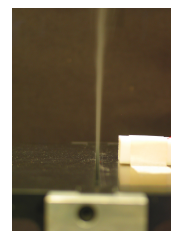


A 2D jet

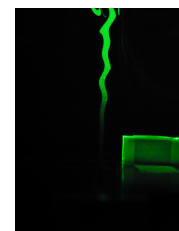


Peacock et al., *Exp. Fluids* 37:22

End view



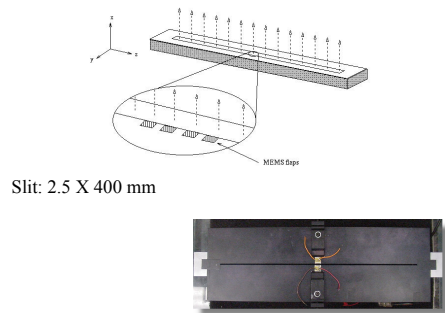
room lighting



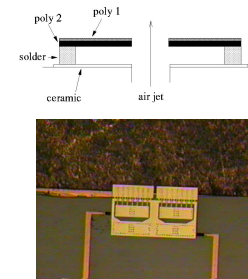
stop-action laser "slice"

aerosolized canola oil

Forcing the jet flow



MEMS actuators

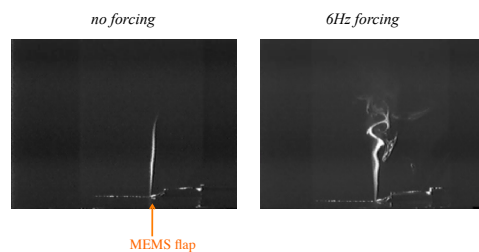


Video : overhead view at 2Hz, 10V

Area of individual flap is 1.0 x 0.25mm

Ma et al., IEEE Trans. Adv. Packaging 26:268

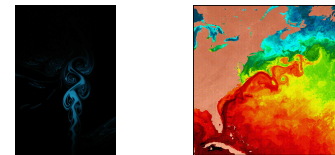
The Butterfly effect in action...



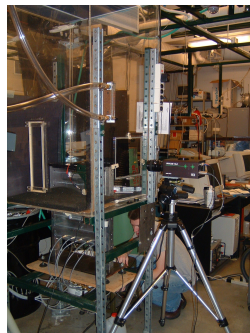
Forcing generates coherent structures that enhance entrainment and mixing

Peacock et al., Exp. Fluids 37:22

Does this have anything to do with reality?



Measurement & isolation:



Communication and chaos:

- Two coupled Lorenz systems will synchronize
- Robust to a small amount of noise
- Use this to transmit & receive information

$$\begin{aligned}x' &= a(y-x) \\ y' &= rx - y - xz \\ z' &= xy - bz\end{aligned}$$

$$\begin{aligned}x' &= a(y-x) \\ y' &= r(x+ex) - y - xz \\ z' &= xy - bz\end{aligned}$$

- Chaotic carrier wave, so hard to intercept or jam

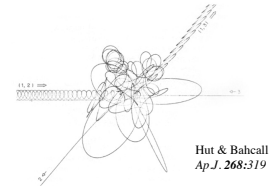
Pecora & Carroll Phys. Rev. Lett 64:821

Another interesting application: chaos in the solar system

- orbits of Pluto, Mars
- Kirkwood gaps
- rotation of Hyperion & other satellites
- ...

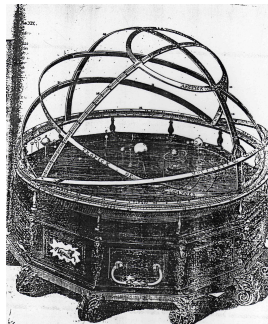
Solar system stability:

- recall: two-body problem not chaotic
- but three (or more) can be...

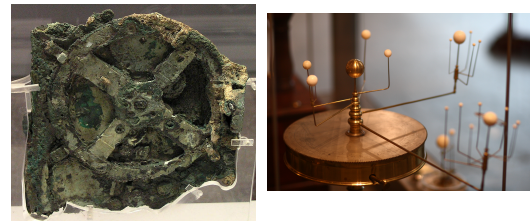


Exploring that issue, circa 1880:

An *orrery*, which is a *mechanical computer* whose gear ratios and circular platters simulate the orbits of the planets



Exploring that issue before the digital computer age...

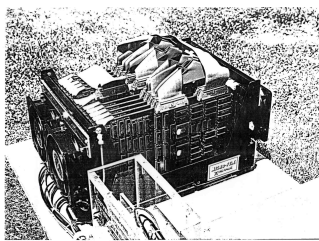


An *orrery*, which is a *mechanical computer* whose gear ratios and circular platters simulate the orbits of the planets

Exploring that issue, circa 1980:

- write the n -body equations for the solar system
- solve them using symplectic ODE solvers on a special-purpose computer

The *digital orrery*
(Wisdom & Sussman)



Numerical Evidence That the Motion of Pluto Is Chaotic

GERALD JAY SUSSMAN AND JACK WISDOM

The Digital Orrery has been used to perform an integration of the motions of the outer planets for 846 million years. This integration indicates that the long-term motion of the planet Pluto is chaotic. Nearby trajectories diverge exponentially with an e-folding time of only about 20 million years.

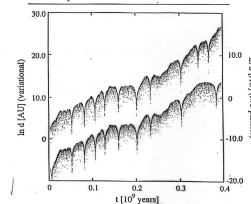


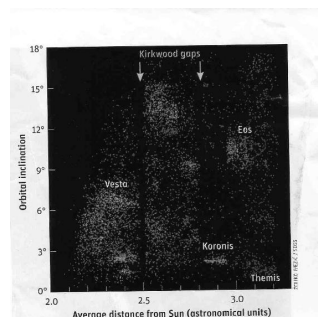
Figure 5: The exponential divergence of nearby trajectories is indicated by the average linear growth of the logarithm of the distance between a family of lines. In the upper trace we see the growth of the vertical distance between a reference trajectory. In the lower trace we see two trajectories diverge with time. The distance between any 10^6 AU and the reference trajectory is approximately 10^6 AU. The vertical spread of trajectories we see are approximately 10^6 AU. The vertical spread of trajectories we see are approximately 10^6 AU. The vertical spread of trajectories we see are approximately 10^6 AU.

Science 241:433

Should we worry?

- No.

Kirkwood gaps:



From *Sky & Telescope*

Chaos and the Kirkwood gaps

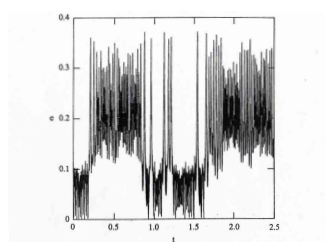


FIGURE 3. Eccentricity of a typical chaotic trajectory over a longer time interval, the time is now measured in millions of years. Bursts of high eccentricity behavior are interspersed with intervals of irregular low eccentricity behavior, broken by occasional spikes.

Wisdom, *Nuclear Phys. B* 2:391

Evidence in favor of the conjecture:

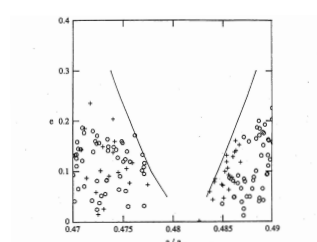


FIGURE 3. Comparison of the actual distribution of asteroids with the outer boundaries of the chaotic zone. There is both a chaotic region and quasi-periodic region in the gap, but trajectories of both types are planet crossing.

Wisdom, *Nuclear Phys. B* 2:391

Chaotic tumbling of satellites:

Voyager and Galileo saw this...

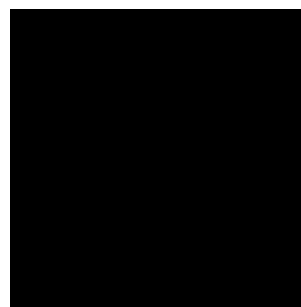


144 August 2000 Sky & Telescope

From *Sky & Telescope*

Ap. J. 97:570
Ap. J. 98:1855

...so did Cassini:



www.nasa.gov/mission_pages/cassini/multimedia/pia06243.html

Chaotic tumbling of satellites:



This happens for **all** satellites at some point in their history, unless they are perfectly spherical and in perfectly circular orbits (pf: KAM theorem; see Wisdom paper on syllabus)

Some of them are still tumbling chaotically because of their geometry, but most (like the earth and its moon) have settled down into tidal equilibria

More chaos in the solar system:

- obliquity of Mars (Touma & Wisdom, *Science* 259:1294)



www.solarviews.com

- etc.

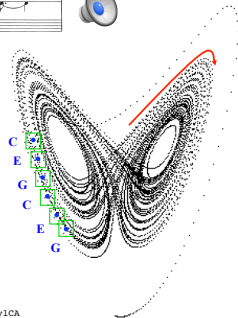
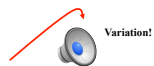
Musical Variations from a Chaotic Mapping



Dabby Chaos 6:95

Pitch sequence:
C, E, G, C, E, G, C, E...

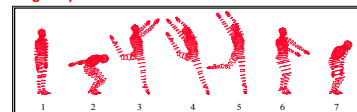
C symbol dynamics



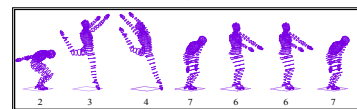
Also fun: <http://www.youtube.com/watch?v=B2XtE9TyICA>

Chaotic variations on movement sequences

original piece



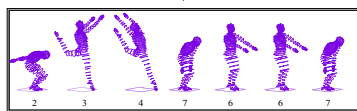
chaotic mapping



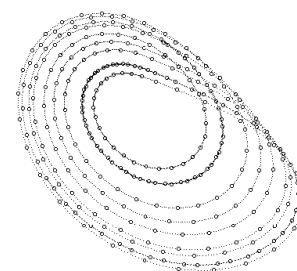
chaotic variation

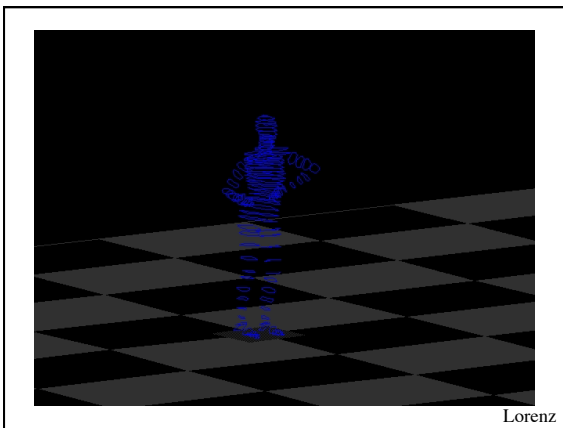
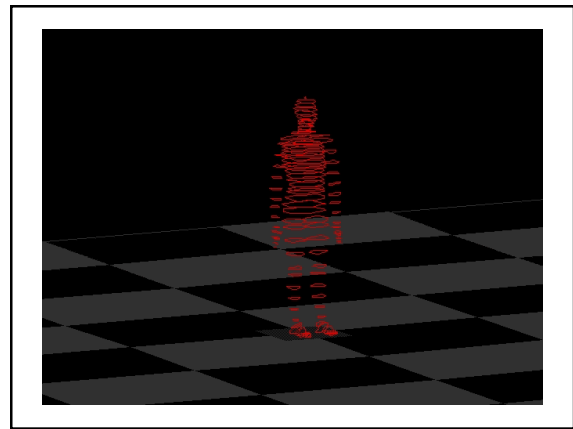
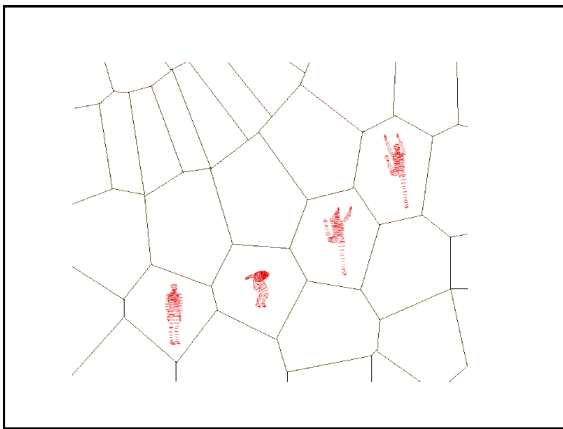
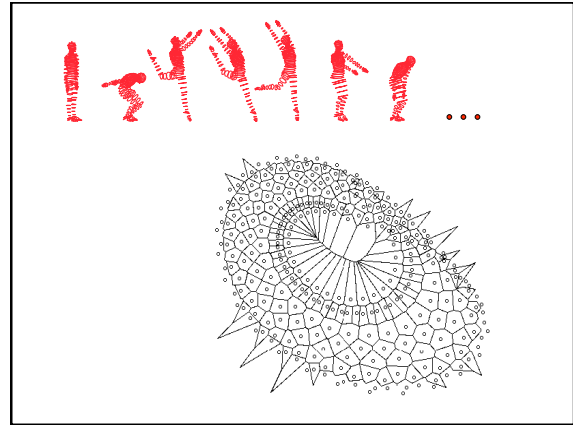
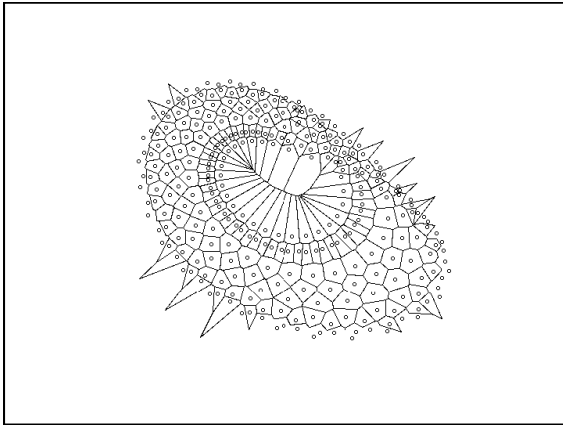
Bradley & Stuart, *Chaos* 8:800

original piece

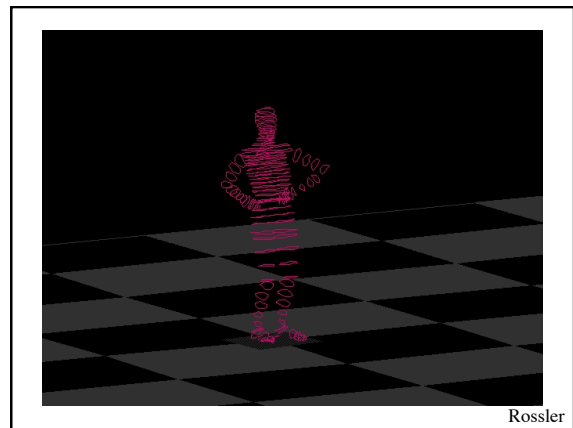


chaotic variation

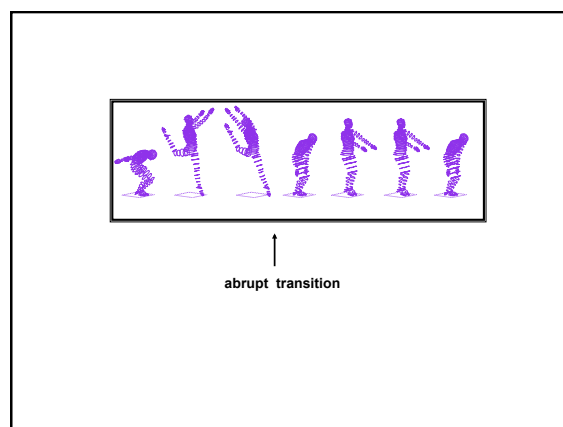
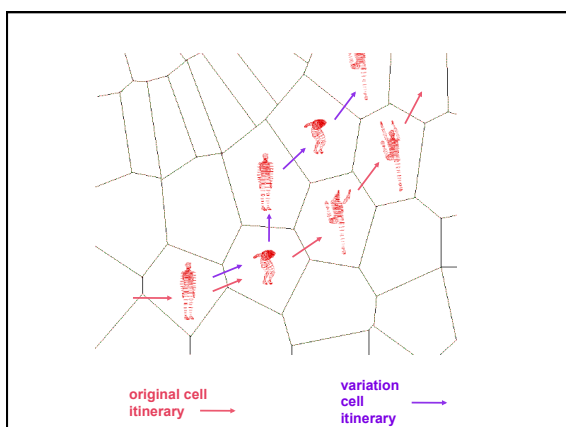
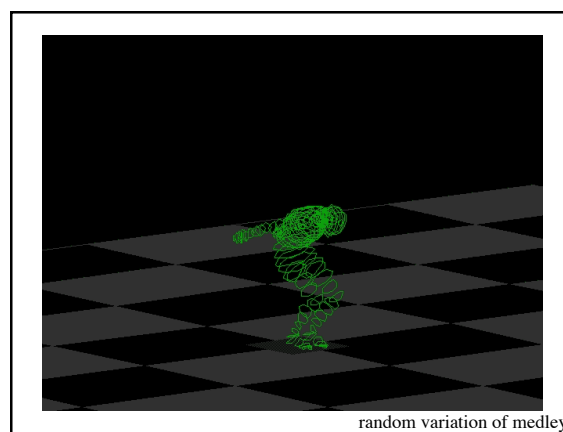
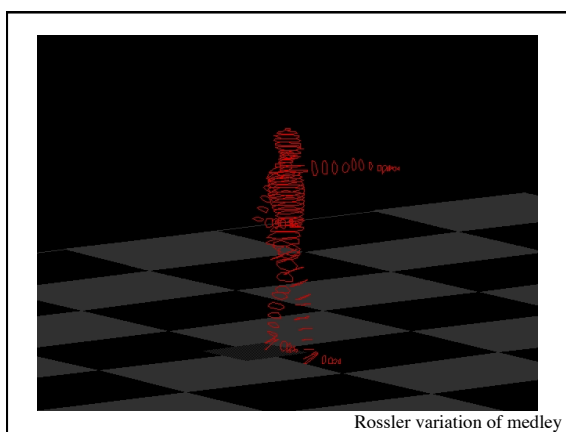
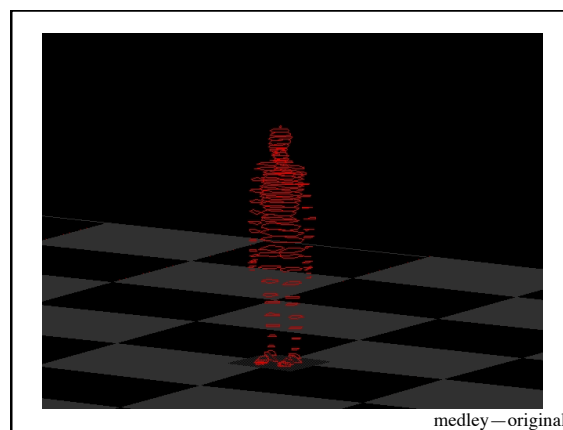
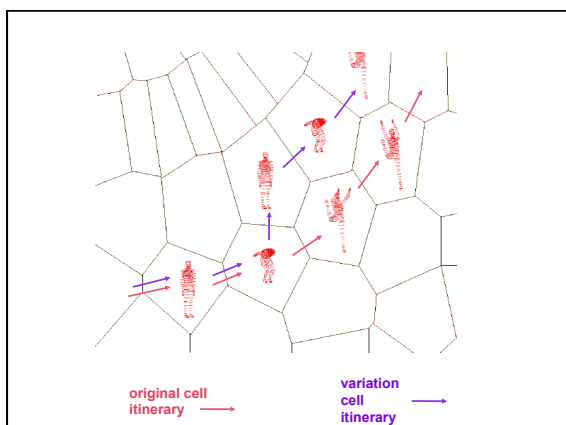


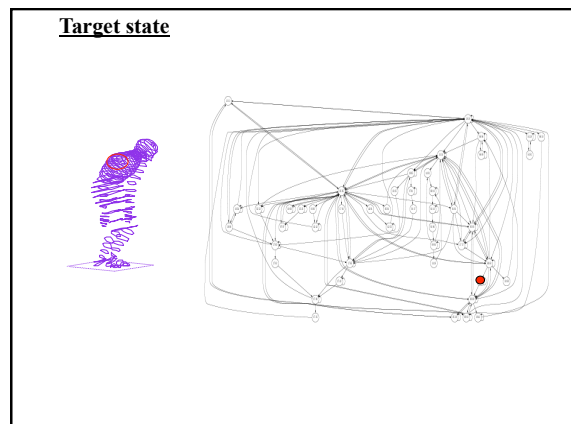
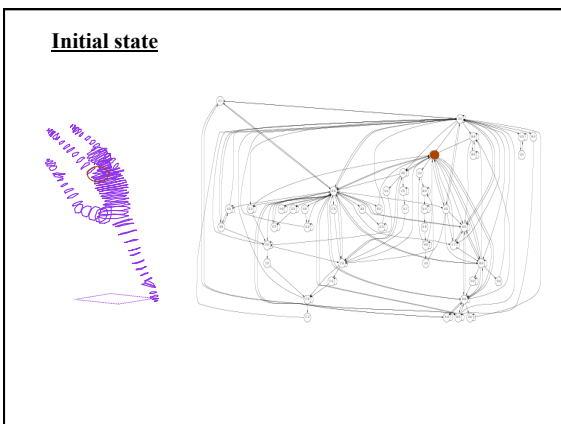
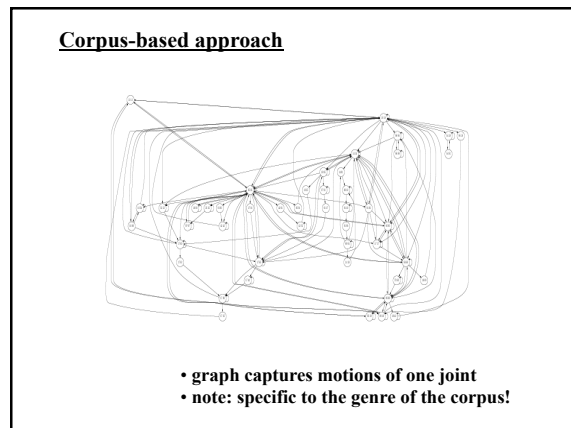
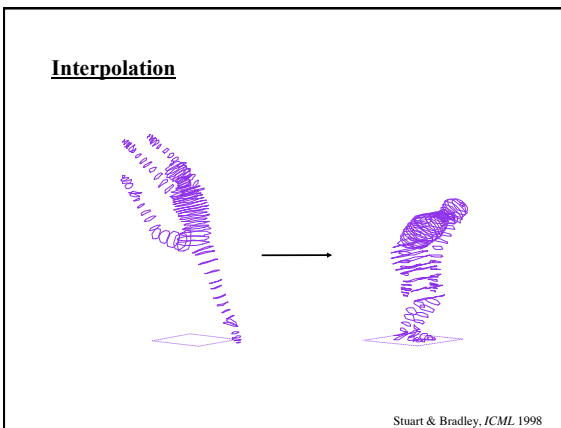
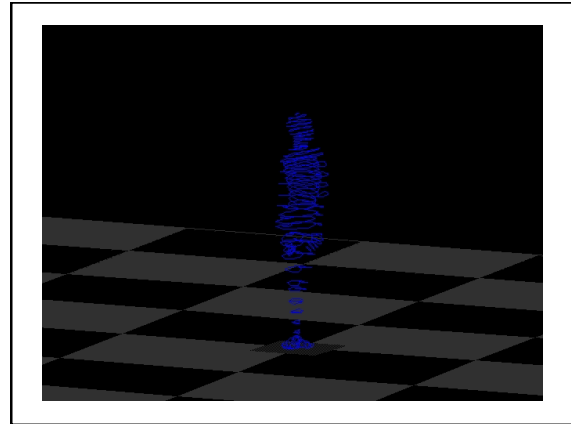
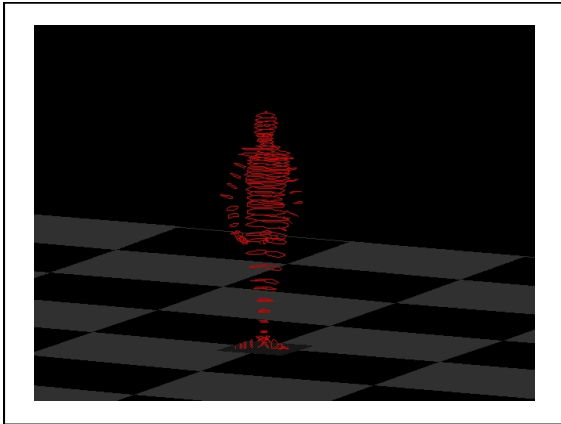


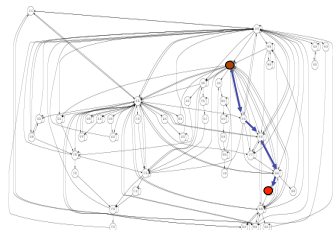
Lorenz



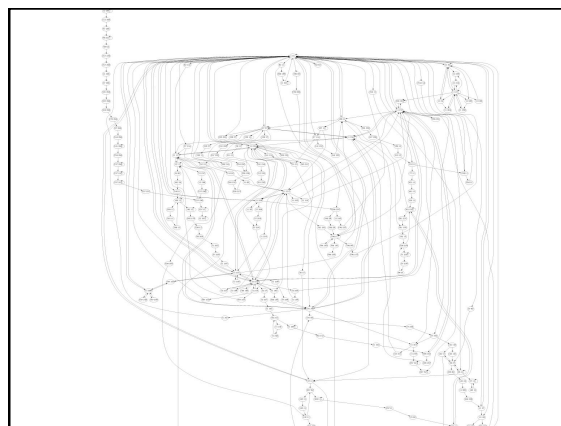
Rossler



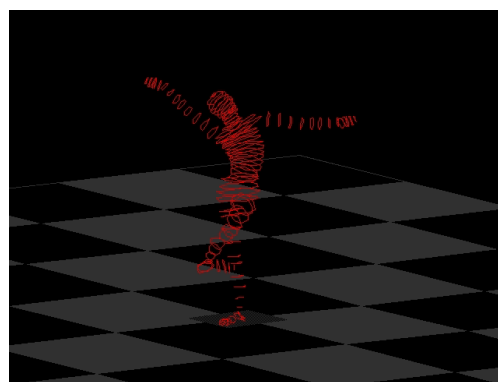


Graph search

...for 44 joints in parallel!



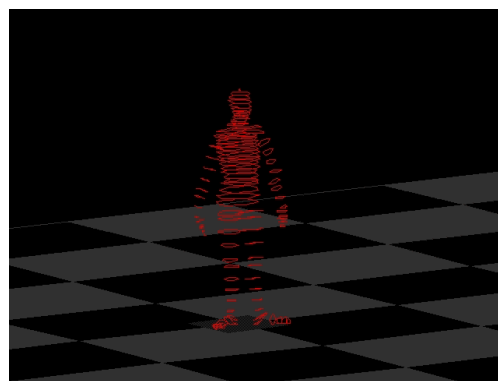
initial



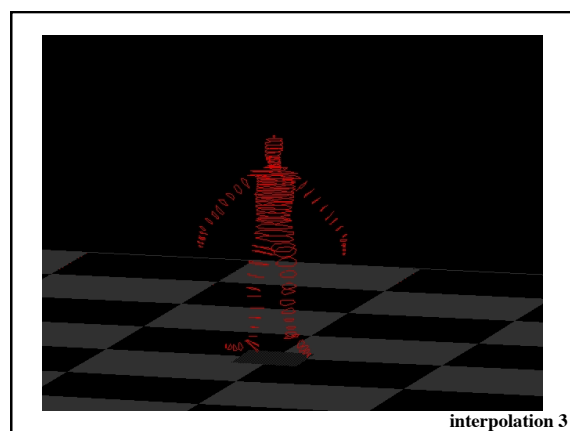
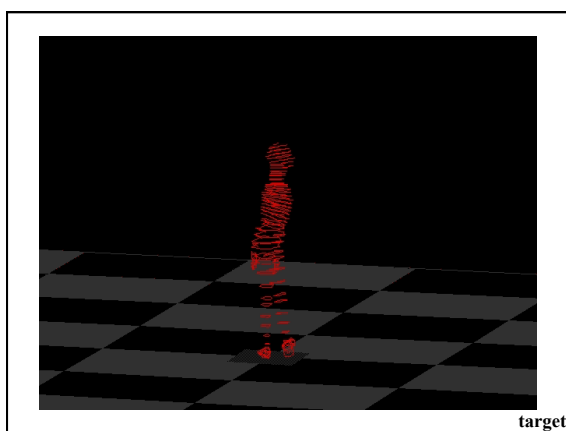
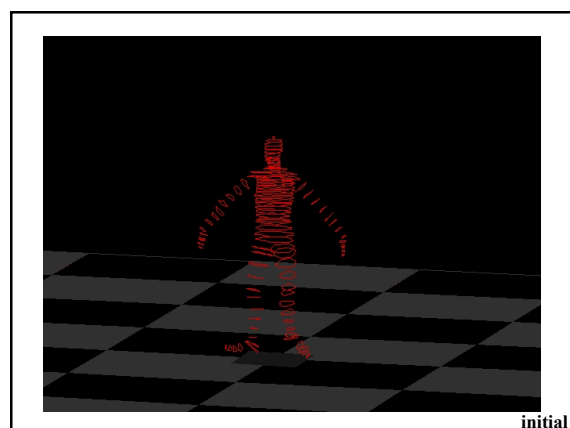
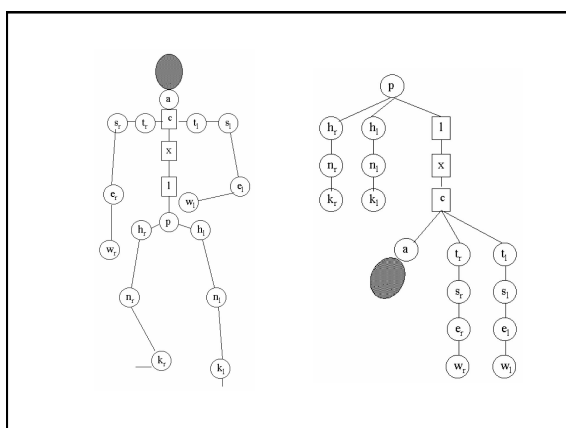
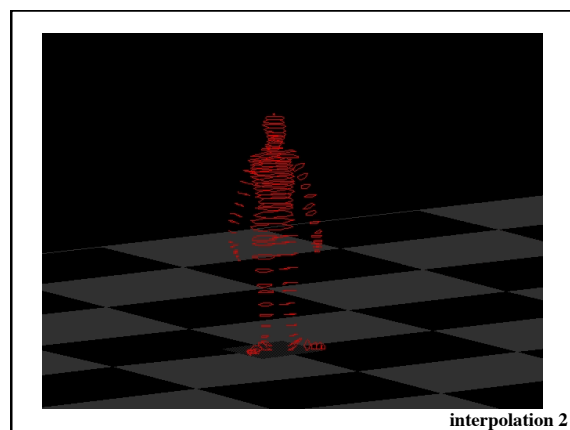
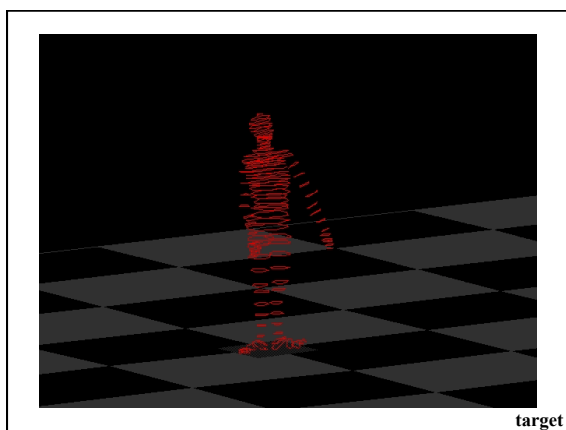
target

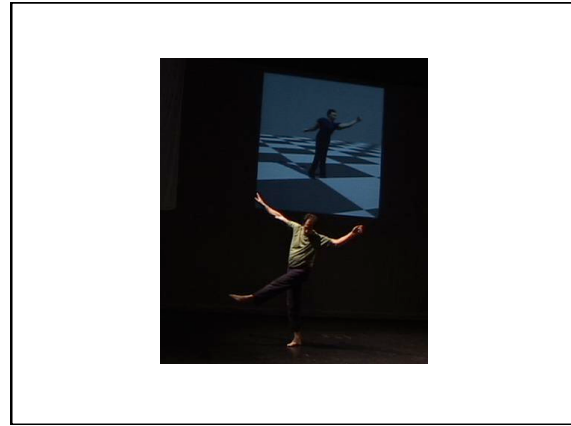
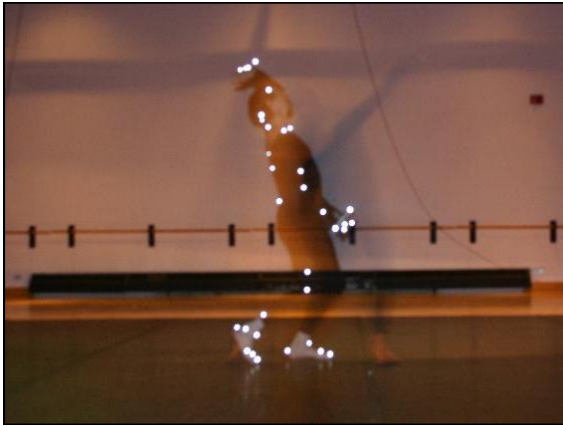


interpolation 1




initial





Con/cantation: (chaotic variations)
A computer-assisted theme and variations performance project



Radcliffe Institute for Advanced Study
Tuesday, April 17th
5pm
Radcliffe Gym
Radcliffe Yard
10 Garden Street
Cambridge, MA 02138
Free Admission

Created by David Capps and Liz Bradley
Video and layout: Angelika von Chamier

Musical algorithms: Josh Stuart
Motion capture and animation: Carnegie Mellon Graphics Laboratory
(Professor Jessica Hodgins, leader; Justin Mastry, motion-capture technician; Mi Mahler, animation and character design)
Cue: David Towlebridge and Evan Sheehan
Inspiration: Diana Dabby

Made possible with support from the Radcliffe Institute for Advanced Study, the National Science Foundation (05-024322), the David and Lucile Packard Foundation, and the Graduate Council on Arts and Humanities at the University of Colorado.

